

NEOs in the thermal infrared (NEATM, TPM) & Pre-impact detectability of a Chelyabinsk-type object

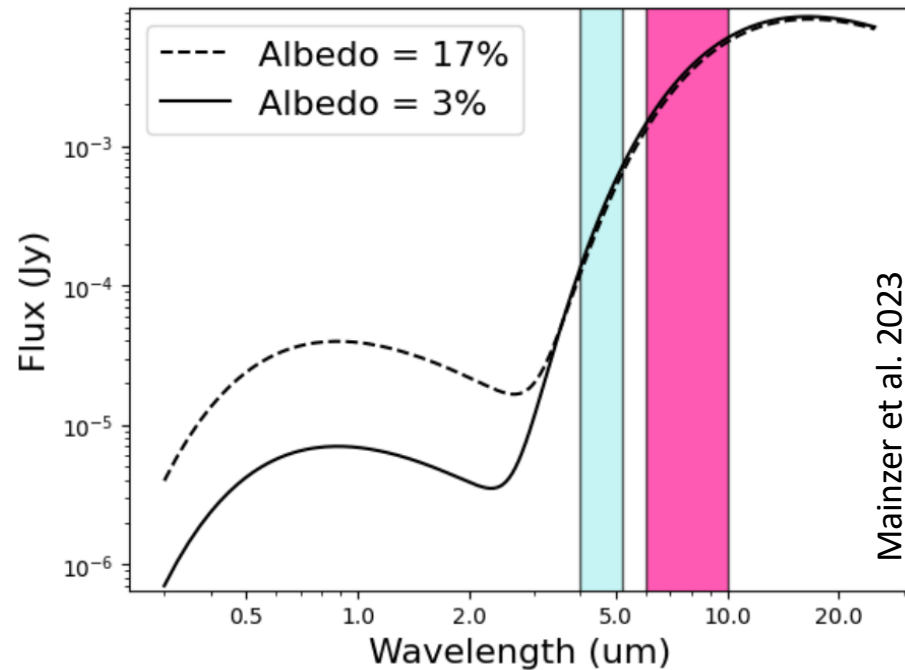
Thomas Müller, MPI für Extraterrestrische Physik, Garching, Germany (tmueller@mpe.mpg.de)



