

Electron modeling in hybrid plasma thruster plume simulations

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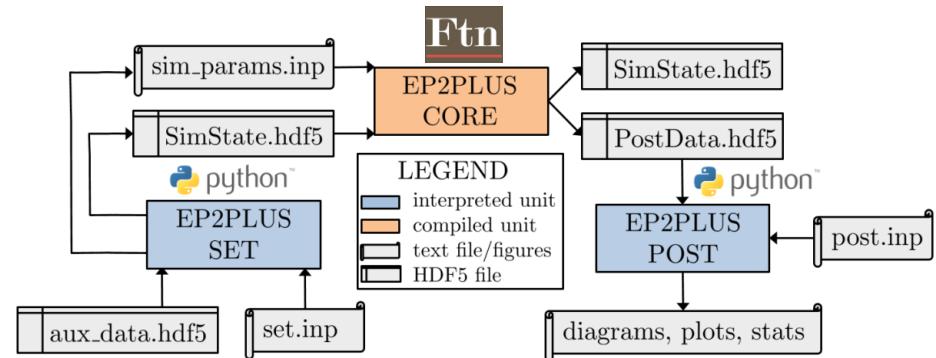
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Contents

- EP2PLUS: Extensible Parallel Plasma PLUme Simulator
- The magnetized electron fluid model
- Benchmark simulations for the electron model
 - Gridded ion thruster plume neutralization
 - Plasma plume expansion with a uniform background magnetic field
- Deformed mesh algorithms and applications
- Conclusions and future studies

EP2PLUS: Extensible Parallel Plasma PLUme Simulator

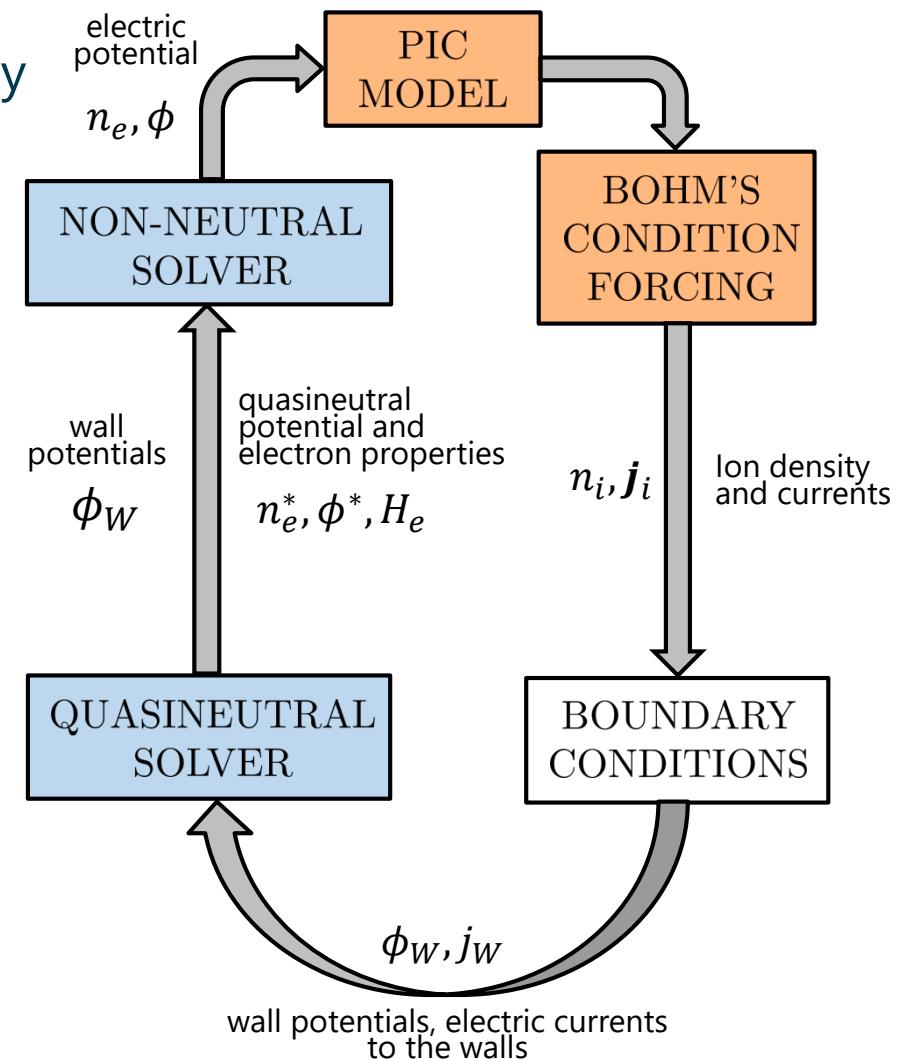
- Hybrid PIC/fluid code
 - PIC ions/neutrals, fluid electrons
- 3D code → asymmetric physics simulation
 - Interaction of plasma thruster plume with the satellite and any downstream object
 - Plume from a cluster of thrusters
 - Asymmetric magnetic field effects
- Industry-level standards
 - HDF5 format I/O files
 - Parallelization with OpenMP
 - Test Driven Design (TDD)
- Most distinguishing features:
 - Electron model enables the computation of both **electric currents** and **magnetic field effects** (ongoing)
 - **Non-neutral code** → interaction with the satellite, peripheral plume regions
 - **Deformed structured meshes** → interaction with cylindrical debris, simulation of plasma thruster discharges, etc...
 - **Easily adaptable to full-PIC**



EP2PLUS: the hybrid loop

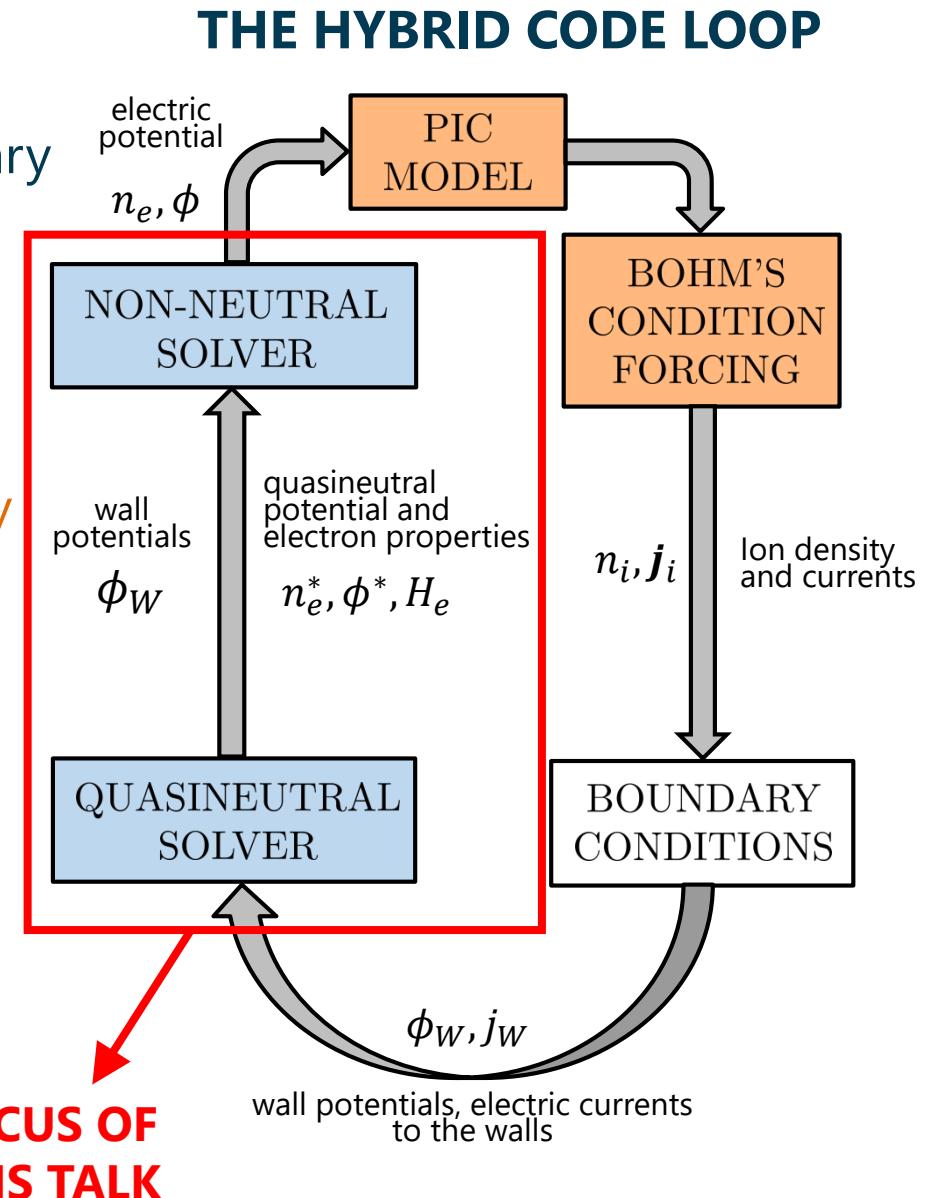
1. PIC sub-model provides ion density and currents
 - Bohm's condition is enforced if necessary
2. Sheath and circuit model compute the electric current and potential at the walls
 - Dielectric and conductive walls
3. Quasi-neutral solver computes the new electric potential, assuming quasineutrality and boundary electric currents
 - Collisions and magnetization can be taken into account as a part of the Bernoulli function $-H_e/e$
4. The electric potential is selectively recomputed with the non-linear Poisson solver in non-neutral regions
 - Non-neutrality criterion

THE HYBRID CODE LOOP



EP2PLUS: the hybrid loop

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The electron fluid model

➤ Assumptions

- Stationary plume properties: $\frac{\partial}{\partial t} = 0$
- Inertialess electrons: $m_e \ll m_i$
- Barotropic electron fluid: $T_e = T_e(n_e) \rightarrow$ BAROTROPY GRADIENT $\nabla h_e = \nabla p_e/n_e \rightarrow h_e = h_e(n_e)$
- Isotropic electrons: $\bar{P}_e = p_e \bar{I}$

