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6th International Conference on Astrodynamics Tools and Techniques

Dynamical Analysis of Rendezvous and Docking with Very Large Space Infrastructures in Non-Keplerian Orbits

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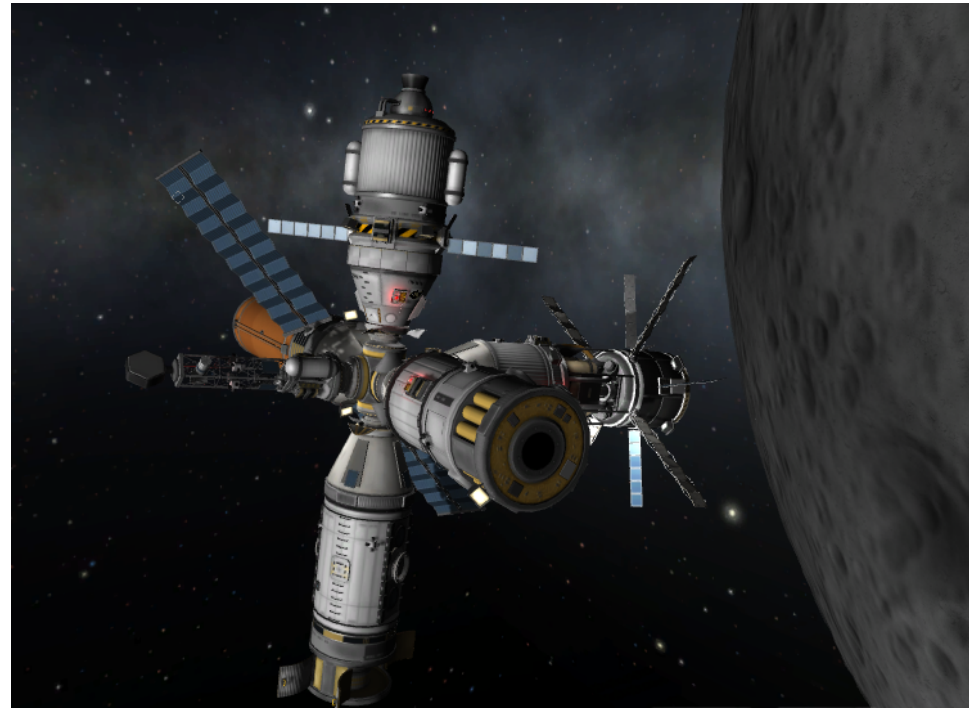
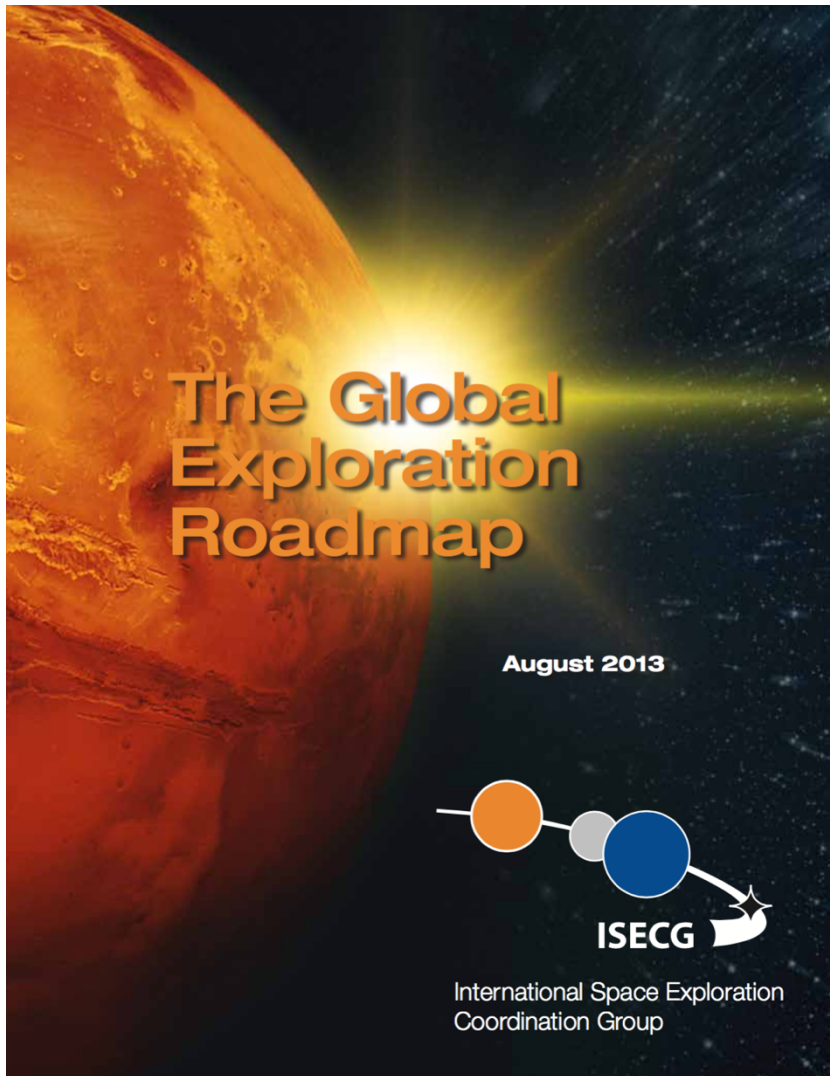
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Framework of Present Work

Framework of Present Work



Courtesy of International Space Exploration Group and Kerbal Space Program.





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Scope of Present Work

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- Definition of a possible rendezvous scenario with a large space infrastructure in Non-Keplerian orbit
 - Dynamical Analysis performed with a simulation tool that includes a coupled orbit-attitude model in a Circular Restricted 3 Body Problem (CR3BP) environment and the flexibility of the structure
- Optimization of rendezvous manoeuvres: transfer and proximity phases
- Preliminary analysis on the effects of flexibility on the coupled dynamics in non-Keplerian orbit
- Preliminary implementation of an astrodynamics tool able to deal with flexible large structures in cis-lunar space





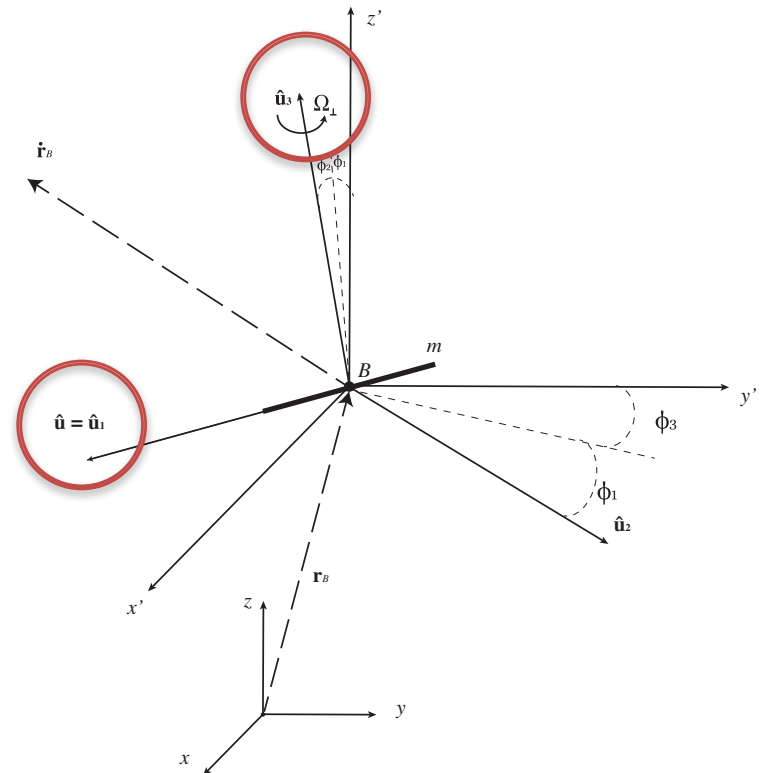
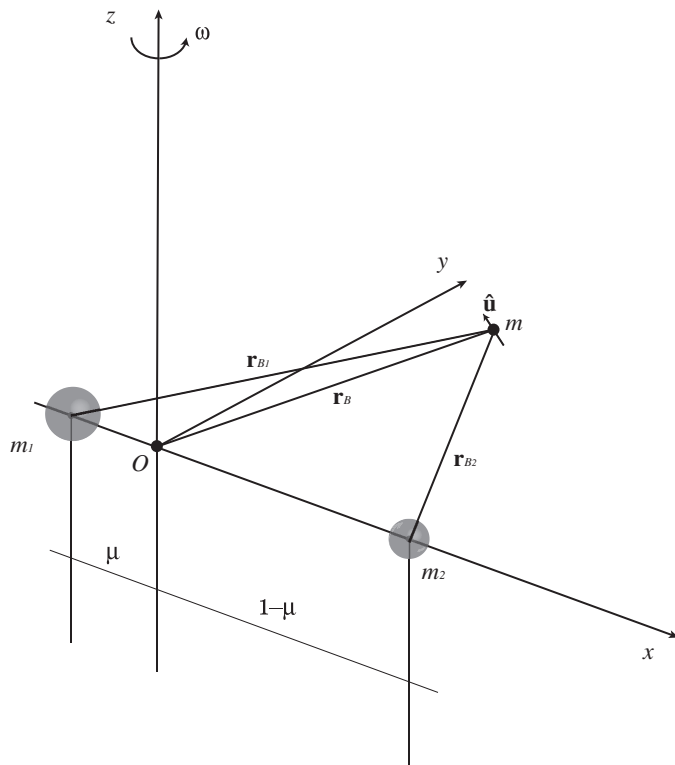
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Theoretical Background

Dynamics in CR3BP

The simulation tool is developed exploiting a “Multi-Body-Friendly” approach; with the idea of further extensions and refinements.