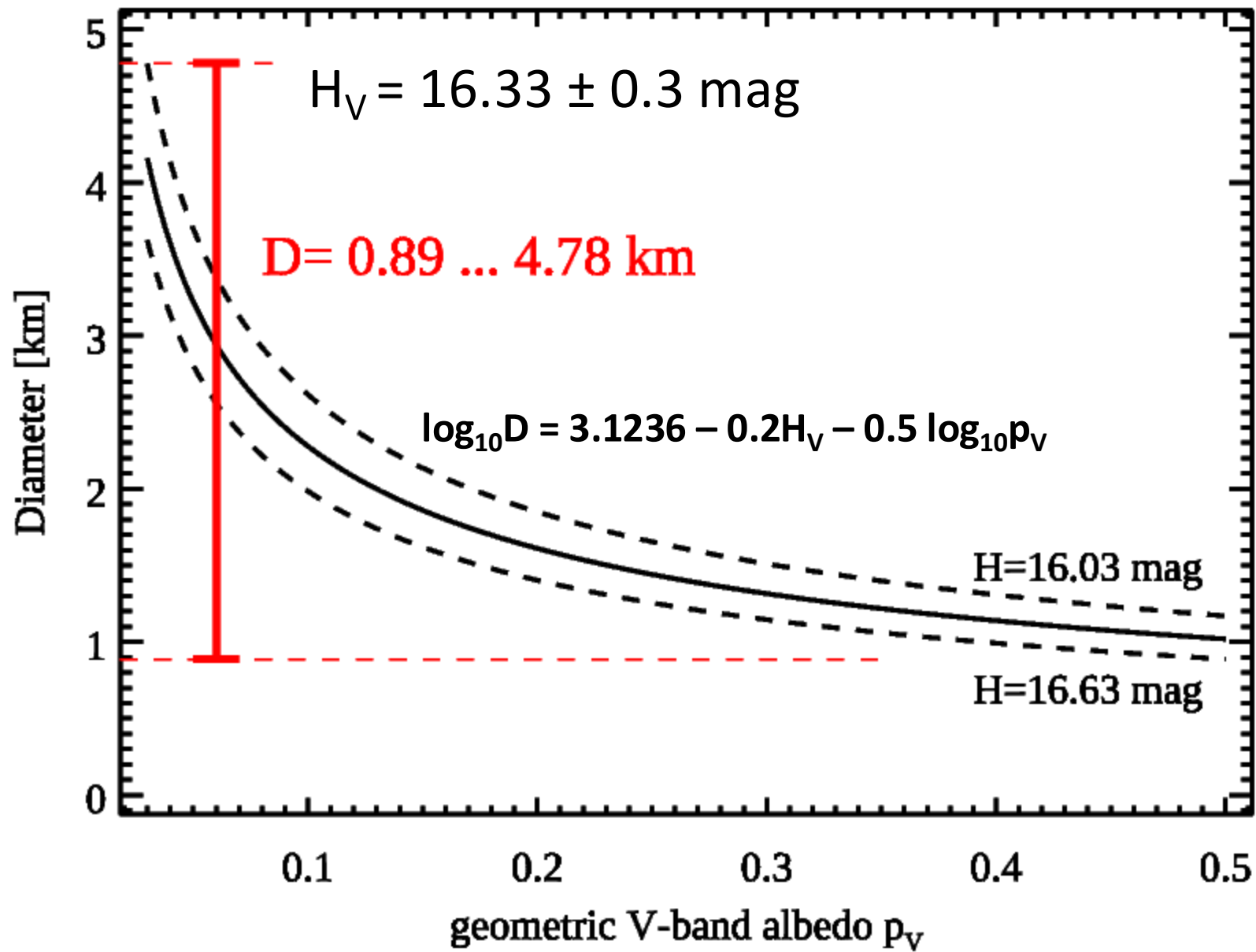
<https://www.youtube.com/watch?v=iCawTYPtehK>