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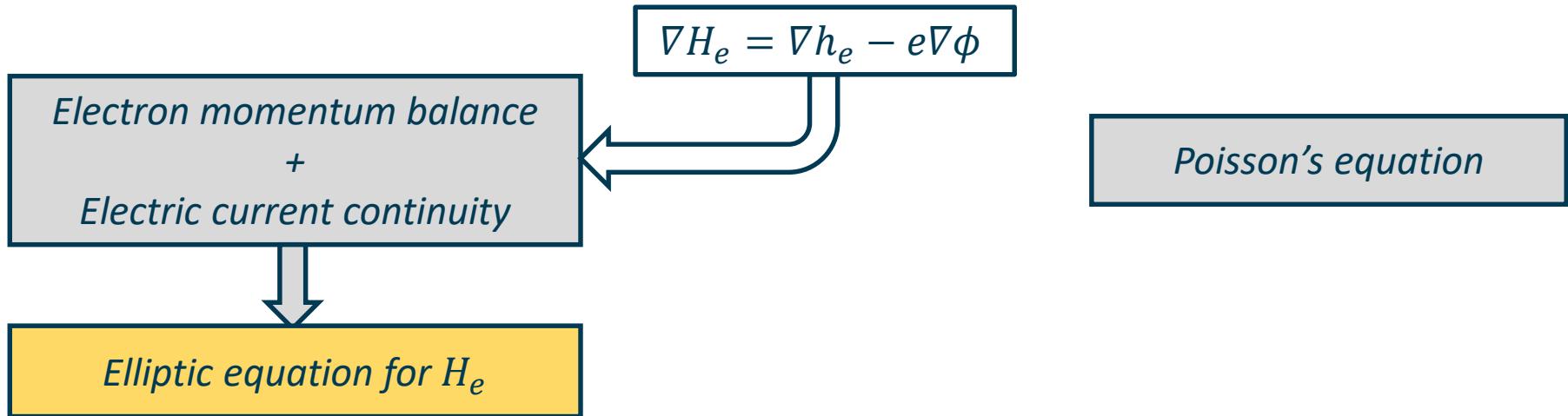
Electron momentum balance
+
Electric current continuity

Poisson's equation

The electron fluid model

➤ Assumptions

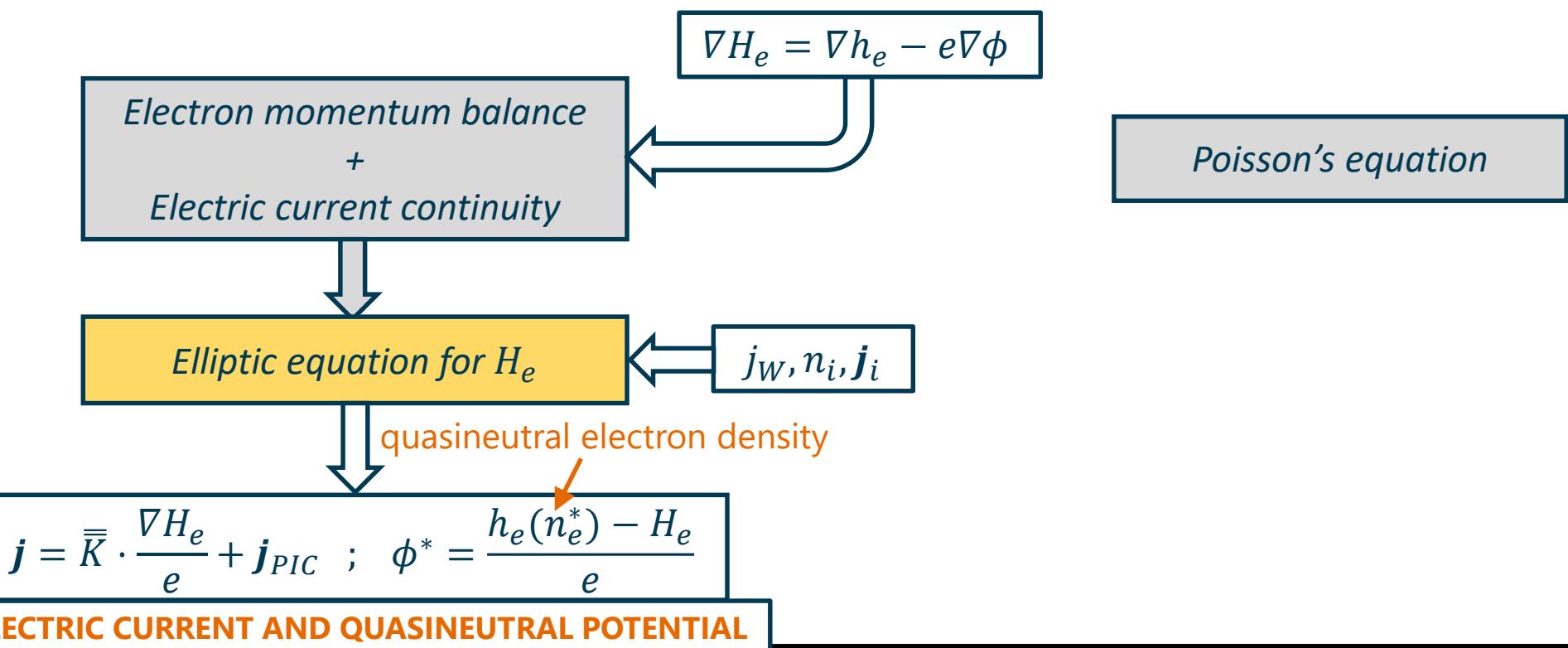
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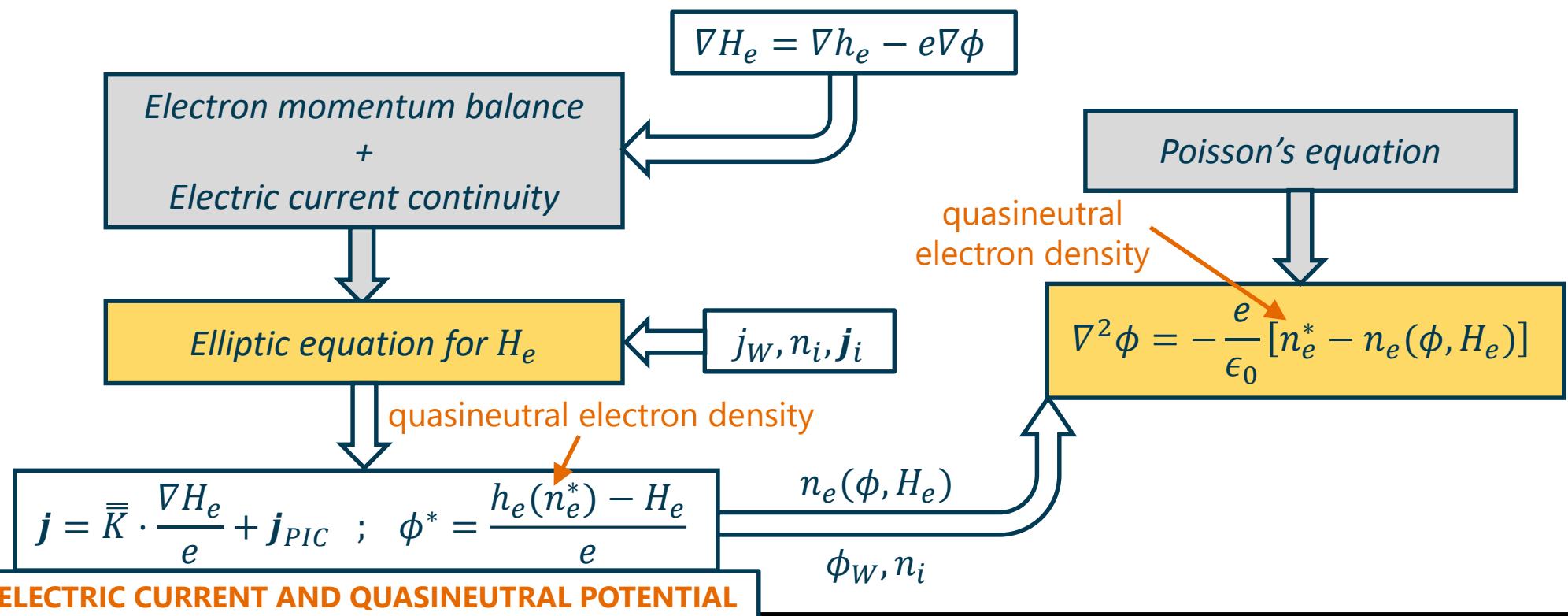
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BAROTROPY GRADIENT