Lagrangian Formulation in the Synodic Reference Frame



Dynamics in CR3BP

$$\mathcal{L} = \mathcal{T} - \mathcal{G}$$

\mathcal{T} kinetic energy

\mathcal{G} generalized potential

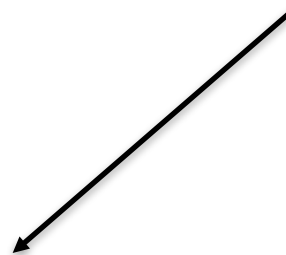
$$\mathcal{T} = \frac{1}{2}m\dot{\mathbf{r}}_B \cdot \dot{\mathbf{r}}_B + \frac{1}{2}\boldsymbol{\Omega} \cdot \mathbf{I}_B \cdot \boldsymbol{\Omega}$$

$$\mathcal{G} = V_{g_1} + V_{g_2} + V_i$$

$$V_{g_i} = -\frac{Gm_i m}{r_{B_i}} + \frac{Gm_i}{r_{B_i}^3} \left[\frac{3}{2}(\hat{\mathbf{r}}_{B_i} \cdot \mathbf{I}_{B_i} \cdot \hat{\mathbf{r}}_{B_i}) - \text{tr}(\mathbf{I}_B) \right]$$

$$V_i = \frac{m}{2}\mathbf{r}_B \cdot [\boldsymbol{\omega} \times (\boldsymbol{\omega} \times \mathbf{r}_B) + 2\boldsymbol{\omega} \times \dot{\mathbf{r}}_B] - \frac{1}{2}\boldsymbol{\omega} \cdot \mathbf{I}_B \cdot (\boldsymbol{\omega} + 2\boldsymbol{\Omega})$$

Effect of finite
dimension of the body



Influence of the Extended Body

The large space infrastructure has been modeled as a rigid rod of length, l , and non-dimensional length $\epsilon_0 = l/r_{12}$

The lagrangian function can be expressed as

$$\mathcal{L} = \mathcal{L}_0 + \epsilon_0^2 \mathcal{L}_2 + \epsilon_0^3 \mathcal{L}_3 + \dots$$

Reducing \mathcal{L} to the main order, the size of the rod disappears from the problem

In this analysis \mathcal{L} has been limited to the second order

The equations of motion are written in non-dimensional form (usual CR3BP non-dimensional formulation)



Equations of Motion

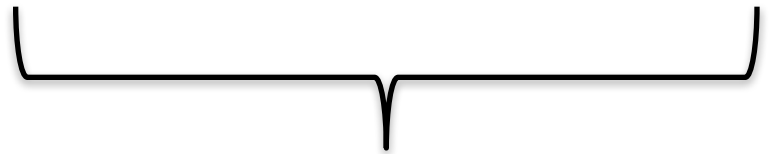
$$\frac{d}{d\tau} \left(\frac{\partial \mathcal{L}_0}{\partial \dot{x}} \right) - \frac{\partial \mathcal{L}_0}{\partial x} = \epsilon_0^2 \frac{\partial \mathcal{L}_2}{\partial x}$$

$$\frac{d}{d\tau} \left(\frac{\partial \mathcal{L}_0}{\partial \dot{y}} \right) - \frac{\partial \mathcal{L}_0}{\partial y} = \epsilon_0^2 \frac{\partial \mathcal{L}_2}{\partial y}$$

$$\frac{d}{d\tau} \left(\frac{\partial \mathcal{L}_0}{\partial \dot{z}} \right) - \frac{\partial \mathcal{L}_0}{\partial z} = \epsilon_0^2 \frac{\partial \mathcal{L}_2}{\partial z}$$

$$\frac{d}{d\tau} \left(\frac{\partial \mathcal{L}_2}{\partial \dot{\theta}} \right) - \frac{\partial \mathcal{L}_2}{\partial \theta} = 0$$

$$\frac{d}{d\tau} \left(\frac{\partial \mathcal{L}_2}{\partial \dot{\varphi}} \right) - \frac{\partial \mathcal{L}_2}{\partial \varphi} = 0,$$



$$\frac{d \mathbf{h}_B}{dt} = \mathbf{m}_B \quad \longleftarrow$$

$$\mathbf{h}_B = I_B \Omega_{\perp} \hat{\mathbf{u}}_3$$

$$\Omega_{\perp} = |\hat{\mathbf{u}} \times \dot{\hat{\mathbf{u}}}| = |\dot{\hat{\mathbf{u}}}|$$

More convenient to describe attitude motion with Newton-Euler formulation:

1-2-3 Euler angles are used

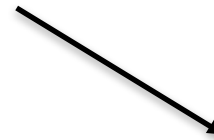


Flexible Dynamics

Flexibility is currently inserted in the model with a lumped parameters technique: spring-mass systems representing pseudo-mode of vibrations.

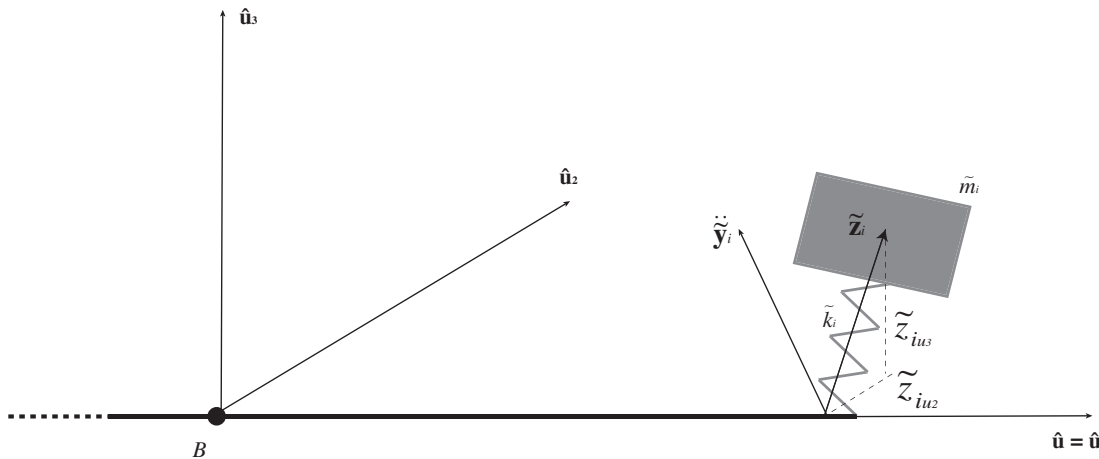
$$\tilde{m}_i(\ddot{\tilde{\mathbf{z}}}_i + \ddot{\tilde{\mathbf{y}}}_i) + \tilde{k}_i\tilde{\mathbf{z}}_i = 0$$

$$\tilde{k}_i = \tilde{\omega}_i^2 \tilde{m}_i$$



Torque exerted on the rod

$$\mathbf{m}_{B_{k_i}} = l_i \hat{\mathbf{u}}_1 \times \tilde{k}_i \tilde{\mathbf{z}}_i$$



$\tilde{z}_{iu1} = 0$: Spring-Mass constrained to be orthogonal to the rod





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Rendezvous Definition

State of the Art

- Autonomous rendezvous in CR3BP is in a preliminary study phase
- Proposed strategies often exploit invariant manifolds: low-cost transfer capabilities
- Existing studies are focused on point-mass spacecrafts (Lizy-Destrex, Murakami et al., Ueda et al.)
- Possible rendezvous strategies (Koon et al).
 - HOI – Halo orbit insertion (Halo To Halo)
 - MOI – Stable manifold orbit insertion (Transfer to Halo)



Proposed Rendezvous Strategy

Departure: the chaser is injected in a unstable manifold of the parking orbit with a first manoeuvre, Δv_1 .

Switching: the chaser is injected, Δv_2 , in the stable manifold of the target's operational orbit. The injection point is at the intersection of unstable and stable manifolds.

Approach: the chaser arrives in proximity of the target and it is moved very close to the operational Halo orbit, Δv_3 .