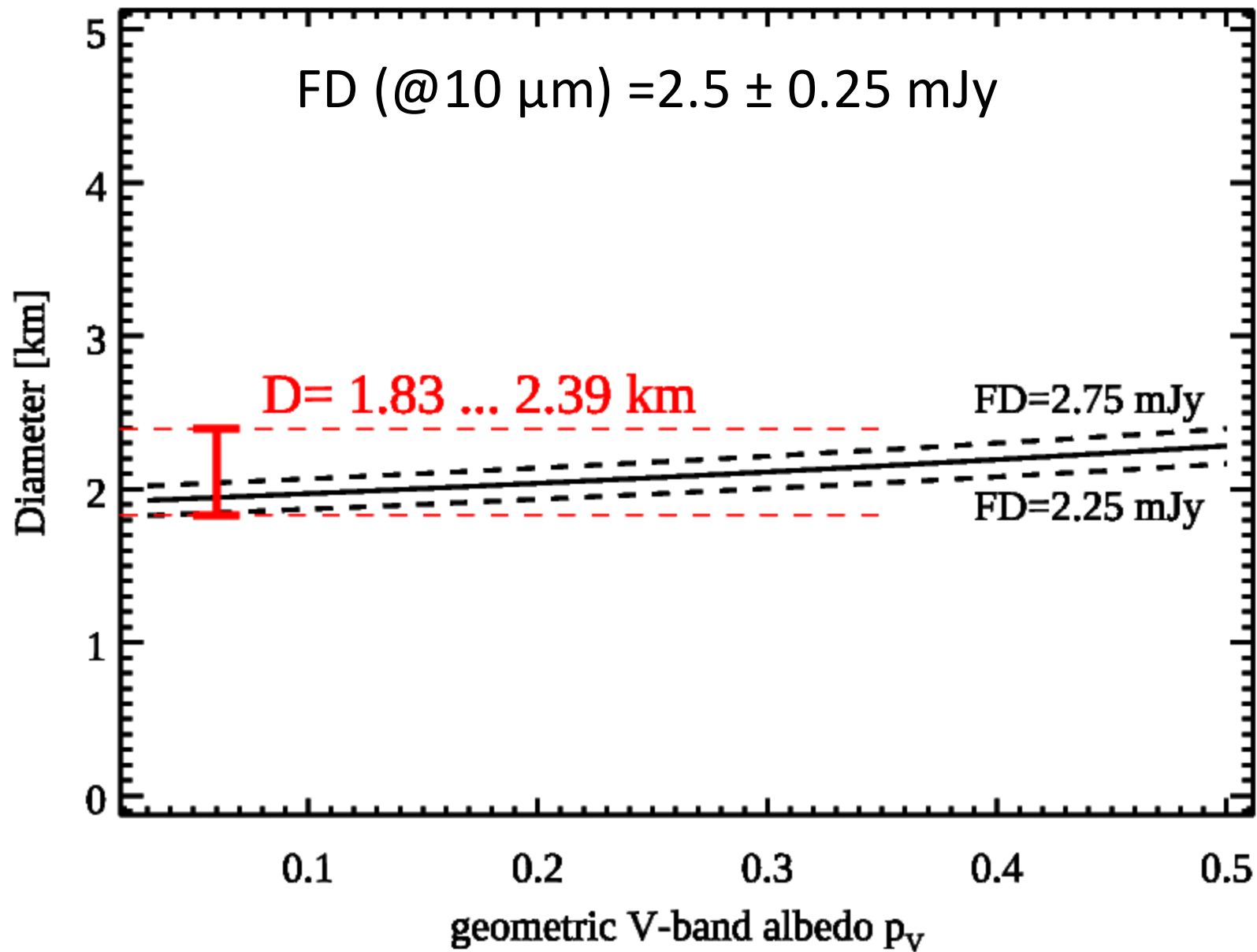
Credit: Aleksandr Ivanov

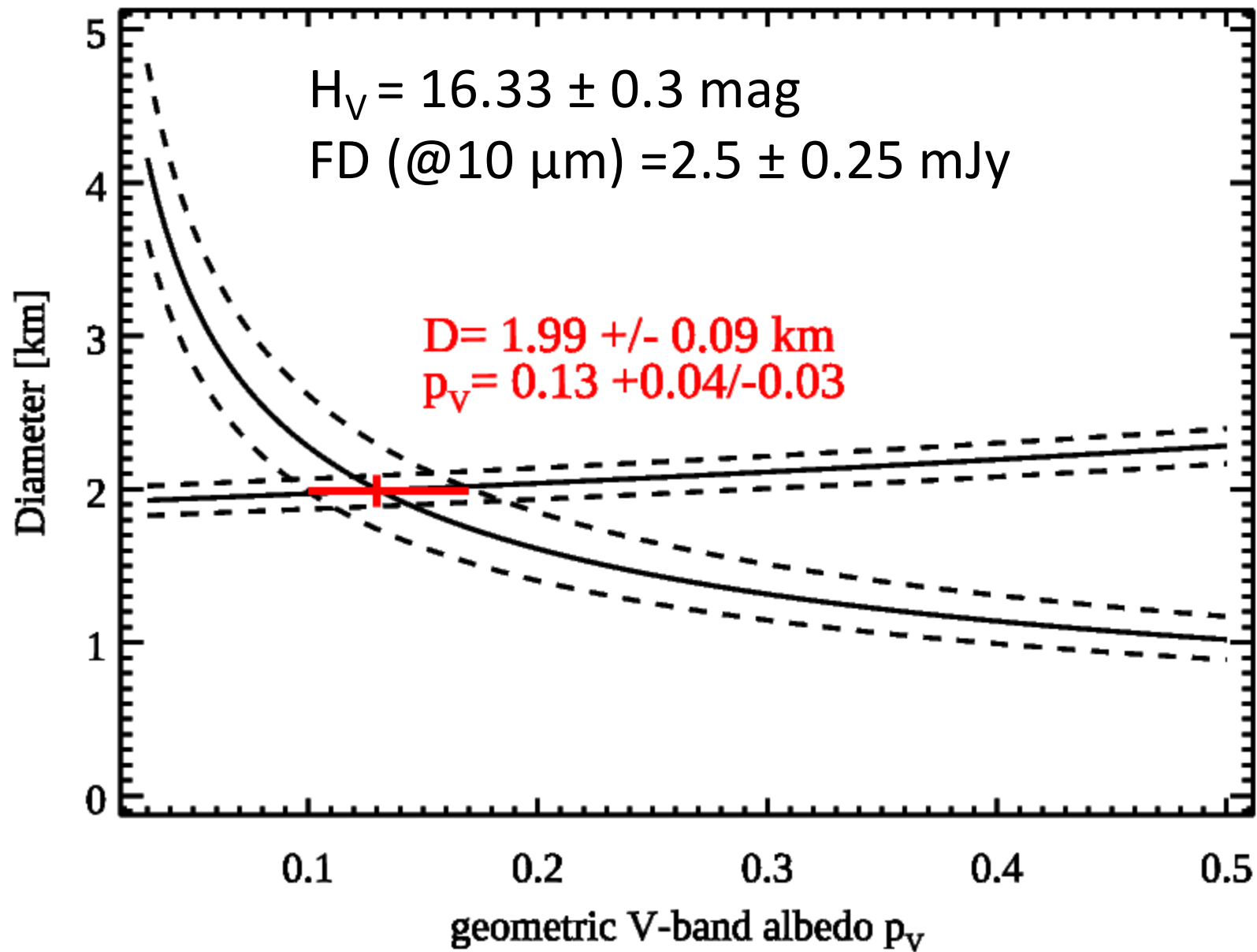
Why Infrared?



- About 90% of the incoming solar radiation is emitted in the IR
 - The risk of source confusion (stars, galaxies) is in the mid-IR about 2 orders of magnitude lower than at visible λ
 - The **high IR-VIS flux ratio** and the **reduced confusion risk** are especially advantageous for observing NEOs which are often viewed at large phase angles and close to the Sun.
 - Visible observations of irregularly-shaped NEOs at high phase angle also suffer from rotational variations of the small, illuminated surface areas
-
- In contrast, thermal IR observations present a different scenario: small, fast-rotating NEOs exhibit **nearly isothermal surfaces** with temperatures ranging from 300 to 400 K. Consequently, the likelihood of early detection is enhanced at IR wavelengths, especially at large phase angles.
 - Additionally, IR measurements provide **valuable constraints on an object's size, albedo, and thermal characteristics** indicative for the surface material strengths







Size determination for MBAs (larger objects, limited phase angle range):

- H-mag only: factor 4-5 uncertainty in size
- IR detection, no H-mag: $\pm 20\text{-}30\%$ via NEATM
- IR detection, with H-mag: $\pm 10\text{-}20\%$ via NEATM
- Multiple/multi-band detections of objects with known orbit and H-mag: $\pm 10\%$ via TPM techniques (if spin-shape solution available: $\pm 5\%$)