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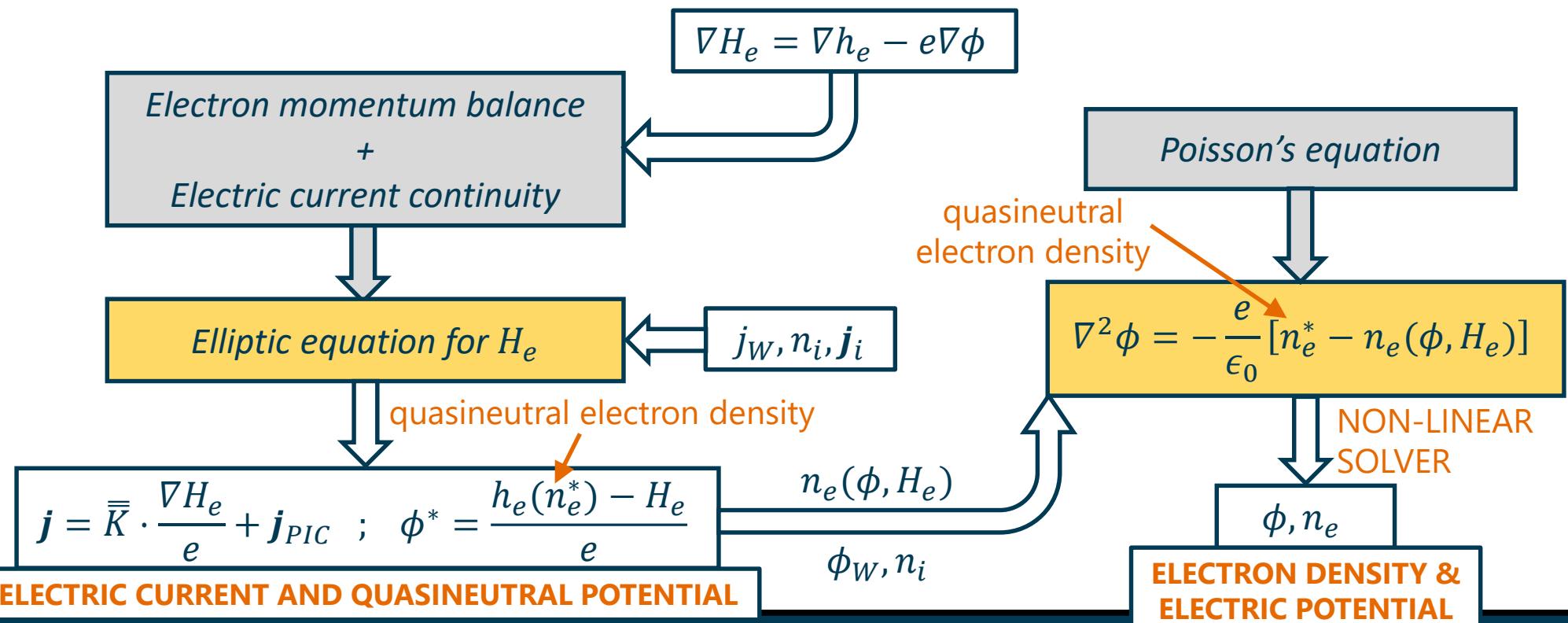
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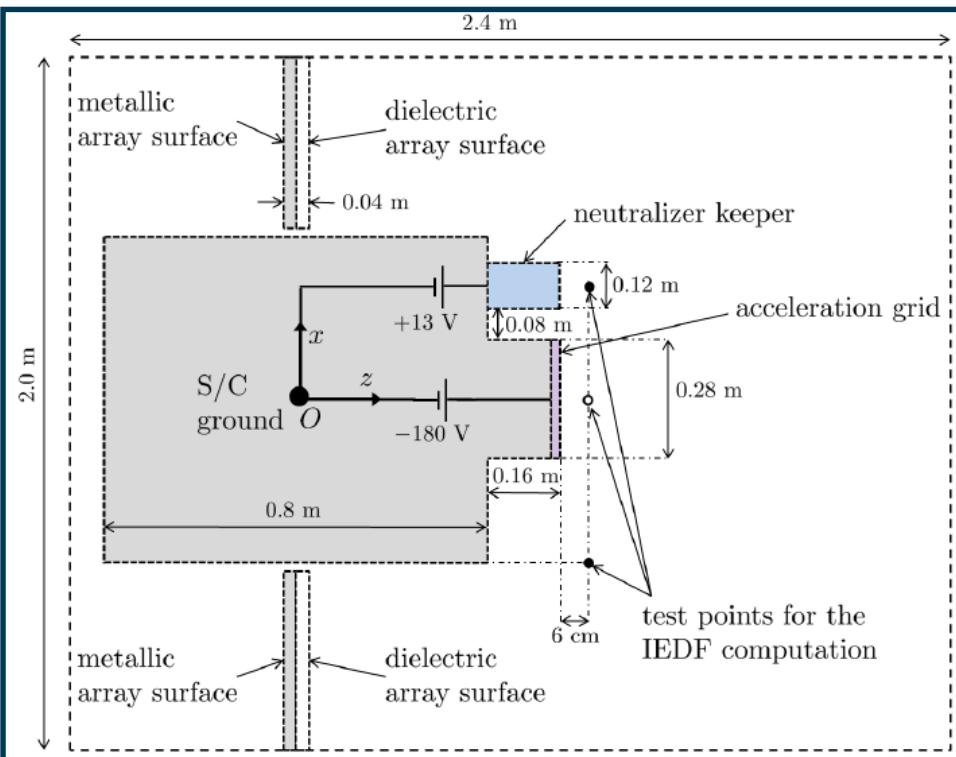
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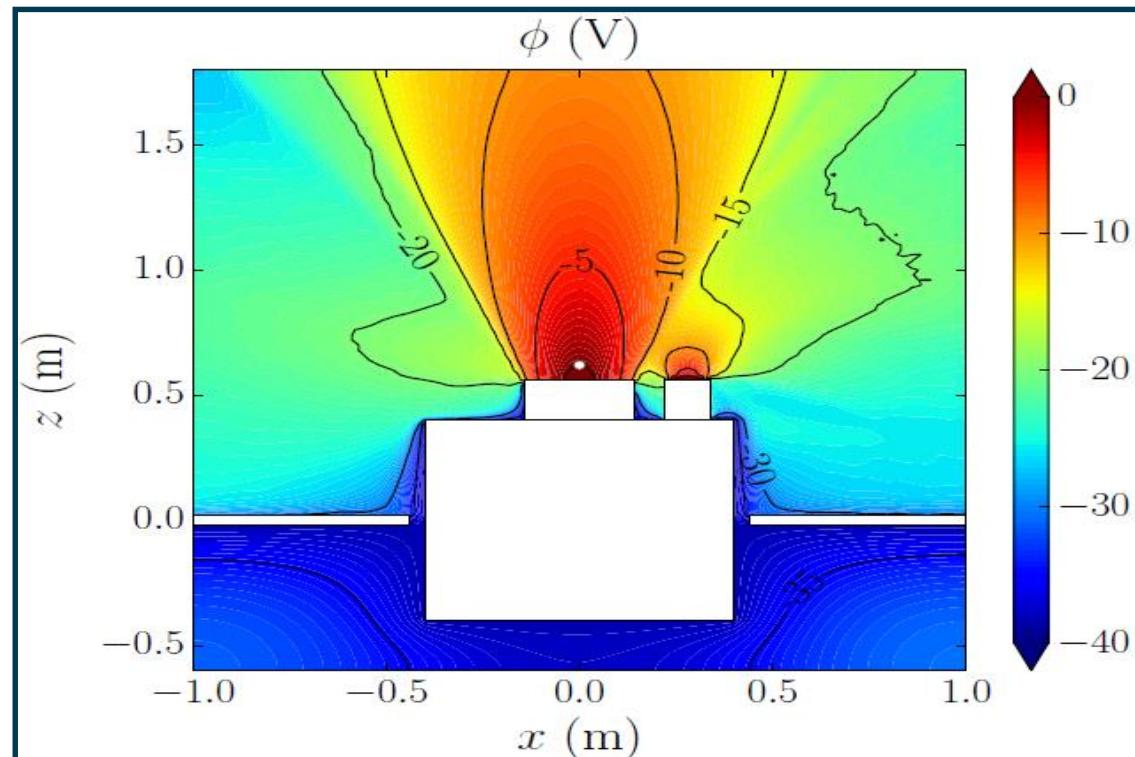
GIT plume neutralization (1)

- Study of the interaction of a plasma plume from a GIT with the satellite and an external neutralizer (**Cichocki PSST 2017**)
 - NSTAR thruster and a neutralizer
 - No magnetic field

EQUIVALENT CIRCUIT



ELECTRIC POTENTIAL

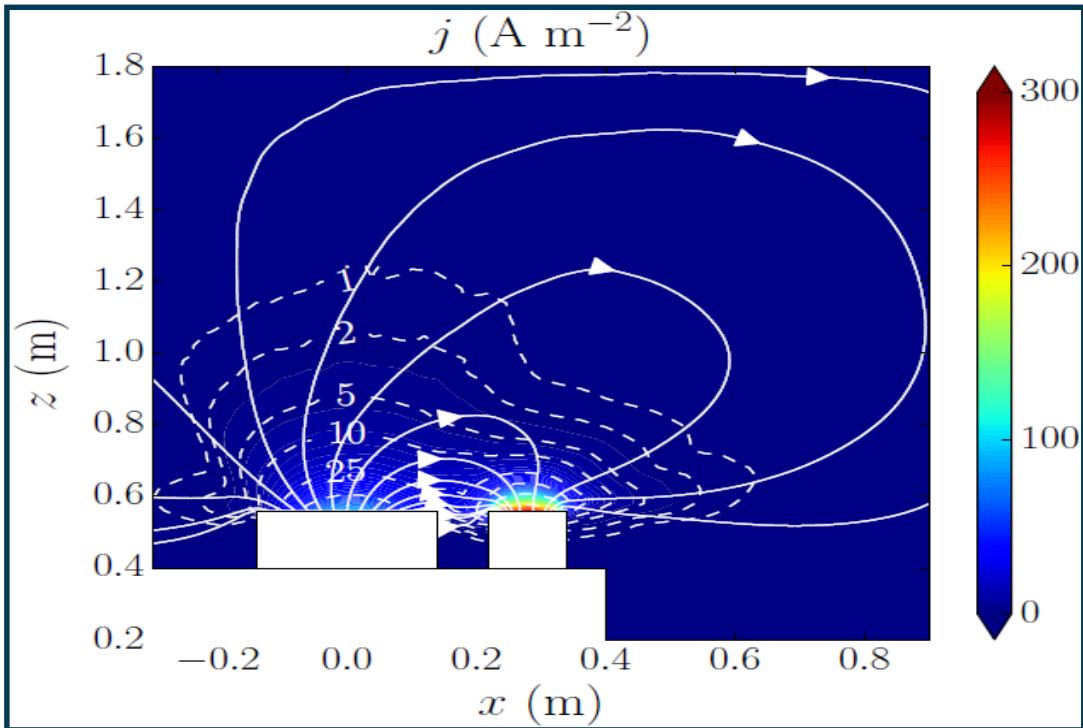


GIT plume neutralization (2)