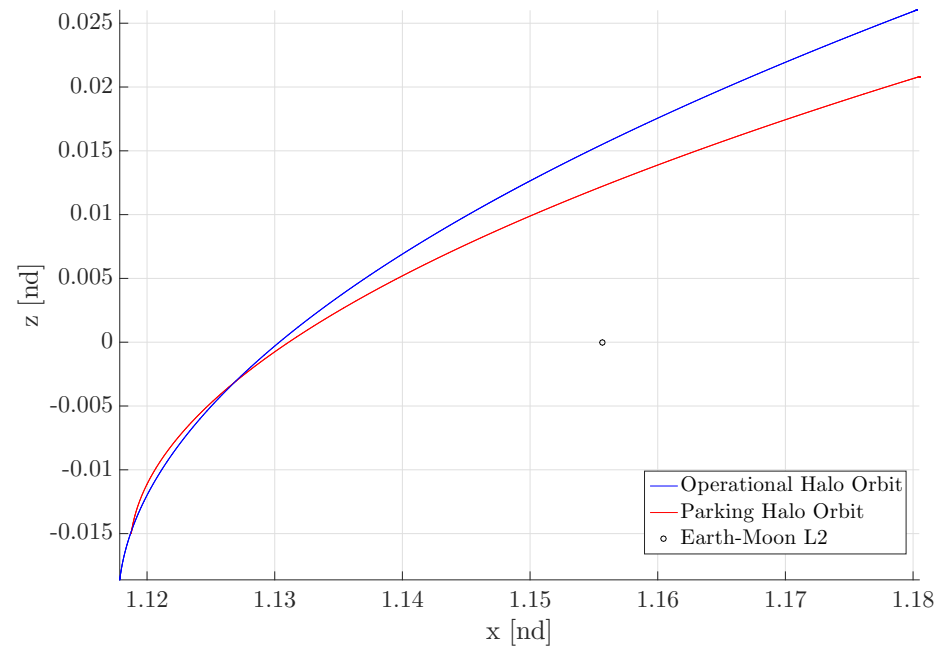
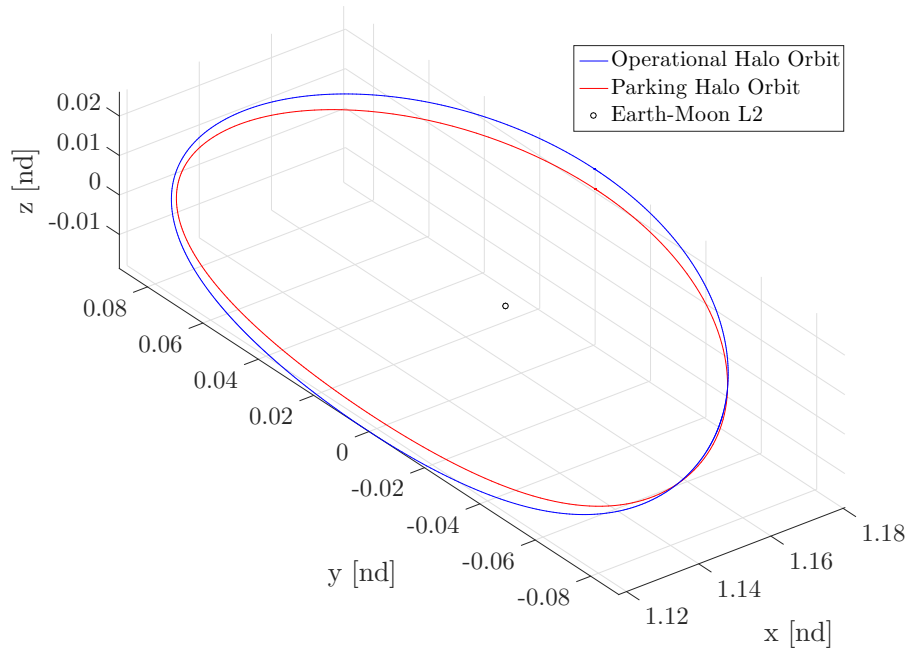
Closing: Δv_4 aligns the chaser with the docking axis of the space station.

Final approach: Δv_5 and Δv_{51} progressively reduce the relative distance between chaser and target. The chaser is maintained aligned with the the docking axis of the target, which is rotating.

Mating: a continuous Δv_6 reduces to zero the relative distance between the two spacecrafts and brings the chaser at the docking port



Operational and Parking Halo Orbit



Orbit	A_z [km]	T [D]	C [-]
Operational Halo	10000	14.808	3.149
Parking Halo	8000	14.813	3.150



Rendezvous Scenario

- Intersection between manifolds (switching point) is assumed to be in the space between Moon and Earth: $x_{sp} = 1 - \mu$
 - The position of the switching point is a free parameter in the rendezvous optimization tool
 - The region where the algorithm look for the switching point is an input from the user
 - Hypothetical mission with a cyclic chaser between parking and operational Halo orbit. (Possibility to encounter a cargo coming from Earth, Moon or LLO)
- Chaser is a point mass
- Target (space station) is an extended body