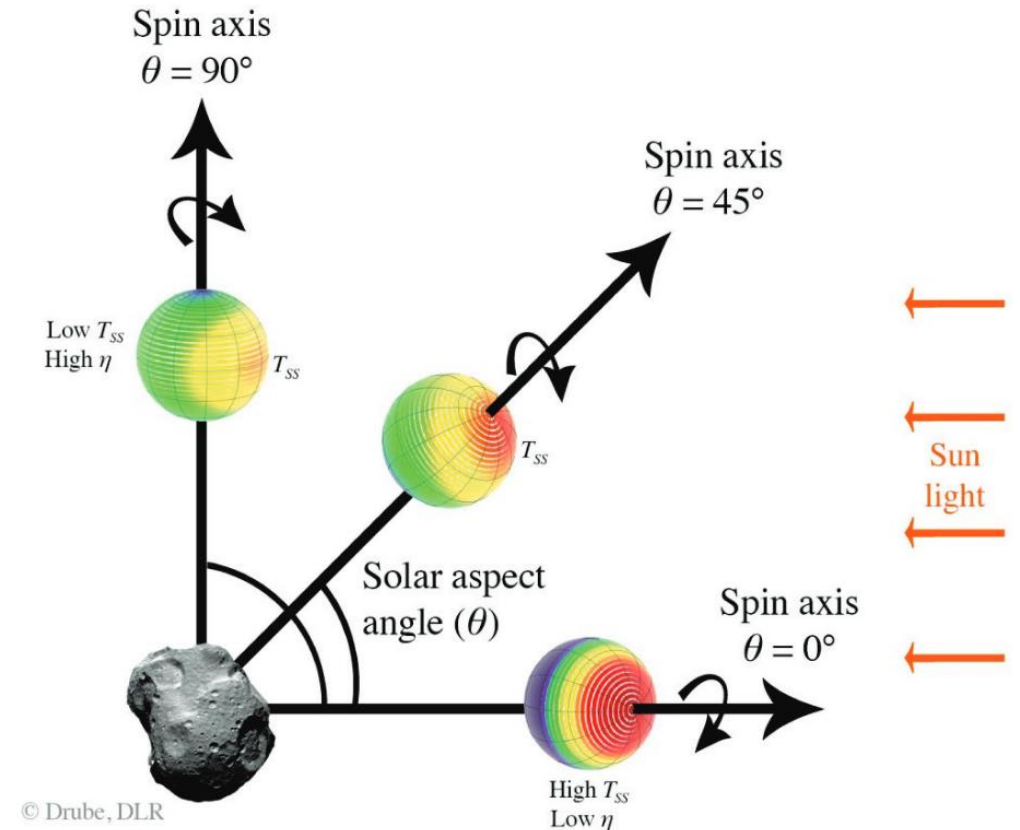
Size determination for NEAs (small, fast(er) rotation, shape, large phase angles):

→ **NEA radiometric size determination easily possible (WISE, NEO Surveyor, NEOMIR, ...), but**

- Lower quality due to large changes in cross section (shape, aspect angle)
- No thermal model validation for objects at large ($>90^\circ$) phase angles
- Unknown surface roughness and thermal inertia effects (bare rock, porous rocks, fine/coarse grain regolith on low-g surface)
- Unknown temperature distribution for small, fast-rotating asteroids
- Thermal properties change for highly eccentric orbits

