GPU aided 2D high-enthalpy flow solver with state-to-state kinetics

Gianpiero Colonna

CNR—Institute for Plasma Science and Technology, Bari

Francesco Bonelli, Davide Ninni, Giuseppe Pascazio Bari Polytechnic, Department of Mechanics, Mathematics and Management

Motivations



In front of bodies moving at hypersonic speed forms a shock wave where temperature can jump from hundreds to many thousands K, inducing vibrational excitation, dissociation and ionization.

It is well known that the system presents regions with strong non-equilibrium, affecting the macroscopic properties of the flow, such as the heat flux to the vehicle surface.

Approaches to treat air in non-equilibrium

Multi-Temperature

5 species17 reactions3 vibrational temperatures

State-to-State

~10² species ~10⁴ reactions 0 vibrational temperatures



CUDA for GP-GPU

The use of new technology can give relevant improvement in StS kinetics in CFD. Graphical processing units allow considerably speed-ups.



5 species State-to-State (StS) model

The State-to-State approach write a relaxation equation for each vibrational level so that it is possible to calculate the distribution of internal states when it departs from the Boltzmann one.

$$\begin{array}{c} \operatorname{Pure} \operatorname{N_2} \\ N_2(v) + N_2 \leftrightarrow N_2(v-1) + N_2 \\ N_2(v) + N \leftrightarrow N_2(v-\Delta v) + N \\ N_2(v) + N_2(w-1) \leftrightarrow N_2(v-1) + N_2(w) \\ N_2(v) + N_2 \leftrightarrow 2N + N_2 \\ N_2(v) + N \leftrightarrow 2N + N \end{array}$$

Pure
$$O_2$$

 $O_2(v) + O_2 \leftrightarrow O_2(v-1) + O_2$
 $O_2(v) + O \leftrightarrow O_2(v - \Delta v) + O$
 $O_2(v) + O_2(w-1) \leftrightarrow O_2(v-1) + O_2(w)$
 $O_2(v) + O_2 \leftrightarrow 2O + O_2$
 $O_2(v) + O \leftrightarrow 2O + O$

$$\begin{array}{c} \operatorname{Mixed} \operatorname{N_2} \\ N_2(v) + O_2 \leftrightarrow N_2(v-1) + O_2 \\ N_2(v_{\max}) + O_2 \leftrightarrow 2N + O_2 \\ N_2(v) + O \leftrightarrow N_2(v-1) + O \\ N_2(v_{\max}) + O_2 \leftrightarrow 2N + O_2 \end{array}$$

 $\begin{array}{c} \operatorname{Mixed} \operatorname{O_2} \\ O_2(v) + N_2 \leftrightarrow O_2(v-1) + N_2 \\ O_2(v_{\max}) + N_2 \leftrightarrow 2O + N_2 \\ O_2(v) + N \leftrightarrow O_2(v-1) + N \\ O_2(v_{\max}) + N \leftrightarrow 2O + N \end{array}$

$$O_2(v) + N_2(w-1) \rightleftharpoons O_2(v-2) + N_2(w)$$

Zeldovich exchange reactions

$$\begin{array}{l} O_2(v) + N \Longleftrightarrow NO + O \\ N_2(v) + O \Longleftrightarrow NO + N \end{array}$$



Flow past a sphere: Nonaka^{*} test case

*S. Nonaka et al. ,JTHT 14 (2), 2000



Computational domain,

with an example of 4 x4 MPI partitioning, along with boundary conditions (left). 152x392 computational grid shown every 2 grid points (right).