	m_T [kg]	l_T [m]
Target	300000	100





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Rendezvous Simulation

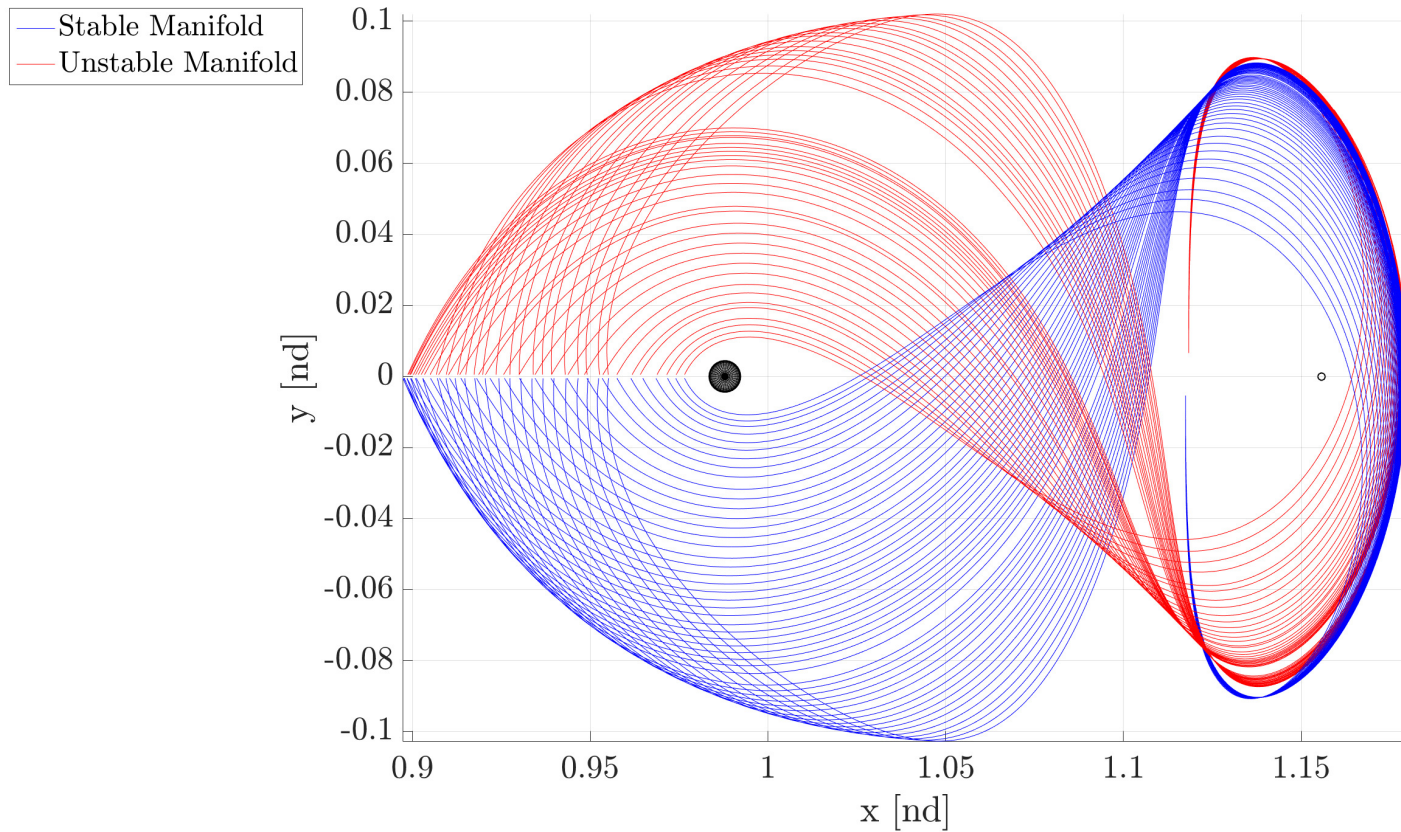
Heteroclinic Transfer

Target and chaser are approximately phased

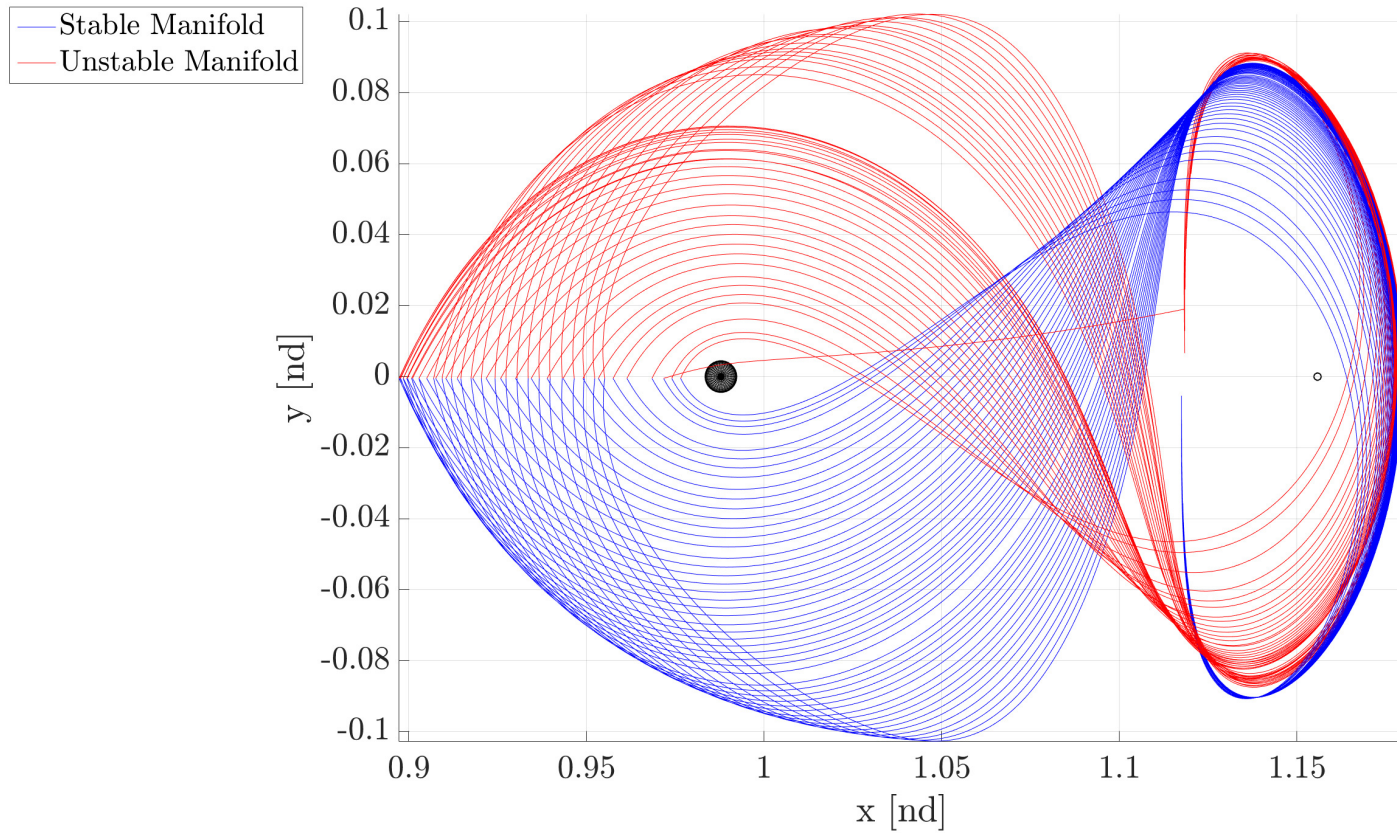
- This condition can be always obtained with a preliminary phasing manoeuvre
- Proximity operations are able to correct errors in phasing
- Invariant manifolds are computed
- Intersection are analysed on a Poincarè section → Sub-Optimal Transfers
- Sub-Optimal transfers are corrected and position continuity along the transfer is enforced
- Best Sub-Optimal transfer is a first guess for the optimization algorithm
- Optimization process varies state vector at the beginning of the transfer → $\min (\Delta v_{transfer})$ → Optimal transfer