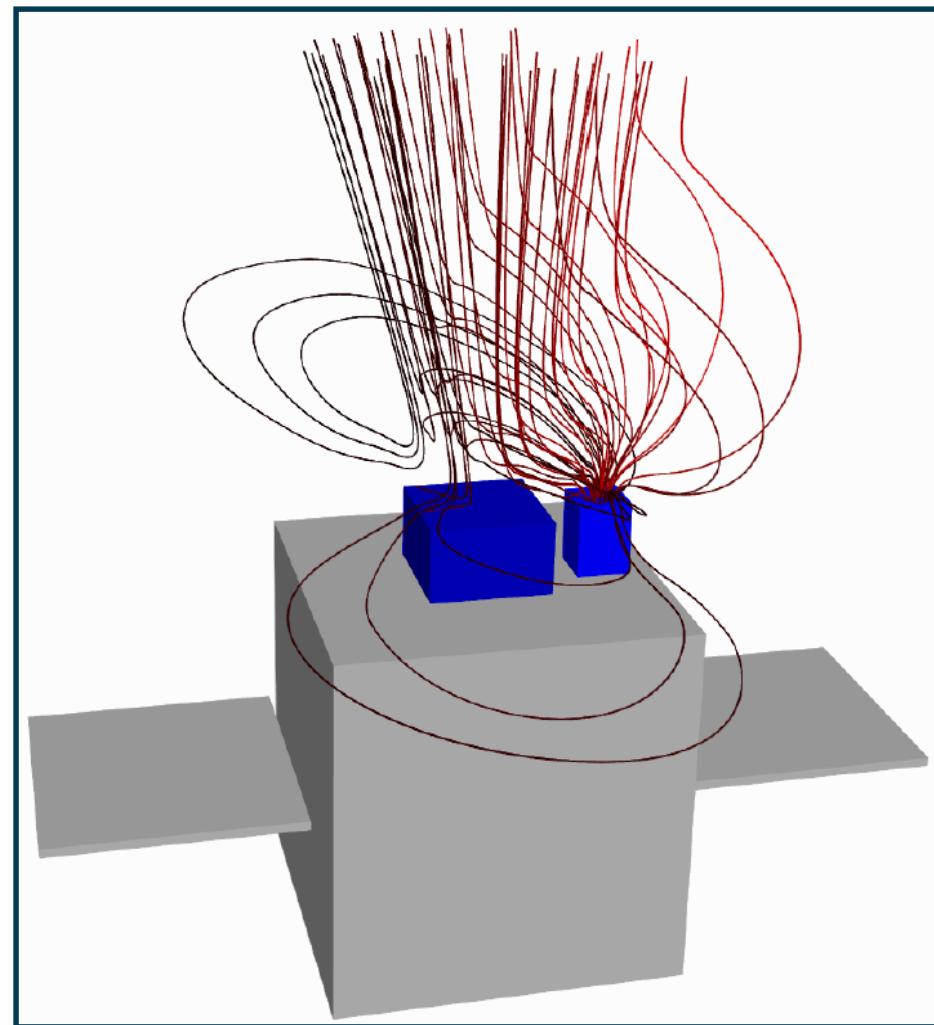
➤ GIT current neutralization was studied:

- Plume current negligible after 0.5 m
- 3D electron streamlines follow minimum resistance path

TOTAL ELECTRIC CURRENT



3D ELECTRON STREAMLINES



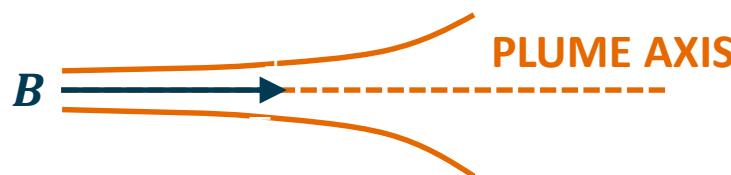
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Plasma plume with uniform magnetic field (1)

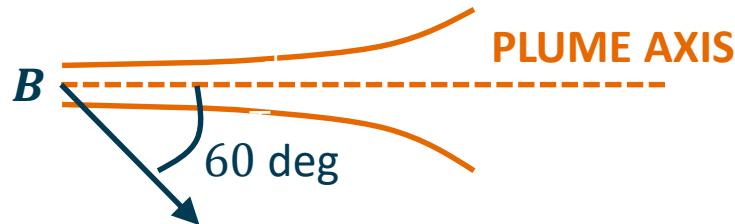
- For a typical plasma thruster plume, various magnetic fields were considered (**Cichocki SP 2018**):

- Unmagnetized case $B = 0$ for reference
- Plume with axial magnetic field



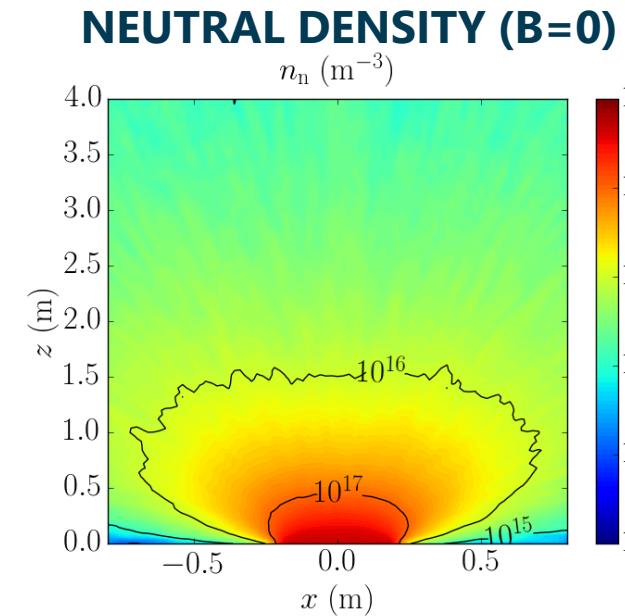
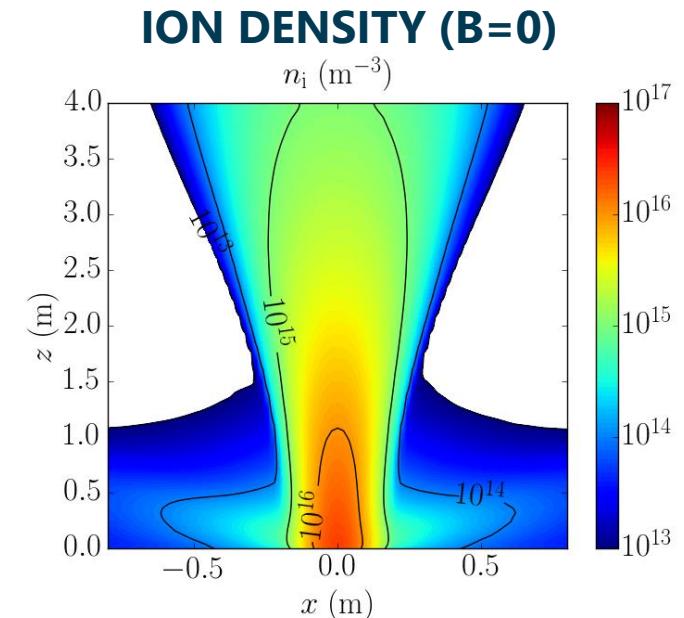
**SYMMETRIC
COMPRESSION,
NOT SHOWN HERE**

- Plume with magnetic field forming an angle of 60 deg with the plume axis



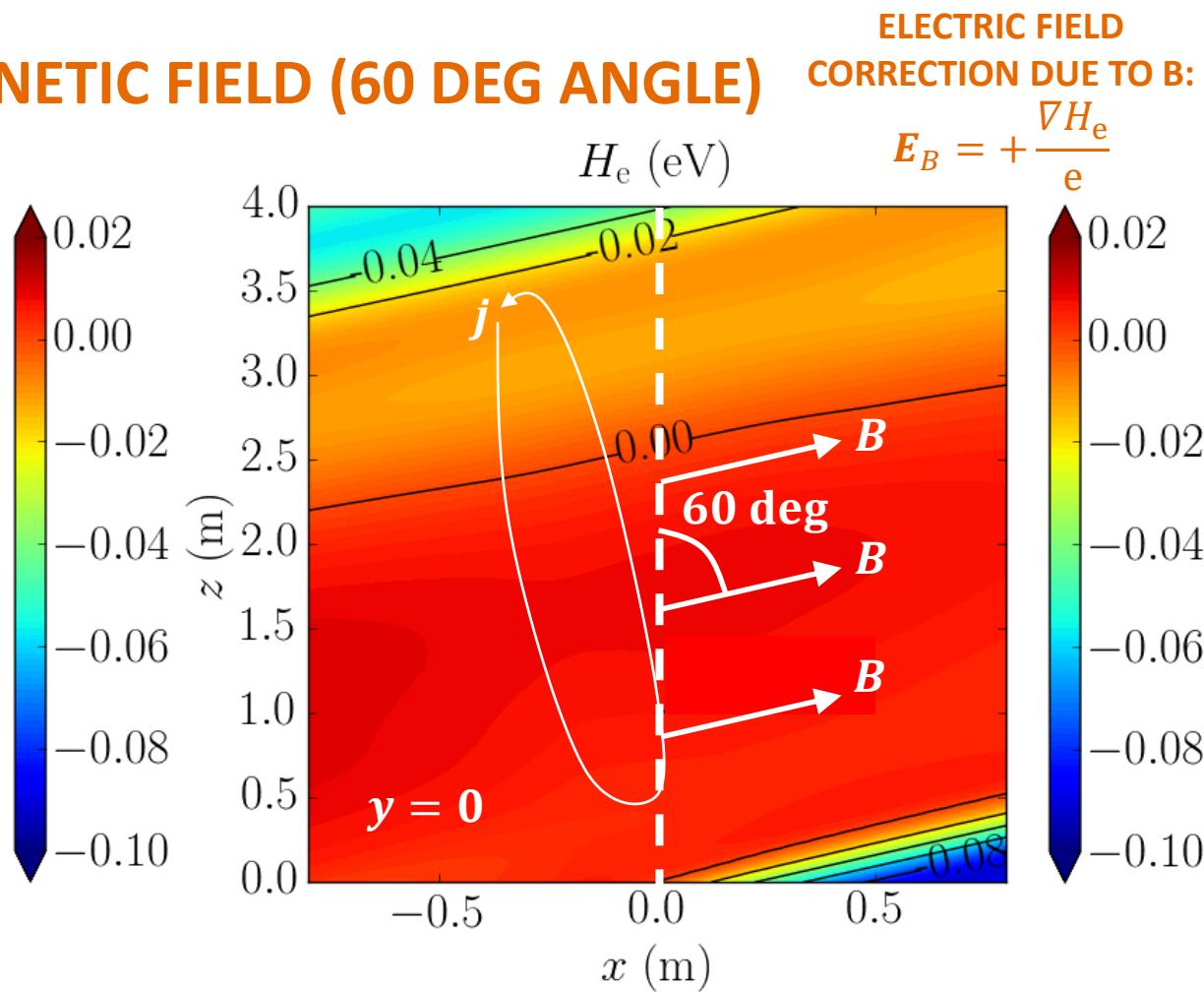
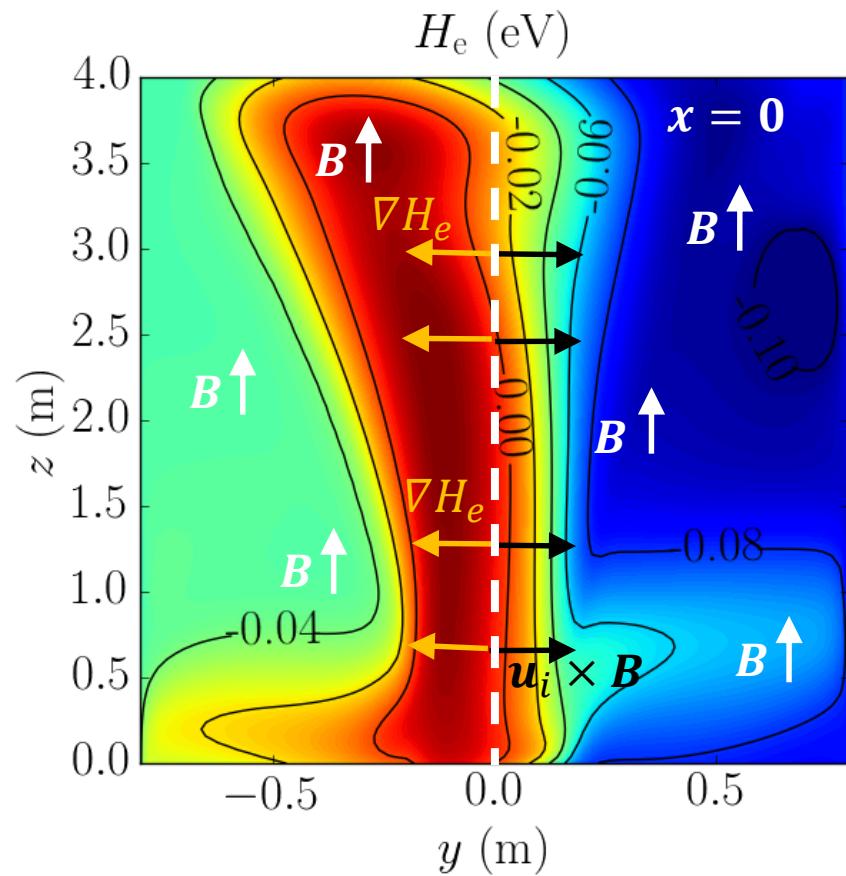
**PREDICTED
DISTORSION**