Nonaka test case



G. Colonna, F. Bonelli, G. Pascazio, Impact of fundamental molecular kinetics on macroscopic properties of high-enthalpy flows: The case of hypersonic atmospheric entry, Physical Review Fluids, 4, 033404 (2019)



Nonaka test case: comparison along stagnation line



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Nonaka test case: highest vibrational level contour plot





Finite rate catalysis model

I. Armenise et al. JTHT 20, 465–476 (2006) M. Barbato et al. JTHT 14, 412–420 (2000)

atom chemisorption (ch)

$$\begin{array}{l} N+^* \to N^* \\ O+^* \to O^* \end{array} \qquad \gamma_{AB} = \frac{\text{Flux of atoms recombining at the surface}}{\text{Flux of atoms impinging on the surface}} \end{array}$$

molecule chemisorption (chdm)

$$\begin{array}{l} N_2 + 2^* \to N^* + N^* \\ O_2 + 2^* \to O^* + O^* \\ NO + 2^* \to N^* + O^* \end{array} \quad \gamma_{NN} = \frac{2(-[N_2][S]^2 k_{chdm}^{N_2} + [N][N^*] k_{ER}^{NN} + [N^*]^2 k_{LH}^{NN})}{Z_N} \\ \end{array}$$

Eley-Rideal (ER)

$$\begin{array}{l} \mathrm{N} + \mathrm{N}^{*} \to \mathrm{N}_{2} +^{*} \\ \mathrm{O} + \mathrm{O}^{*} \to \mathrm{O}_{2} +^{*} \\ \mathrm{N} + \mathrm{O}^{*} \to \mathrm{NO} +^{*} \\ \mathrm{O} + \mathrm{N}^{*} \to \mathrm{NO} +^{*} \end{array} \quad \gamma_{OO} = \frac{2(-[O_{2}][S]^{2}k_{chdm}^{O_{2}} + [O][O^{*}]k_{ER}^{OO} + [O^{*}]^{2}k_{LH}^{OO})}{Z_{O}} \\ \gamma_{NO} = \frac{(-[NO][S]^{2}k_{chdm}^{NO} + [N][O^{*}]k_{ER}^{NO} + [O][N^{*}]k_{ER}^{ON} + [N^{*}][O^{*}]k_{LH}^{NO})}{Z_{N}} \end{array}$$

Langmuir–Hinshelwood (LH)

thermal desorption (td)

$$O^* \to O +^* \\ N^* \to N +^* \qquad \qquad Z_A = [A] \sqrt{kT/(2\pi m_A)}$$



CIRA: SCIROCCO Plasma Wind Tunnel test case

F. Bonelli et al., Effect of finite-rate catalysis on wall heat flux prediction in hypersonic flow, Phys. Rev. Fluids 6, 033201

DEMISE

eesa



SCIROCCO Plasma Wind Tunnel test cases

StS-1: recombing molecules have the same distribution of incoming ones

StS-2: recombing molecules have uniform distributions

StS-3: recombing molecules populate the highest vibrational level

Exp.	Park	StS-1	StS-2	StS-3	Park FC
	(err. %)				

q_{probe}	1543	1708	1873	1816	1774	2160
$[kW/m^2]$		(10.8%)	(21.38%)	(17.69%)	(14.97%)	(39.99%)

 $[kW/m^2]$ Park StS-1 StS-2 StS-3 Park FC 1146.6roto-translational 10821213.3 1516.01763.8(% contribution)(63.34%)(83.48%)(99.41%) (53.07%)(64.78%)diffusive 1000.5606.10660.17 300.010.38(35.25%) (16.52%) (0.59%)(46.31%) (% contribution)(35.48%)vibrational 20.1813.3(% contribution)(1.18%)(0.62%)

Total heat flux at the stagnation point

Decomposition of the total heat flux at the stagnation point



CIRA: SCIROCCO Plasma Wind Tunnel test case stagnation line profiles

F. Bonelli et al., Effect of finite-rate catalysis on wall heat flux prediction in hypersonic flow, Phys. Rev. Fluids 6, 033201



DESIGN FOR

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CIRA: SCIROCCO Plasma Wind Tunnel test case – vibrational distributions





DOUBLE WEDGE TEST CASES

D. Ninni et. Acta astronautica 191 (2022) 178



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DOUBLE WEDGE TEST CASES: StS vs, Park

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Conclusions

- The StS model has been implemented in 2D fluid dynamic code accelerated by GPU (speedup ~ 100)
- The code has been applied to blunt body including also state-resolved surface processes and kinetic equation for active surface site occupation.
- Application to unsteady flows as double wedge.
- Comparison with Park multi-temperature models shows differences in all the test cases
- 3D version of the code is under construction,