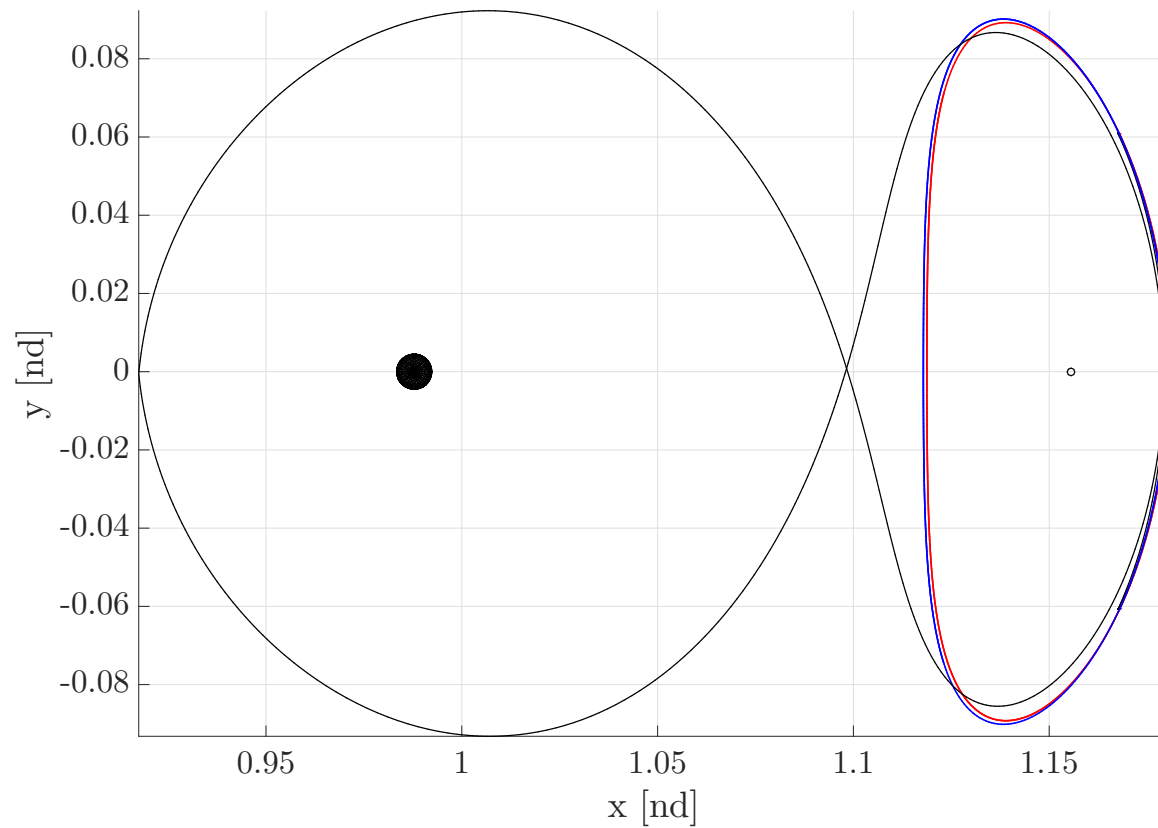
Heteroclinic Transfer



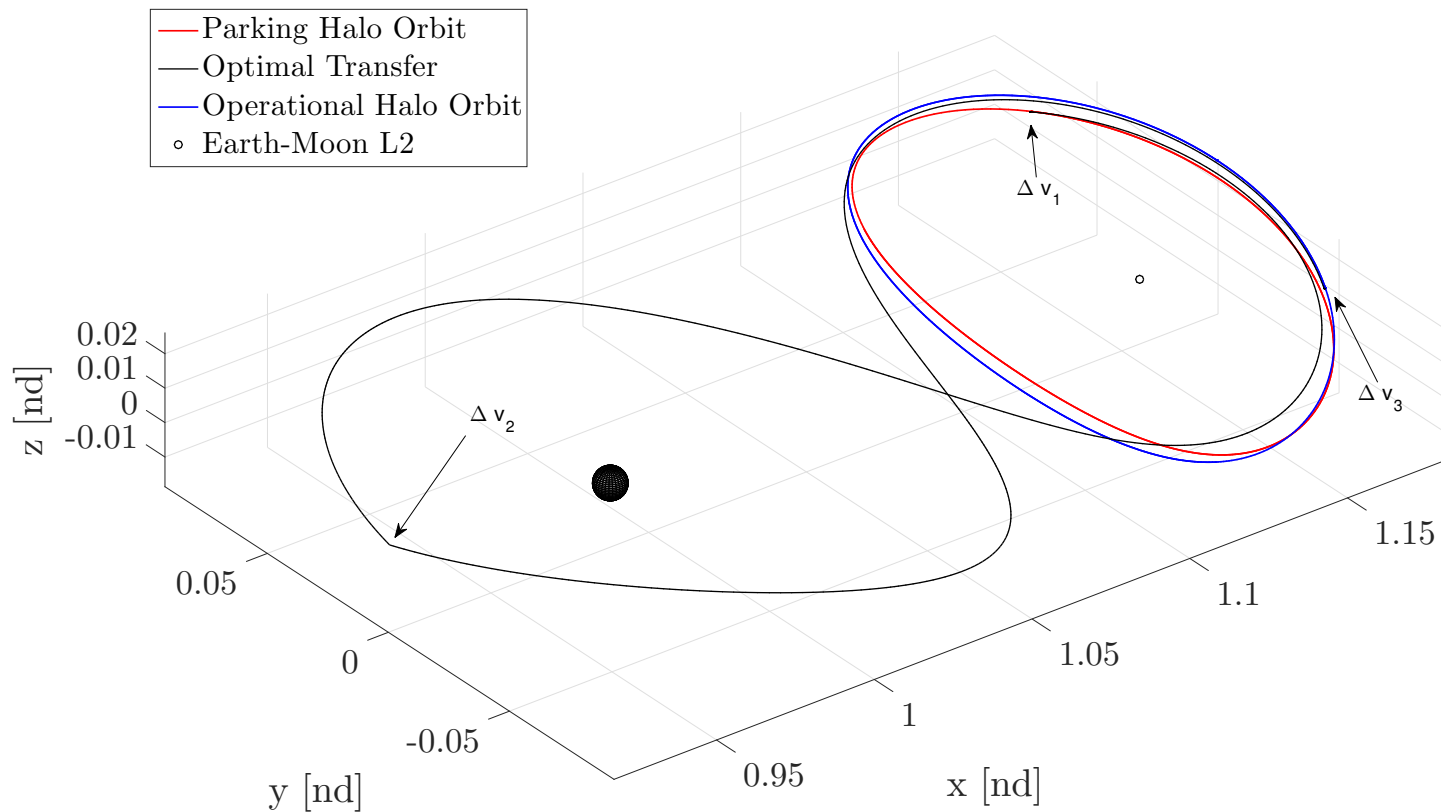
Heteroclinic Transfer



Heteroclinic Transfer



Heteroclinic Transfer



$t_{transfer}$ [d]	Δv_1 [m/s]	Δv_2 [m/s]	Δv_3 [m/s]	$\Delta v_{transfer}$ [m/s]
26.14	5.49	152.29	0.51	158.29

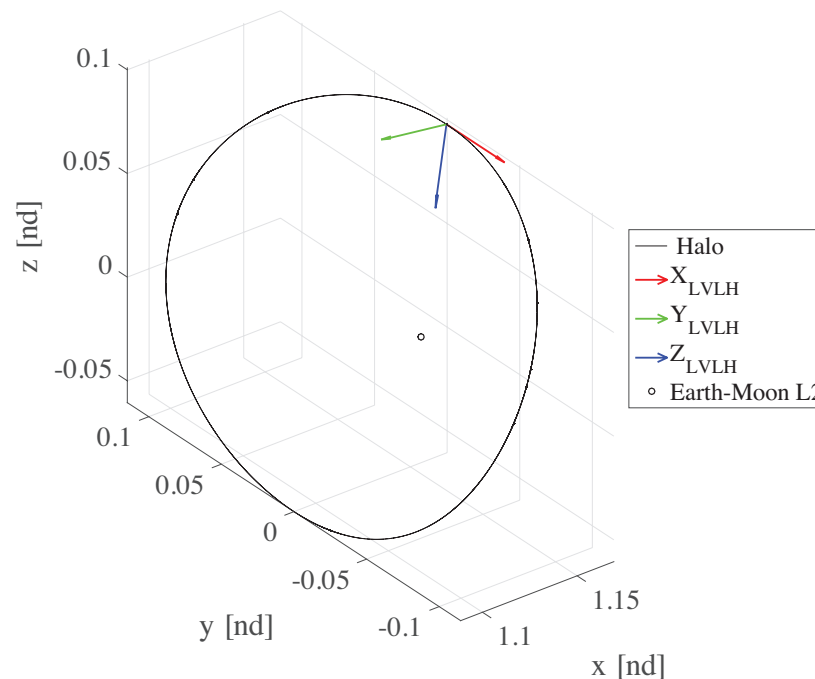


Proximity Operations

The transfer is assumed to be concluded when:

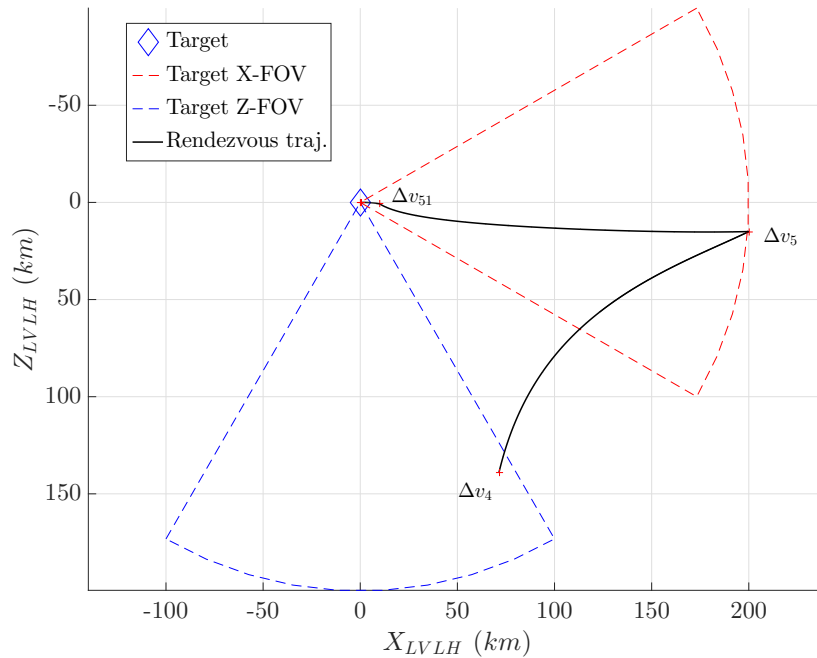
- Relative distance is $\sim 10^2$ km
- Chaser in view along Z_{LVLH}