- Magnetic field intensity: **0.1 Gauss**



Plasma plume with uniform magnetic field (2)

OBLIQUE MAGNETIC FIELD (60 DEG ANGLE)

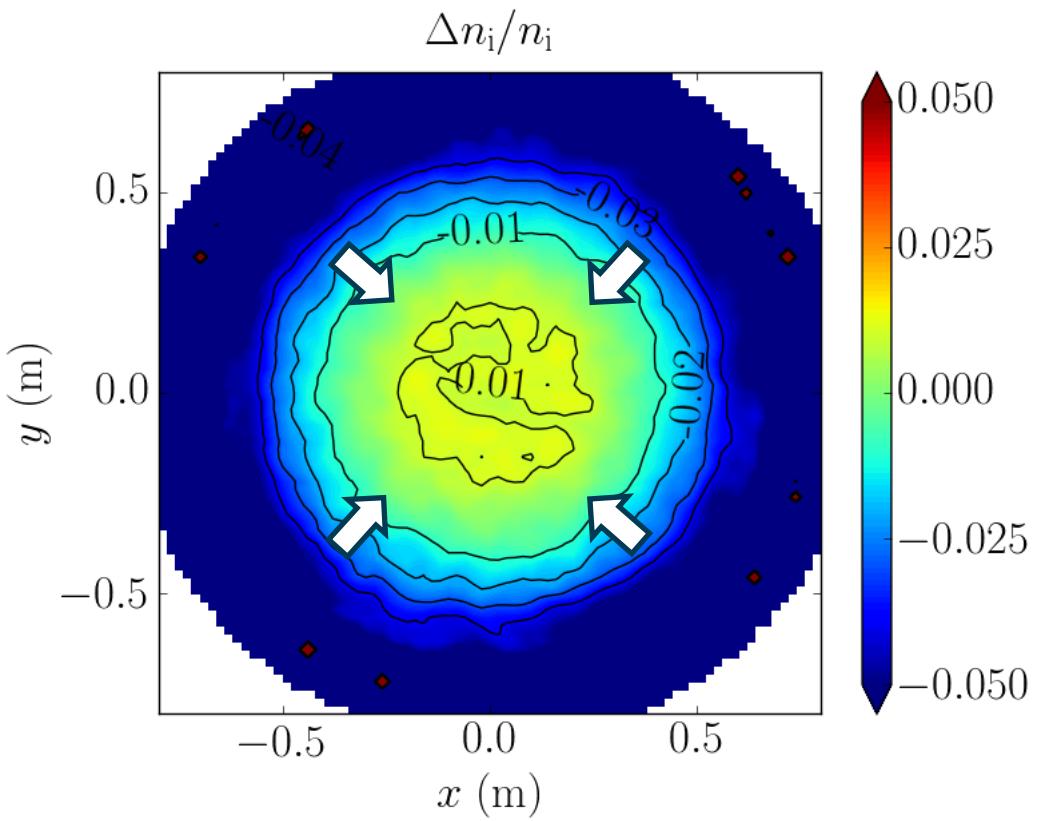


- $\frac{\partial H_e}{dy} \gg \frac{\partial H_e}{dx}$: compression is only in plane $y - z \rightarrow$ distortion of plume cross section
- Along plume centerline: $u_i \times B \approx -\nabla H_e/e \rightarrow$ no net deflection of plasma plume

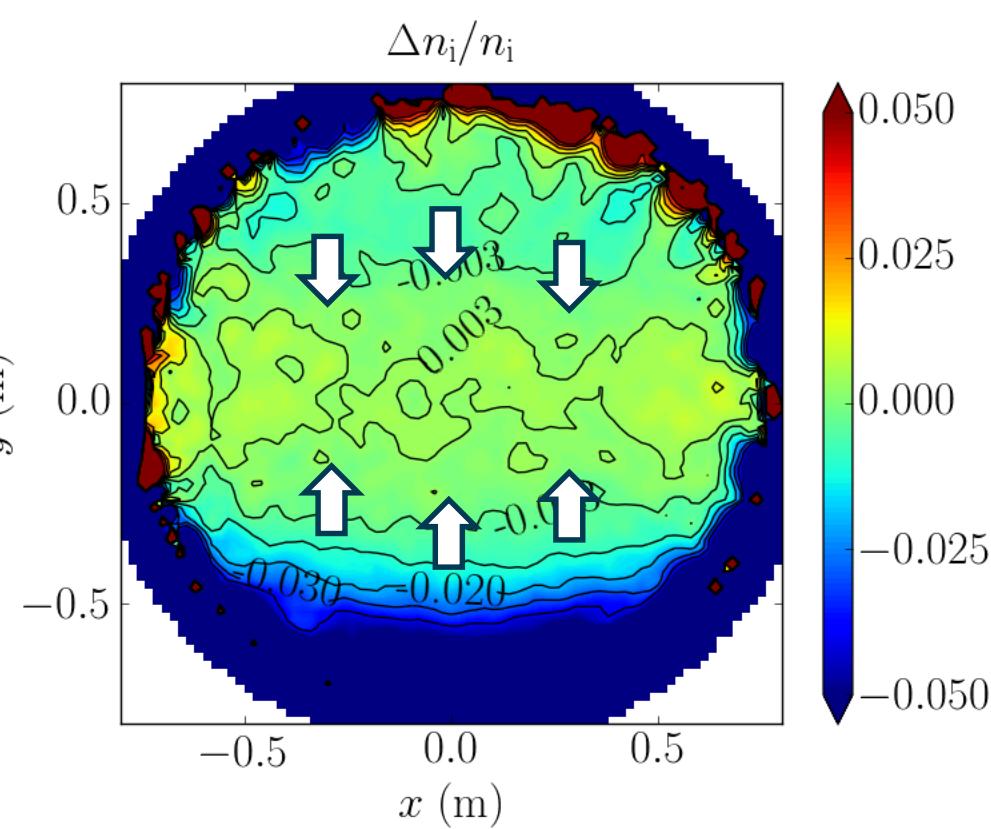
Plasma plume with uniform magnetic field (3)

- Ion density at a final cross section (4 m downstream) was compared with the unmagnetized case. The plume cross section
 - compresses in the axial field case, and
 - distorts in the oblique field case

AXIAL FIELD



OBLIQUE FIELD (60 DEG)