NEATM (Harris 1998) application: Ryugu

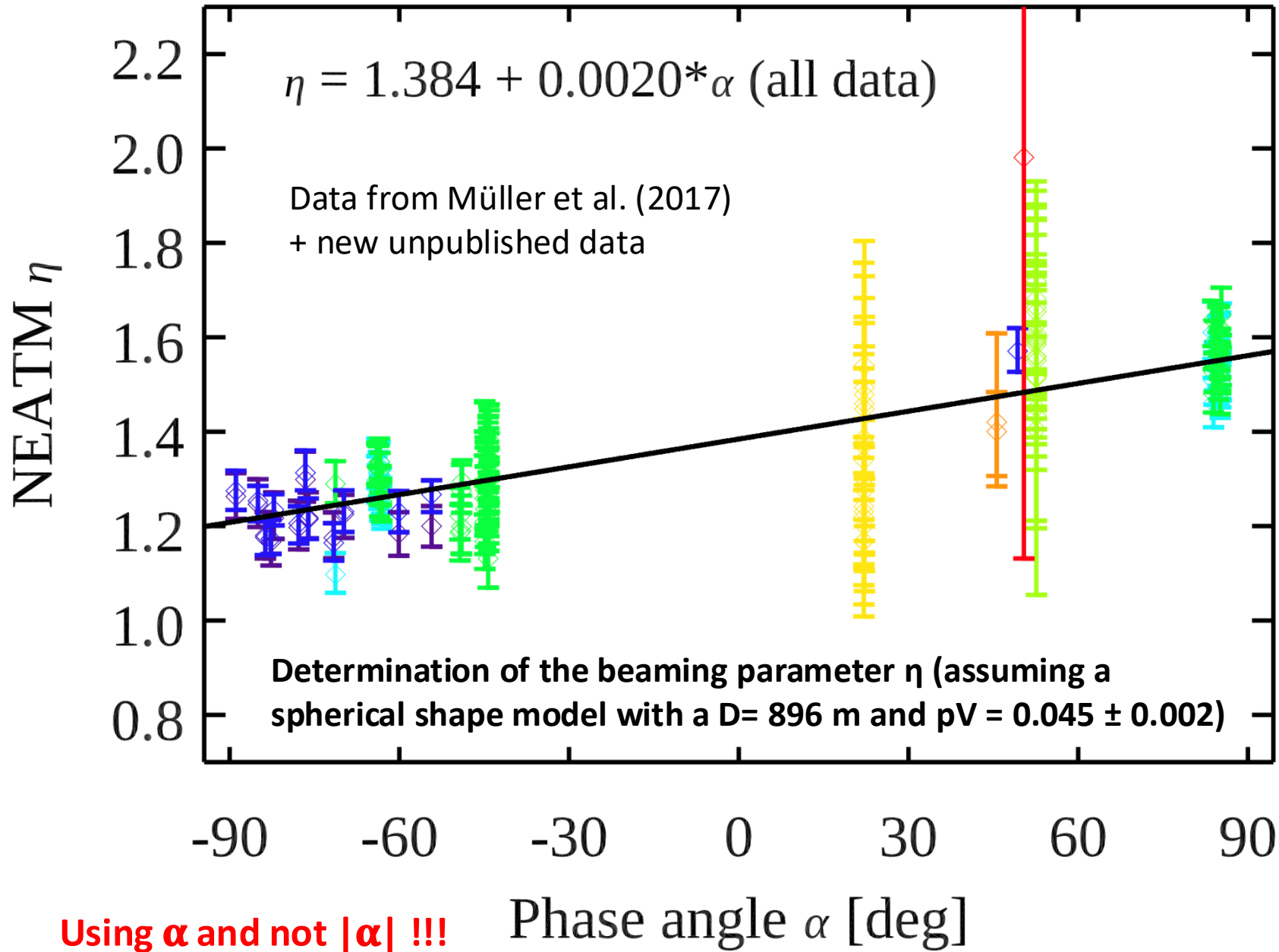
- Assuming a spherical shape model and only day-side emission is considered (using the absolute phase angle $|\alpha|$)
- $T = [S (1 - A) \cos(\Theta) / (\eta \epsilon \sigma)]^{1/4}$ where A is the Bond albedo, ϵ is the bolometric emissivity (assumed to be 0.9), σ is the Stefan-Boltzmann constant, S is the solar insolation at the distance of the object, Θ is the angle away from the sub-solar point, and η is the beaming parameter.
- The reflected sunlight (at zero phase angle) is determined by the visible geometric albedo p_V . It is connected to the Bond albedo A through the phase integral q , as $A = q p_V$
- q can be calculated from the slope parameter G via $q = 0.290 + 0.684 G$ (Bowell et al. 1989)