Proximity operations are analysed in EML2-LVLH frame

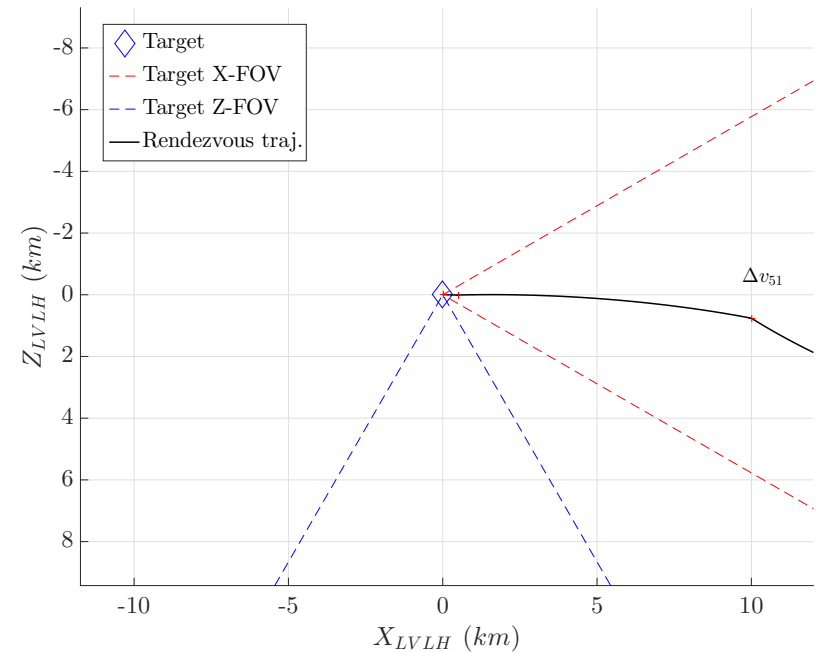


Proximity Operations

Closing Phase



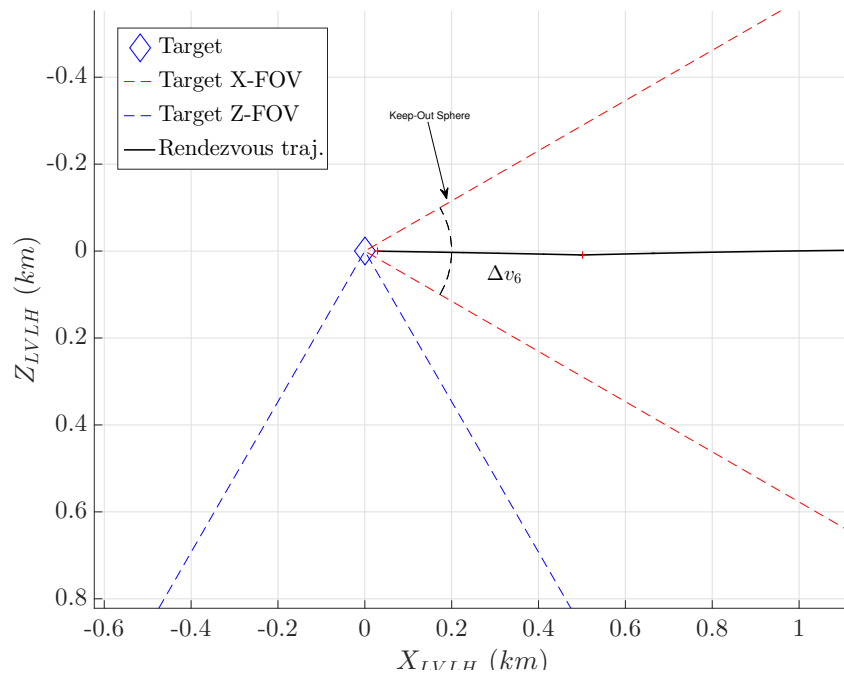
Final Approach Phase



$t_{proximity}$ [d]	Δv_4 [m/s]	Δv_5 [m/s]	Δv_{51} [m/s]
3.23	1.27	3.41	2.52

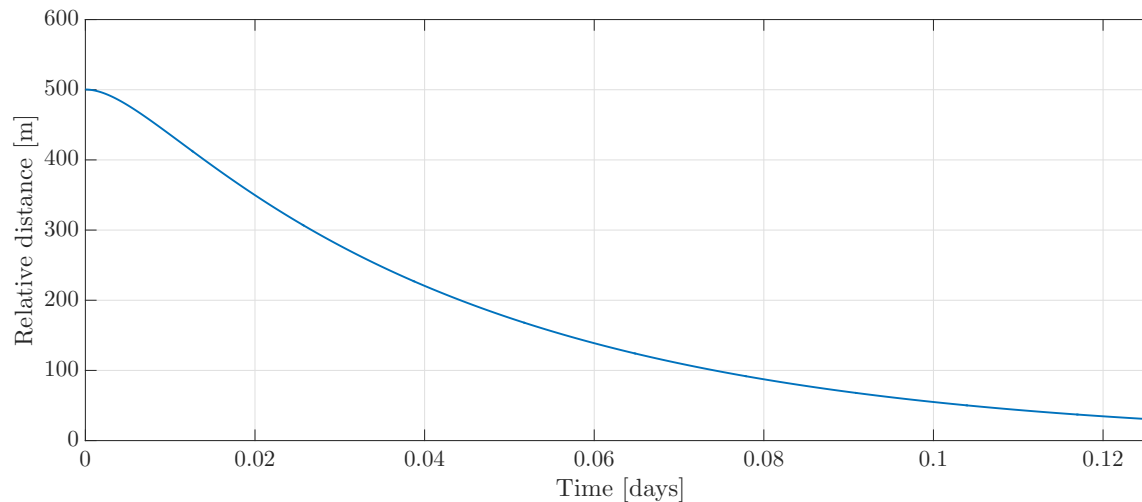


Mating Phase



Continuous thrust before docking

LQR trajectory on linearised CR3BP: target dynamics is computed with the flexible orbit-attitude model



t_{mating} [h]	Δv_6 [m/s]
~3	0.44

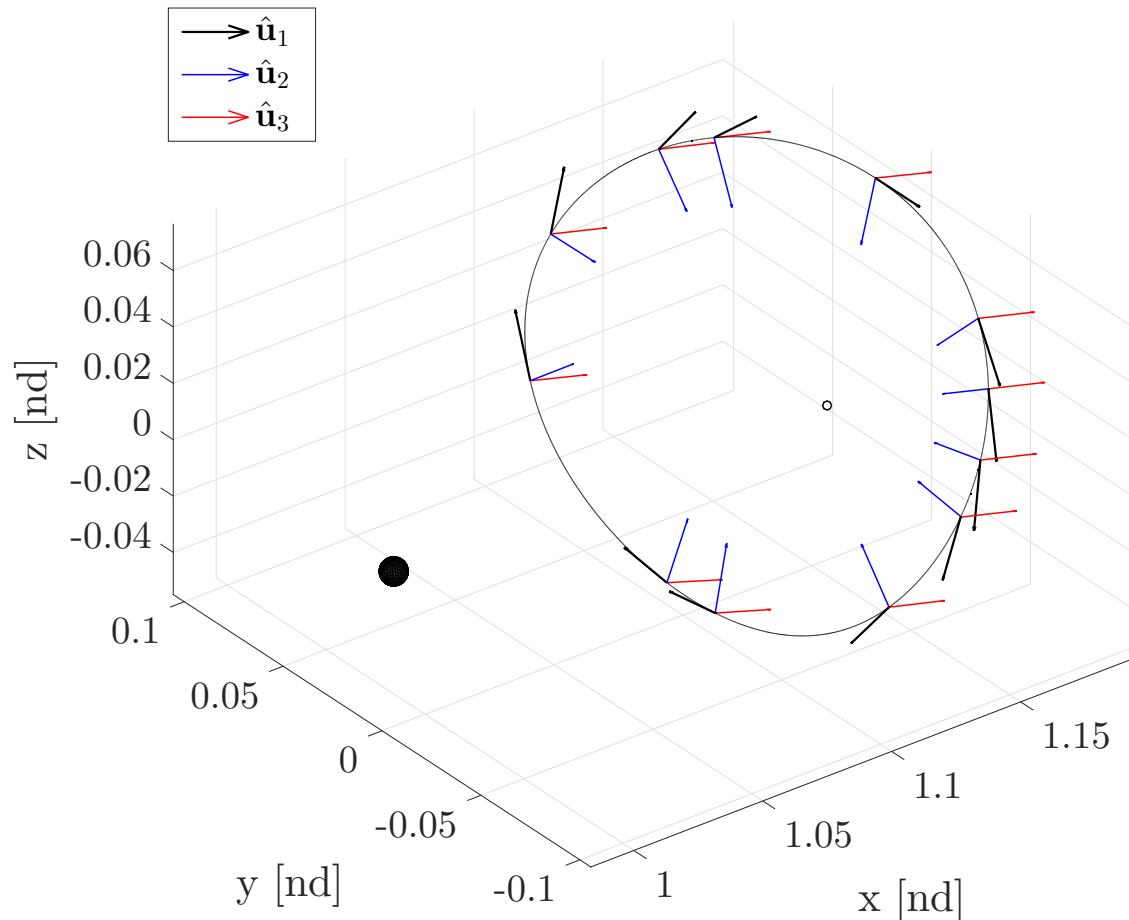




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Flexible Orbit-Attitude Analysis

Orbit-Attitude Motion



Quasi-periodic attitude motion: almost 1 rotation in the first orbital period

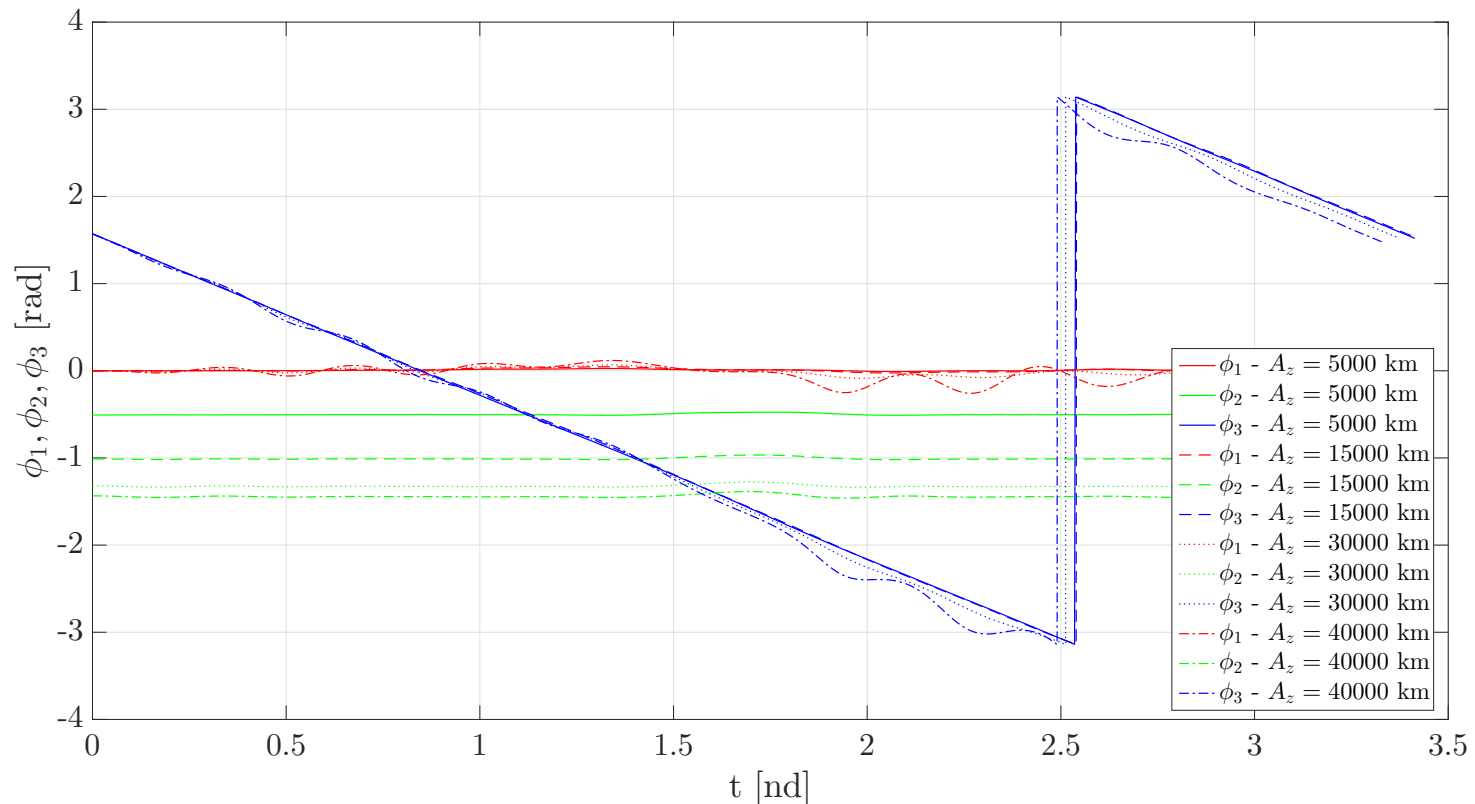
Attitude dynamics (in terms of stability of the periodic motion) is very sensitive to orbital dynamics