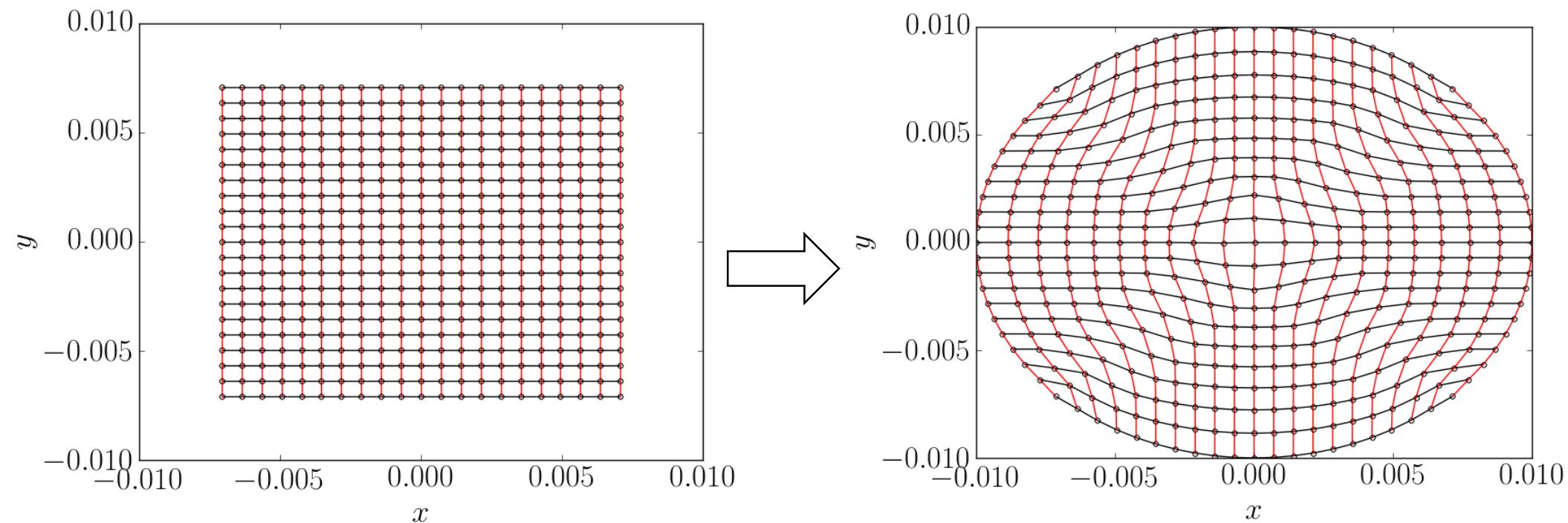


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Deformed meshes

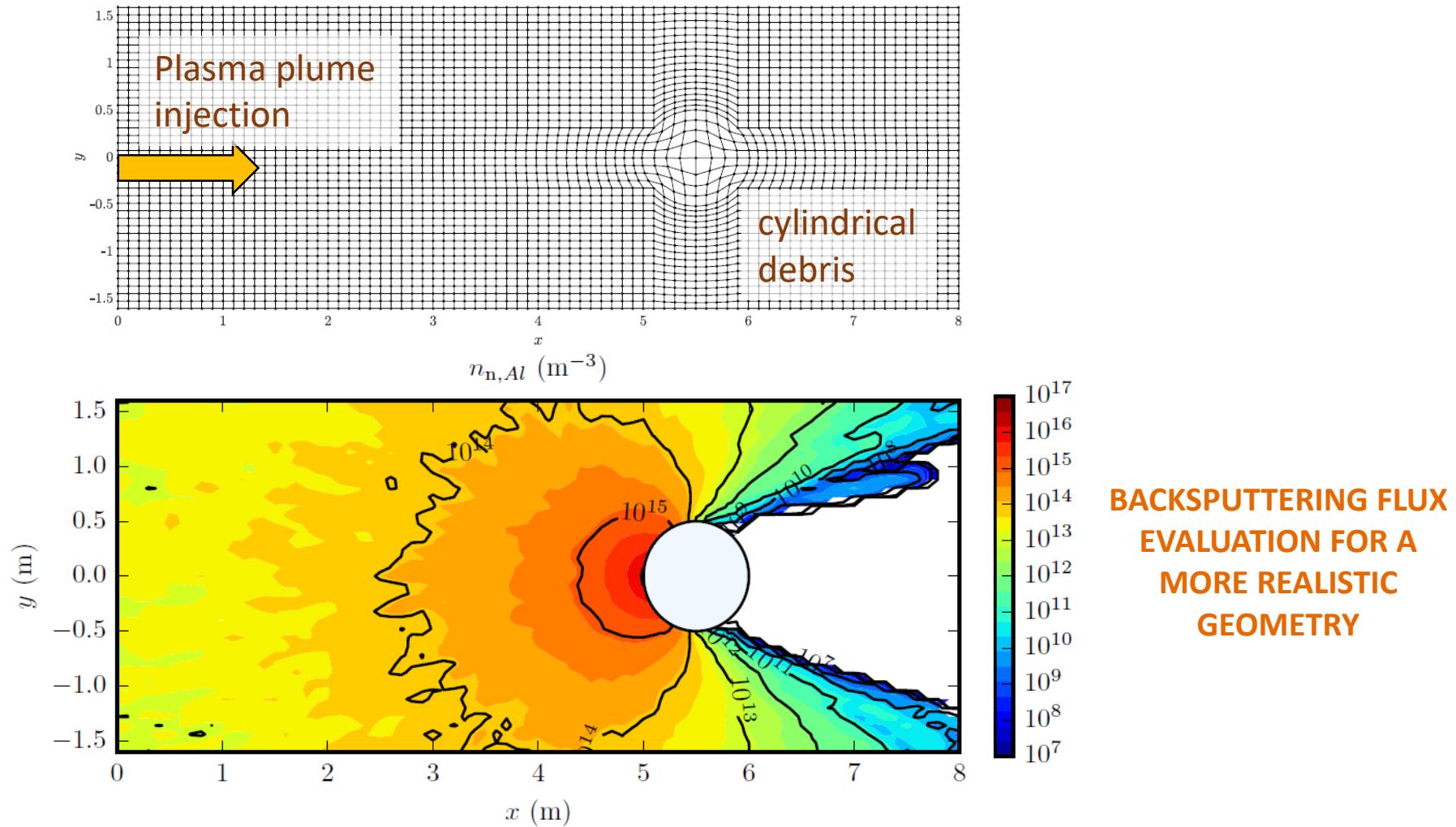
- A cylinder cross section is obtained from the deformation of a square (**J.M. Catalan's Master thesis**)
 - Coordinate parallel to the axis is unperturbed



- New types of simulations are enabled
 - Plasma thruster discharge simulations (inside a cylinder)
 - Ion-beam-shepherd interaction with a cylindrical debris
 - Ion beam optics with circular holes

Deformed mesh applications

- The interaction of a plasma plume with a cylindrical debris can be studied with such a deformed mesh
 - Improvement of simulation capabilities for the IBS scenario (**Cichocki AA 2018**)



BACKSPUTTERING FLUX
EVALUATION FOR A
MORE REALISTIC
GEOMETRY

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Conclusions and future studies

- EP2PLUS electron fluid model permits computing
 - Electric currents
 - Electric potential correction due to both collisions and magnetization
- Benchmark simulations for the electron model:
 - Neutralization of the current in an unmagnetized GIT plume
 - Effects of a uniform background magnetic field on plume expansion
- Newly developed deformed meshes extend the applications of the code
 - Cylindrical discharge chambers and target debris objects (IBS)
- Future work
 - Study of the near plume of HETs
 - Use Akiles opensource kinetic electron model
 - Full-PIC simulations of the plume expansion

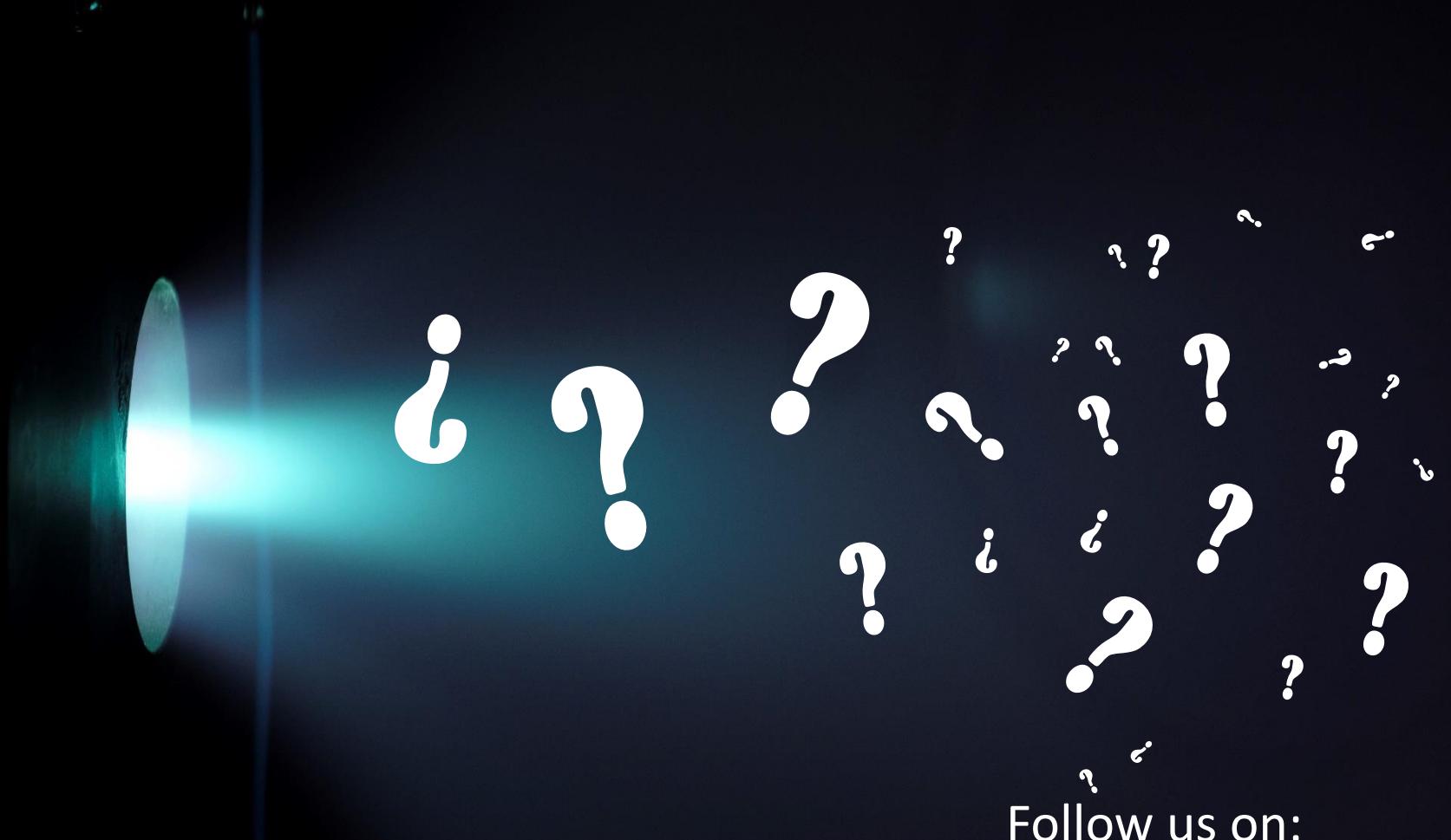


<https://github.com/ep2lab/akiles>

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Thank you! Questions?