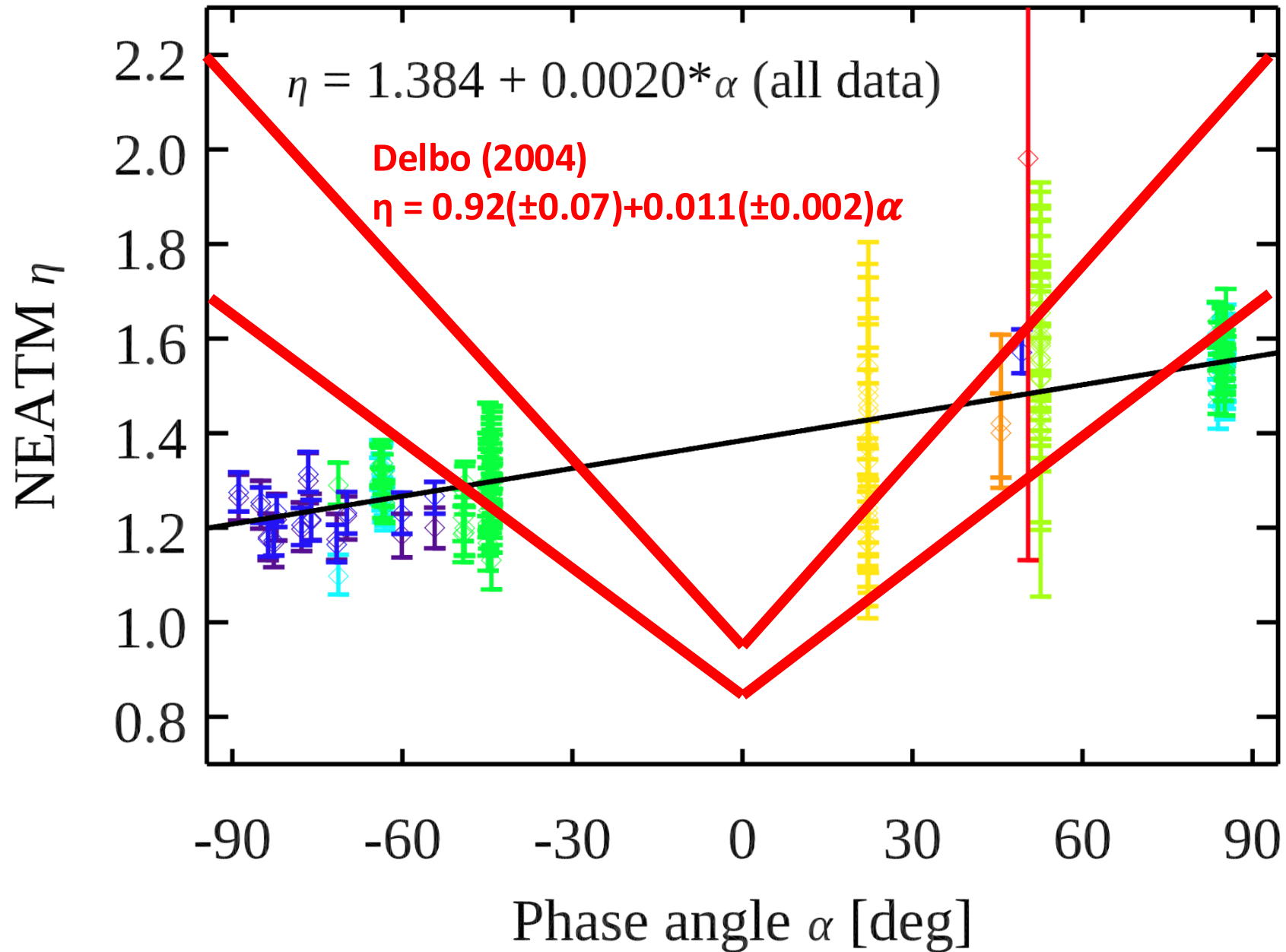
Application to all available remote observations from Spitzer-IRAC