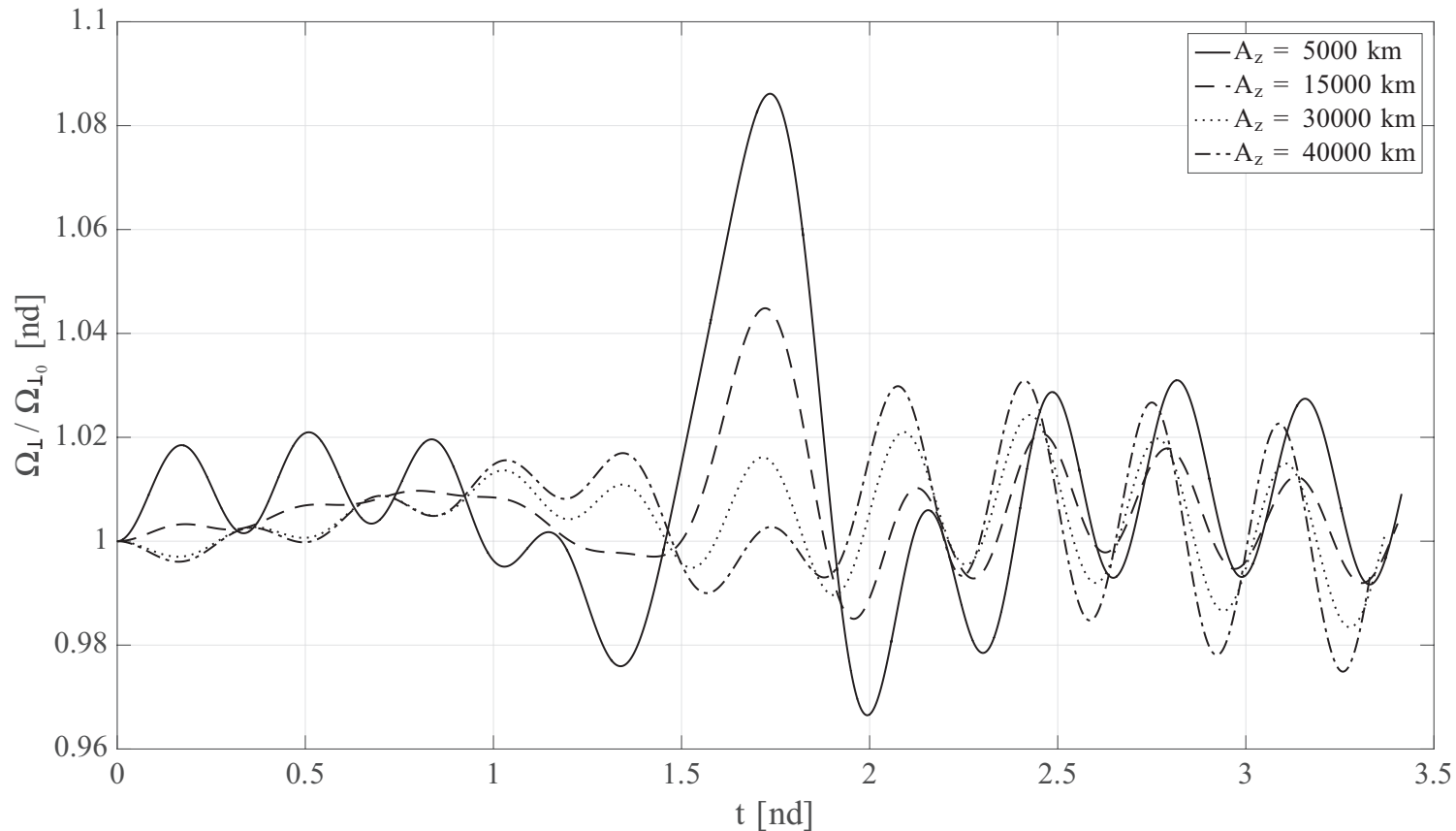
Orbit-Attitude Dynamics Effect on Flexible Dynamics

Influence of the orbital frequency on flexible (spring-mass) frequencies.

- Spring-Mass have $\tilde{\omega} = 50 [-]$
- Halo orbits with $A_z = 5000, 15000, 30000, 40000$ [km]



Orbit-Attitude Dynamics Effect on Flexible Dynamics



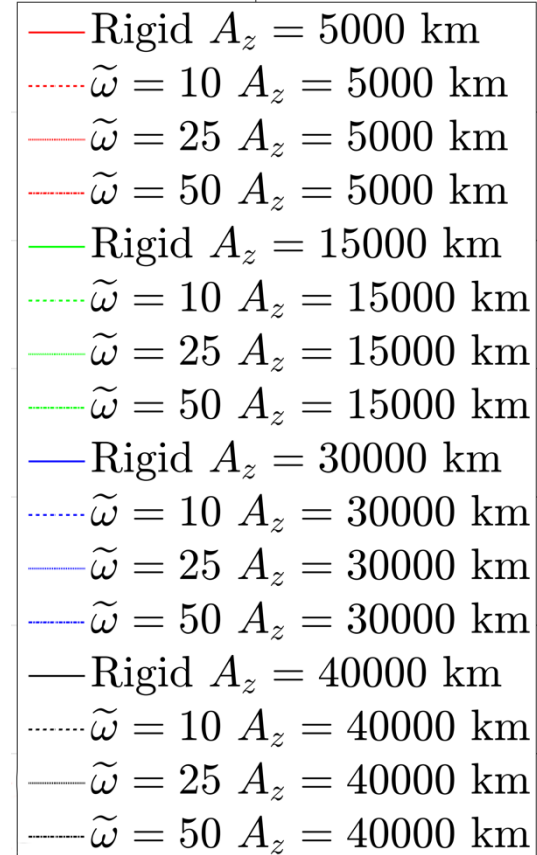
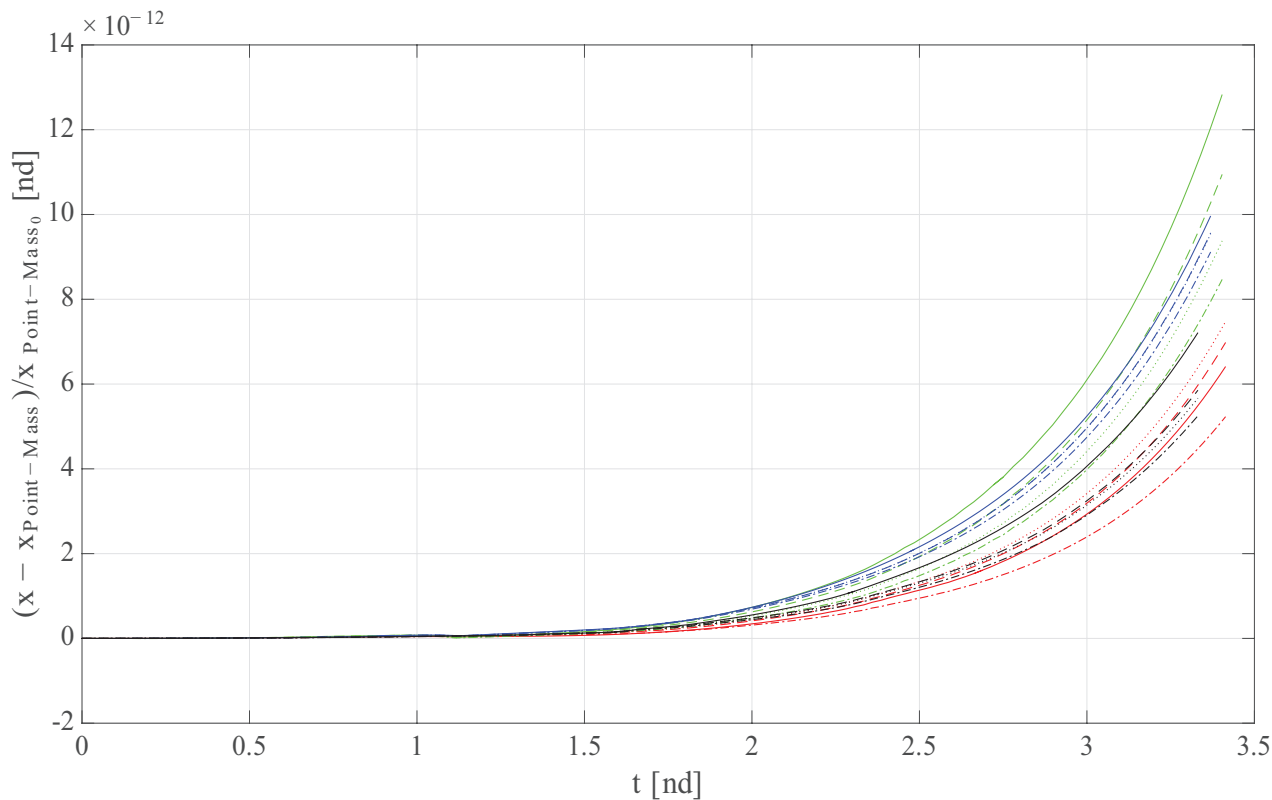
Orbits with large A_z have particular influence on orientation angles, while the others strongly act on Ω_{\perp}