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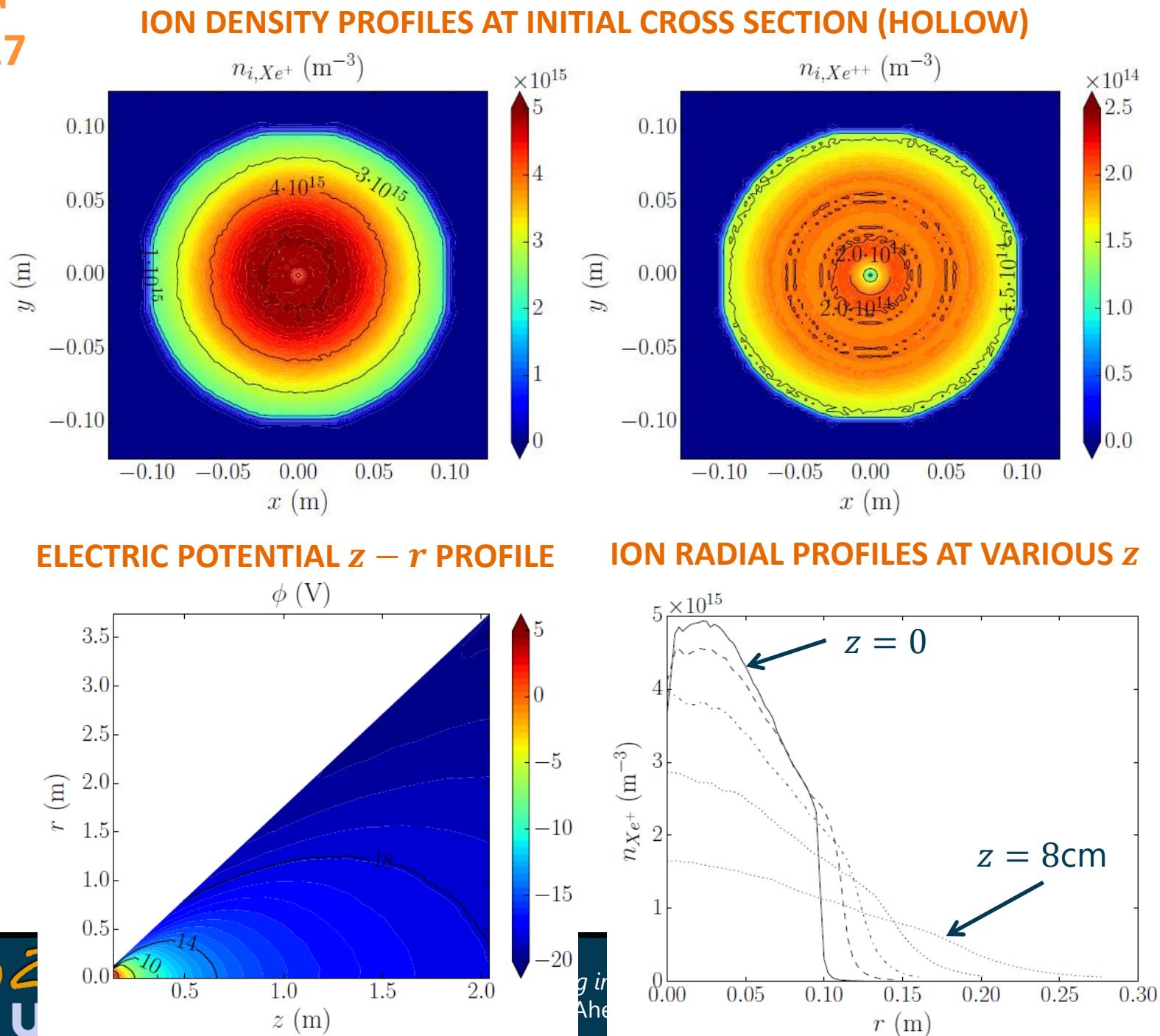
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Simulation of a HEMPT thruster plume

Kahnfeld
IEPC 2017



The electron fluid model (1)

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BAROTROPY GRADIENT

$$\nabla h_e = \nabla p_e/n_e \rightarrow h_e = h_e(n_e)$$

► Solved equations:

Electron momentum balance + electric current continuity:

$$\mathbf{0} = -\nabla p_e - en_e(\mathbf{E} + \mathbf{u}_e \times \mathbf{B}) - \sum_{s=1}^L \nu_{es} m_e n_e (\mathbf{u}_e - \mathbf{u}_s) \quad \cap \quad \nabla \cdot \mathbf{j} = 0$$

Poisson's equation for the electric potential:

$\nabla^2 \phi = -\frac{e}{\epsilon_0} (n_e^* - n_e)$ where $n_e^* = \sum_{s=1}^L Z_s n_s$ is the **quasineutral** electron density

The electron fluid model (2)

- In terms of a Bernoulli's function H_e , momentum balance becomes:

$$\nabla H_e = -e \nabla \phi + \nabla h_e \rightarrow \mathbf{j} = \bar{\bar{K}} \cdot \frac{\nabla H_e}{e} + \mathbf{j}_{PIC}$$

Conductivity tensor

function of heavy species properties and magnetic field

- Applying current continuity, we get an elliptic equation for H_e :

$$\nabla \cdot \mathbf{j} = 0 \rightarrow \bar{\bar{K}} \cdot \nabla \nabla H_e + \nabla H_e \cdot (\nabla \cdot \bar{\bar{K}}) = -e \nabla \cdot \mathbf{j}_{PIC}$$

- Once H_e is solved for, quasineutral electric potential is then obtained as:

$$\phi^* = \frac{h_e(n_e^*) - H_e}{e}$$

$\mathbf{H}_e \equiv \mathbf{0}$ for collisionless unmagnetized electrons!!

- Electric potential and electron densities are recomputed in non-neutral regions with a non-linear Poisson's solver:

- Quasineutral nodes are considered as Dirichlet points
- $\nabla^2 \phi = -\frac{e}{\epsilon_0} [n_e^* - n_e(\phi, H_e)]$ where $n_e(\phi, H_e) = h_e^{-1}(e\phi + H_e)$

Detailed electron fluid model (1)

- Considered equations:
 - Current continuity

$$\nabla \cdot \mathbf{j} = -\frac{\partial \rho_c}{\partial t} = 0$$

Stationary conditions

- Electron momentum balance equation

$$0 = -\nabla p_e - en_e(\mathbf{E} + \mathbf{u}_e \times \mathbf{B}) - \sum_{s=1}^L \nu_{es} m_e n_e (\mathbf{u}_e - \mathbf{u}_s)$$

- Introducing

- The gyrofrequency $\omega_{ce} = eB/m_e$
- The Hall parameter $\chi = \omega_{ce}/\nu_e$
- The scalar electron conductivity $\sigma_e = \frac{e^2 n_e}{m_e \nu_e}$
- The driving current density $\mathbf{j}_d = \mathbf{j}_i - \frac{en_e}{\nu_e} \sum_{s=1}^L \nu_{es} \mathbf{u}_s$
- The electron momentum balance can be written as:

$$\mathbf{j} + \chi(\mathbf{j} \times \mathbf{1}_b) = \sigma_e \left(\frac{\nabla p_e}{en_e} - \nabla \phi \right) + (\mathbf{j}_d + \chi \mathbf{j}_i \times \mathbf{1}_b)$$

unit vector along
the magnetic field

Detailed electron fluid model (2)

- Solving for \mathbf{j} and introducing the Bernoulli function H_e :

$$H_e = h_e - e\phi$$

- We finally obtain:

$$\mathbf{j} = \boxed{\bar{\mathbf{K}}} \cdot \frac{\nabla H_e}{e} + \boxed{\frac{\bar{\mathbf{K}}}{\sigma_e} \cdot (\mathbf{j}_d + \chi \mathbf{j}_i \times \mathbf{1}_b)}$$

we call this \mathbf{j}_{PIC}

Conductivity tensor

CONDUCTIVITY TENSOR

$$\bar{\mathbf{K}} = \sigma_e \begin{bmatrix} 1 & \chi b_z & -\chi b_y \\ -\chi b_z & 1 & \chi b_x \\ \chi b_y & -\chi b_x & 1 \end{bmatrix}^{-1}$$

- With the use of the current continuity $\nabla \cdot \mathbf{j} = 0$, we finally obtain:

ELLIPTIC PDE FOR H_e

$$\nabla \cdot \mathbf{j} = 0 \rightarrow \bar{\mathbf{K}} : \nabla \nabla H_e + \nabla H_e \cdot (\nabla \cdot \bar{\mathbf{K}}) = -e \nabla \cdot \mathbf{j}_{PIC}$$

- Once H_e is known, we obtain the electric potential as:

$$\phi = \frac{h_e - \boxed{H_e}}{e}$$

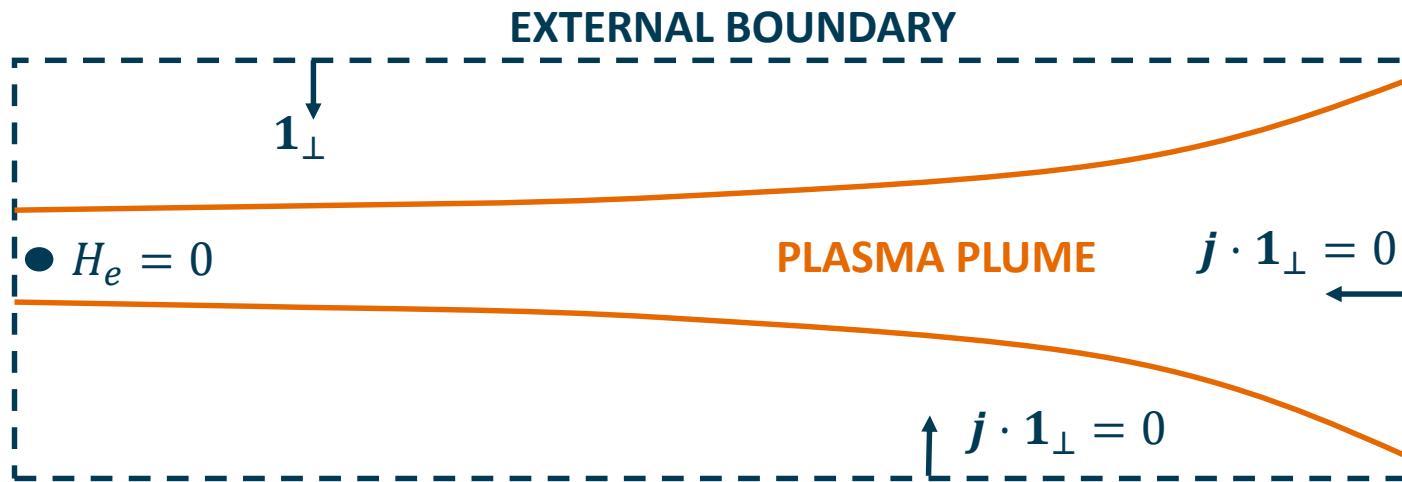
Correction with respect to the
collisionless, unmagnetized solution

- Parameters affecting the solution:

- Hall parameter χ
- Scalar conductivity σ_e

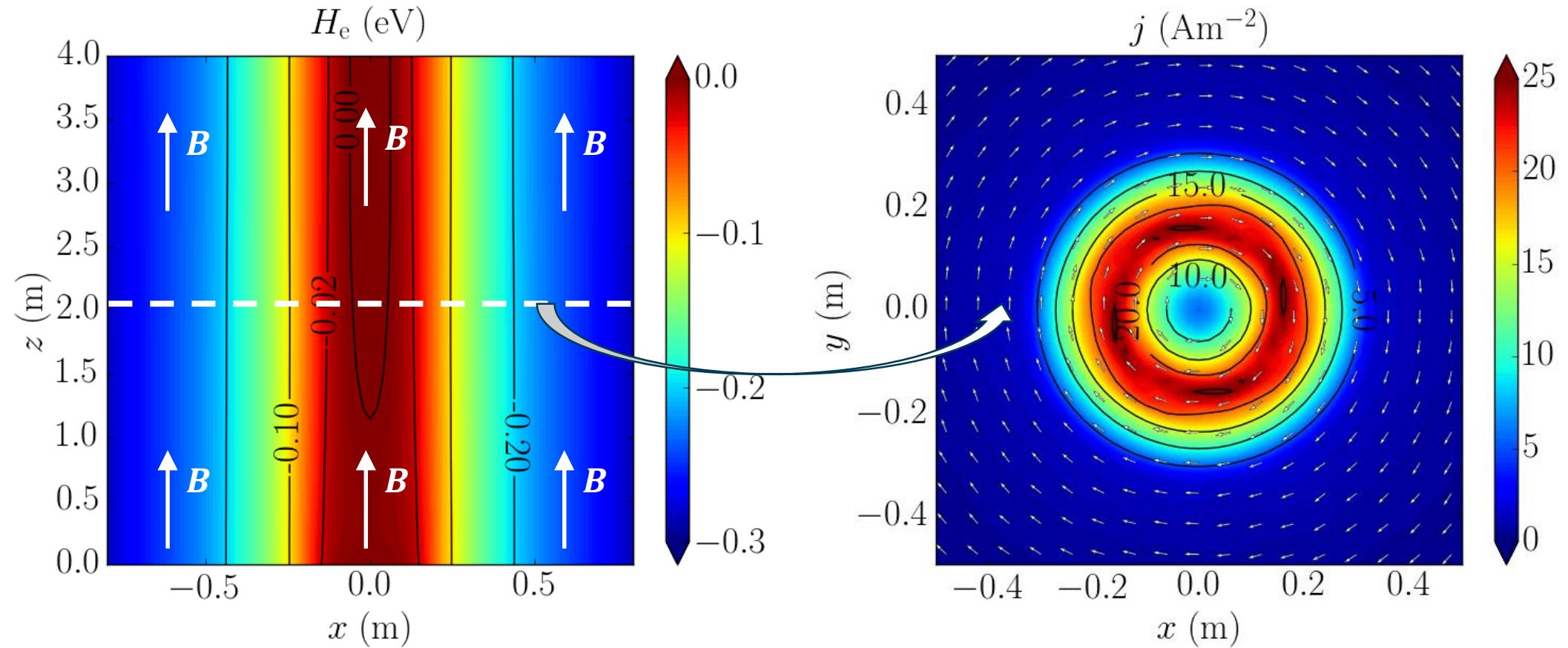
Detailed electron fluid model (3)

- $\nabla \cdot \mathbf{j} = 0$ urges (through Gauss theorem) that:
 - Total net current through simulation boundaries be zero
- We consider a stronger closure:
 - Zero normal current at all boundary nodes: $\mathbf{j} \cdot \mathbf{1}_\perp = 0$
- So, H_e is solved with the following boudary conditions:
 - $H_e = 0$ at a reference point where $\phi = 0, n_e = n_{e0}, T_e = T_{e0}$
 - $(\bar{\mathbf{K}} \cdot \nabla H_e) \cdot \mathbf{1}_\perp = -e \mathbf{j}_{PIC} \cdot \mathbf{1}_\perp$ at the external boundaries



Axial field case of the plasma plume expansion

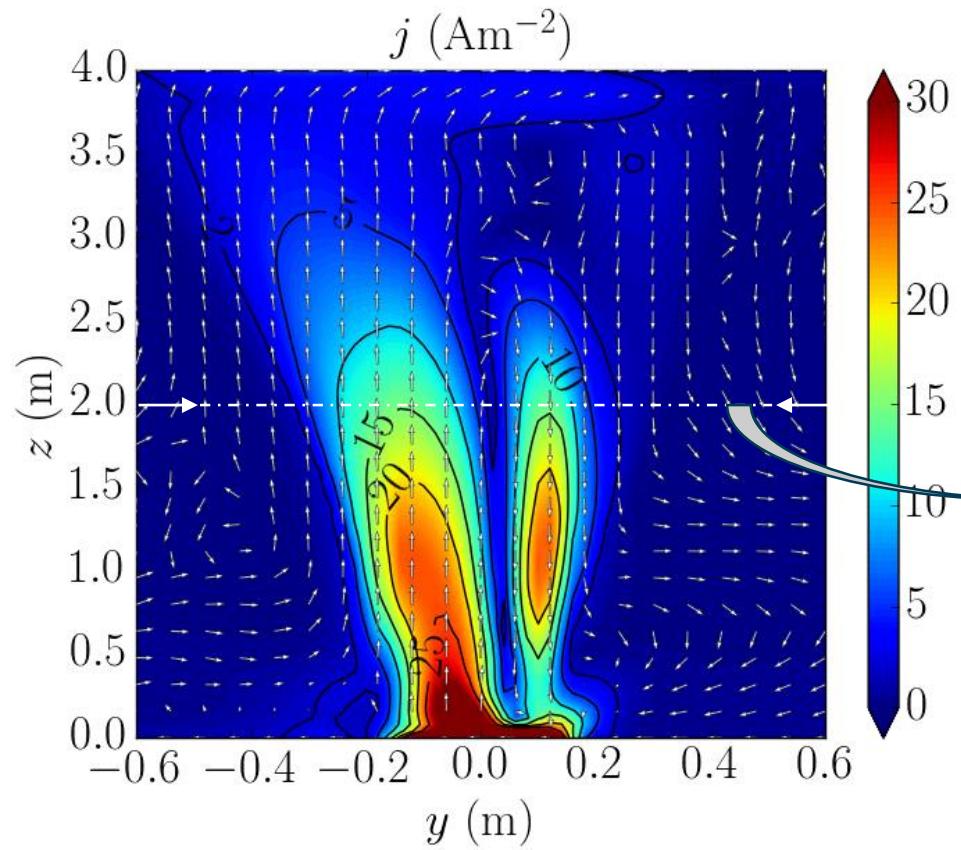
AXIAL MAGNETIC FIELD



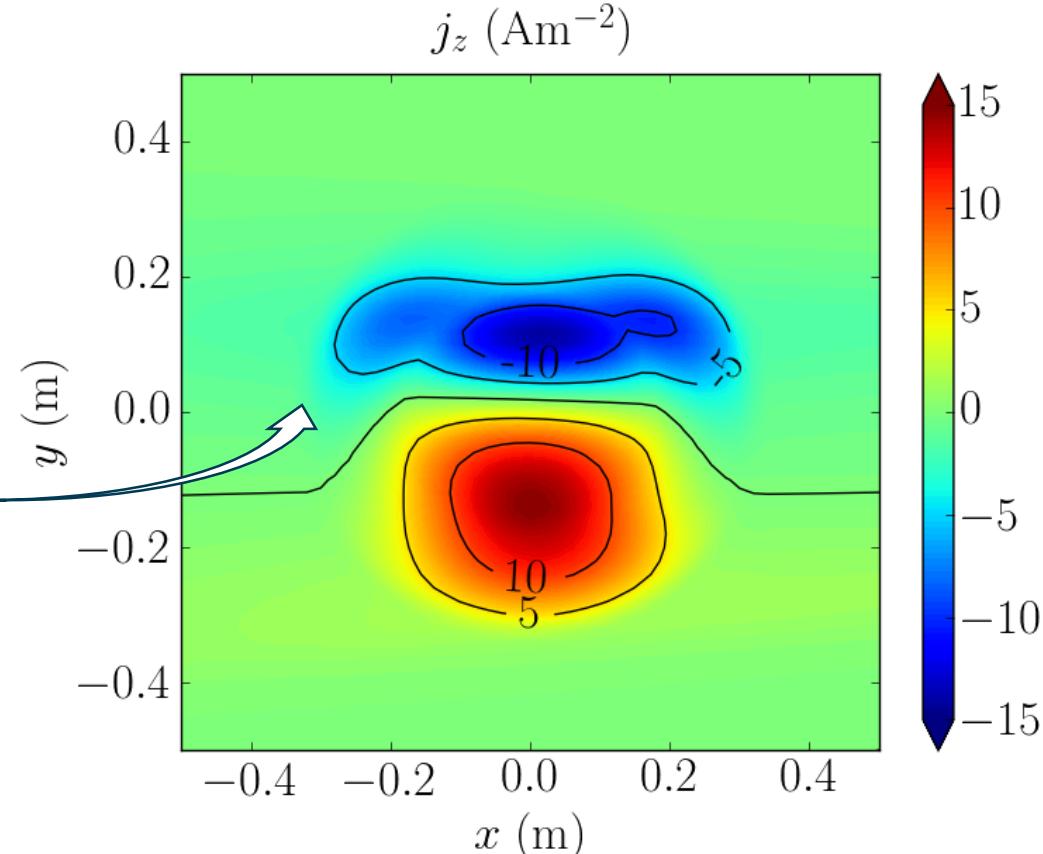
- The potential correction is axisymmetric → diamagnetic current loop
- Lower radial potential fall → **plume compression** (radial confinement)

Currents in the oblique magnetic field case (60 deg)

ELECTRIC CURRENT AT $x = 0$



Z-CURRENT DENSITY AT $z = 2 \text{ m}$



- Boundary conditions affect the currents nearby
 - Need to define a sufficiently large simulation box

Non-neutrality criteria in EP2PLUS

- Non-neutrality criterion for plasma nodes:

**RELATIVE CHARGE
DENSITY ESTIMATION**

$$\varepsilon_n = \left| \frac{n_e^* - n_e}{n_e^*} \right|^{1/2} = \left| \frac{\epsilon_0 \nabla^2 \phi^*}{e n_e^*} \right|^{1/2} < \varepsilon_{max}$$

- Non-neutrality criterion for material cell-faces:

**DEBYE LENGTH TO
CELL SIZE RATIO**

$$\varepsilon_f = \frac{1}{\Delta l} \sqrt{\frac{\epsilon_0 T_e^*}{e^2 n_e^*}} < \varepsilon_{max}$$

RELATIVE CHARGE DENSITY