Natural Frequencies Effect on Orbit-Attitude Dynamics

Same Halo orbits of the previous analysis but different $\tilde{\omega}$

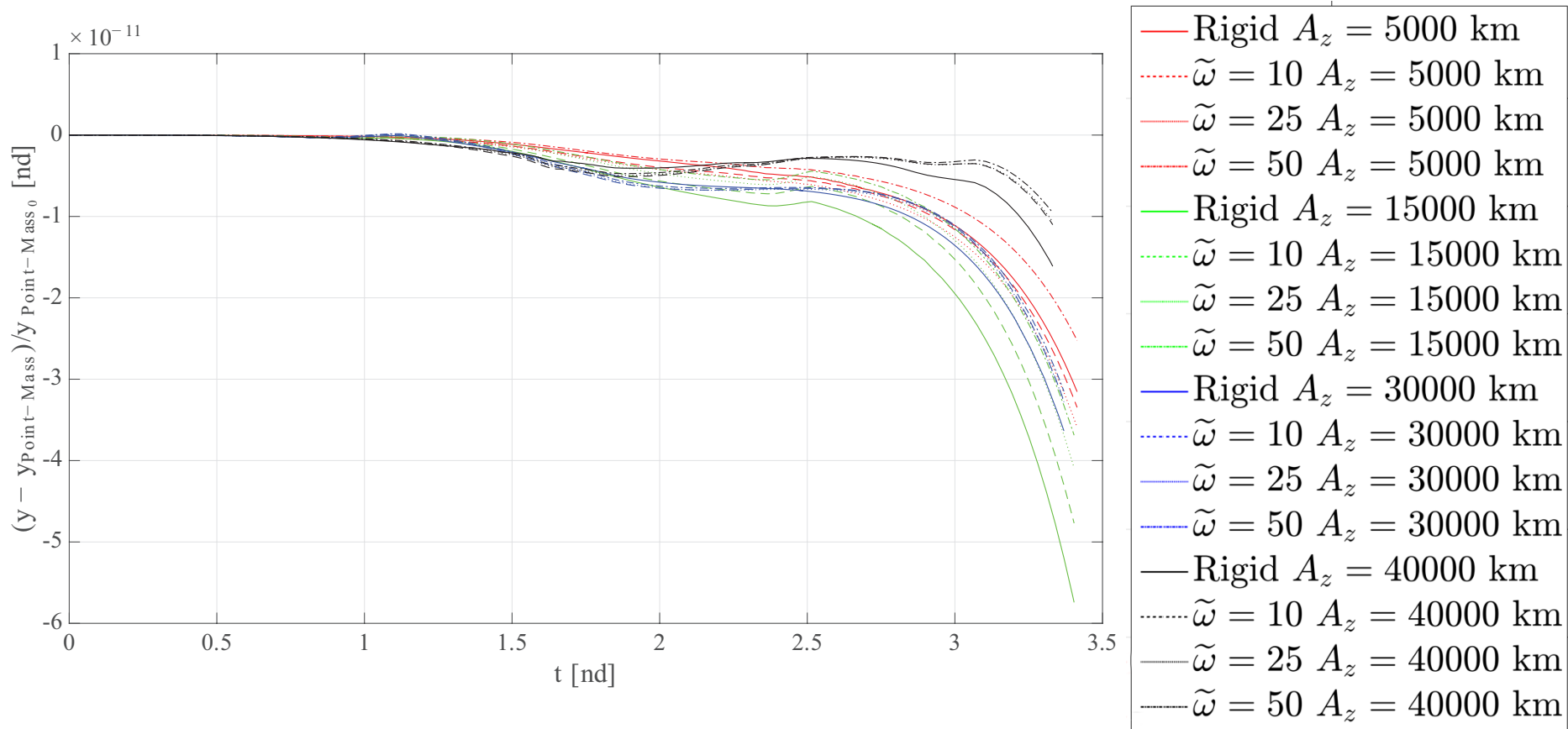
- Error with respect to point-mass dynamics



Natural Frequencies Effect on Orbit-Attitude Dynamics

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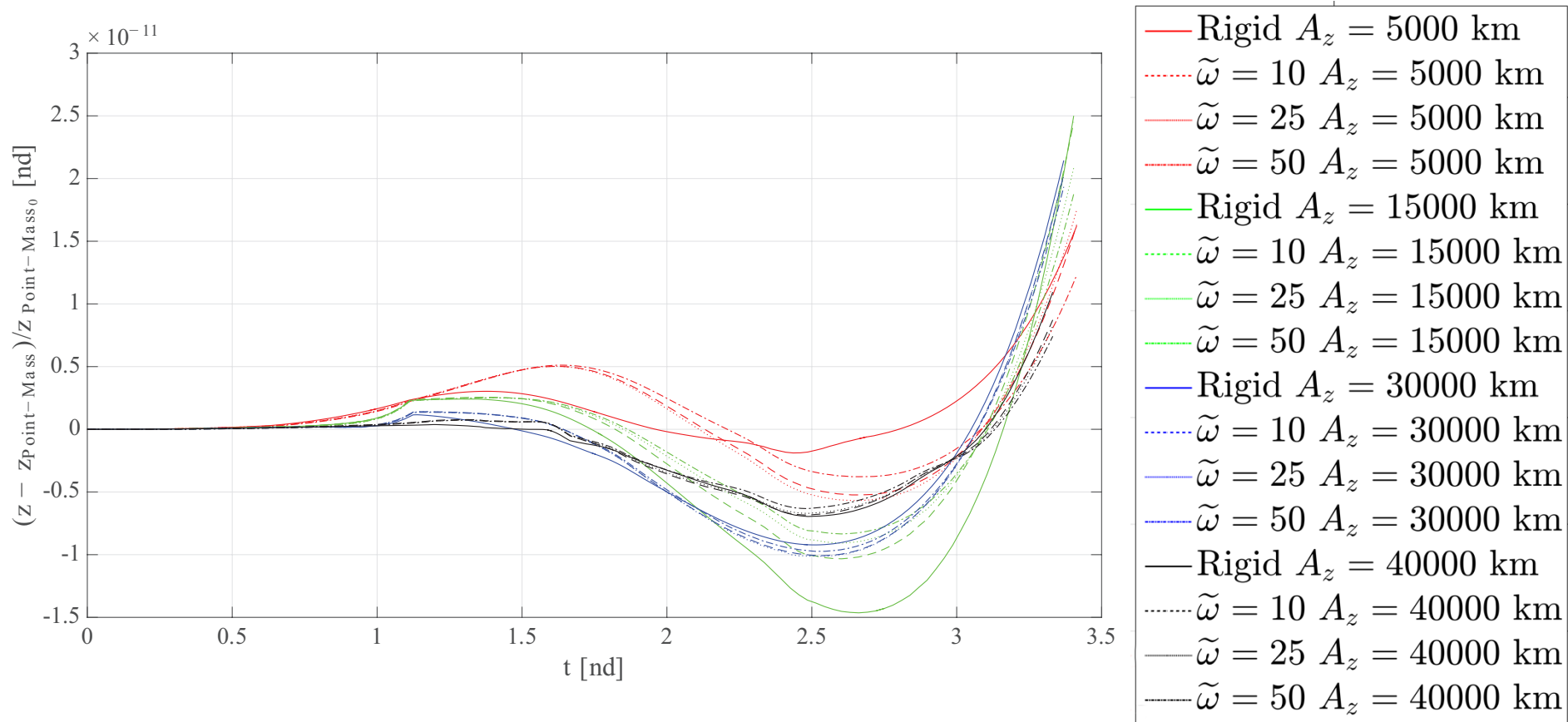
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Natural Frequencies Effect on Orbit-Attitude Dynamics

Same Halo orbits of the previous analysis but different $\tilde{\omega}$

- Error with respect to point-mass dynamics





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Concluding Remarks

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- The effects of flexibility on orbit-attitude dynamics (and viceversa) must be analysed isolating each single effect
 - Outline meaningful trends and particular coupling effects
- The effects of flexible orbit-attitude dynamics cannot be neglected
 - In particular for attitude stability and proximity operations
- Rendezvous with a large and flexible space infrastructure in Non-Keplerian orbit deserves particular attention
 - Refined models must be used to simulate the dynamics
- Cyclic mission with a passage on the Earth side of the Moon is feasible in terms of needed Δv



Future Works

- Enhance the flexibility modelling approach
 - Distributed parameters technique
 - Refined lumped parameters approach will be compared with distributed parameters technique
 - Assess the fidelity of the structural model
- Increase the fidelity of simulations
 - Disturbing phenomena will be included (SRP, 4th body ...)
- Add more structural elements to the large flexible body
 - Multi-body approach





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