(ch1,ch2)/-IRS, WISE-W1/-W2, Subaru-COMICS, AKARI-IRC, Herschel-PACS

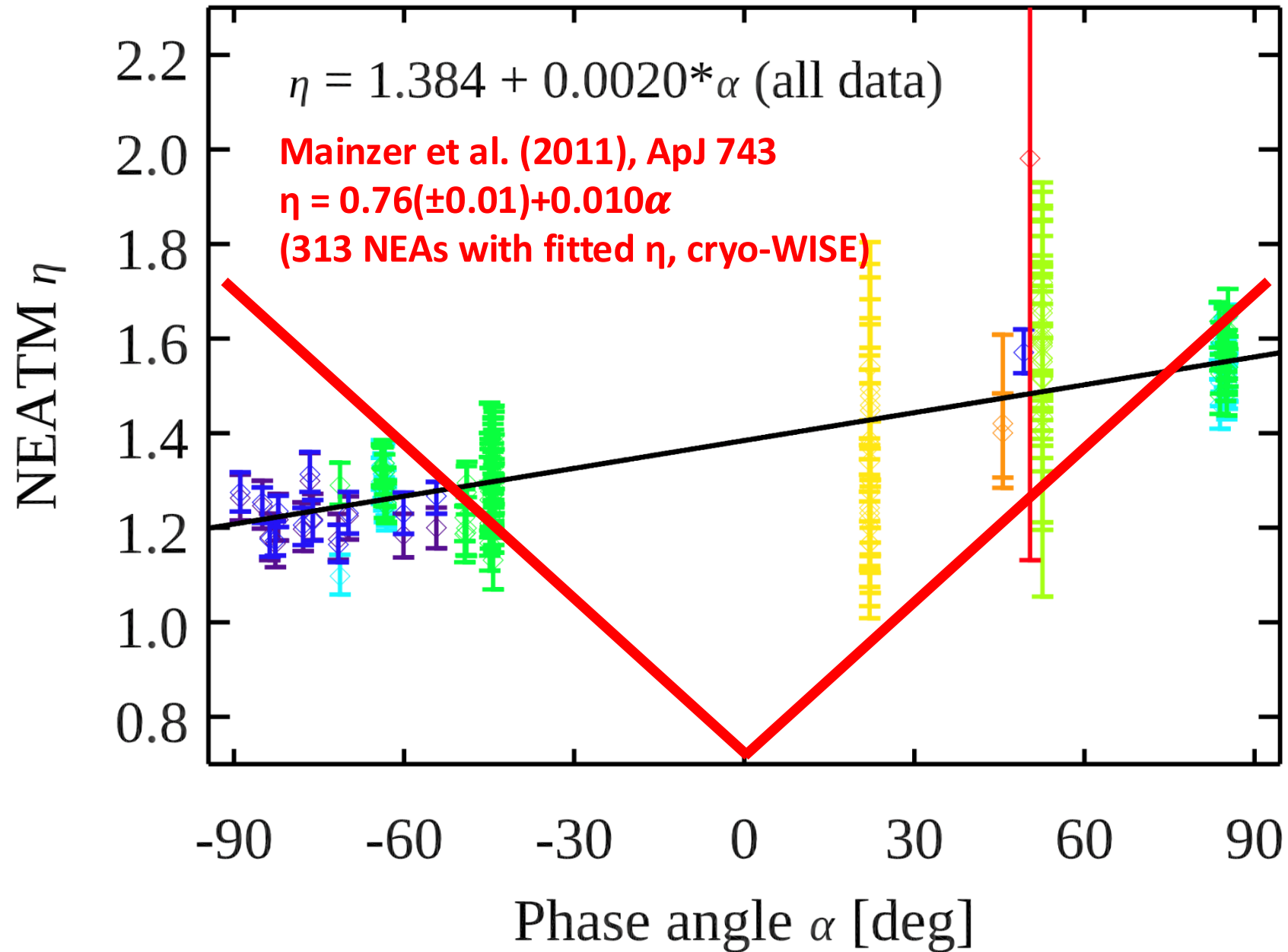
NEATM Beaming phase function



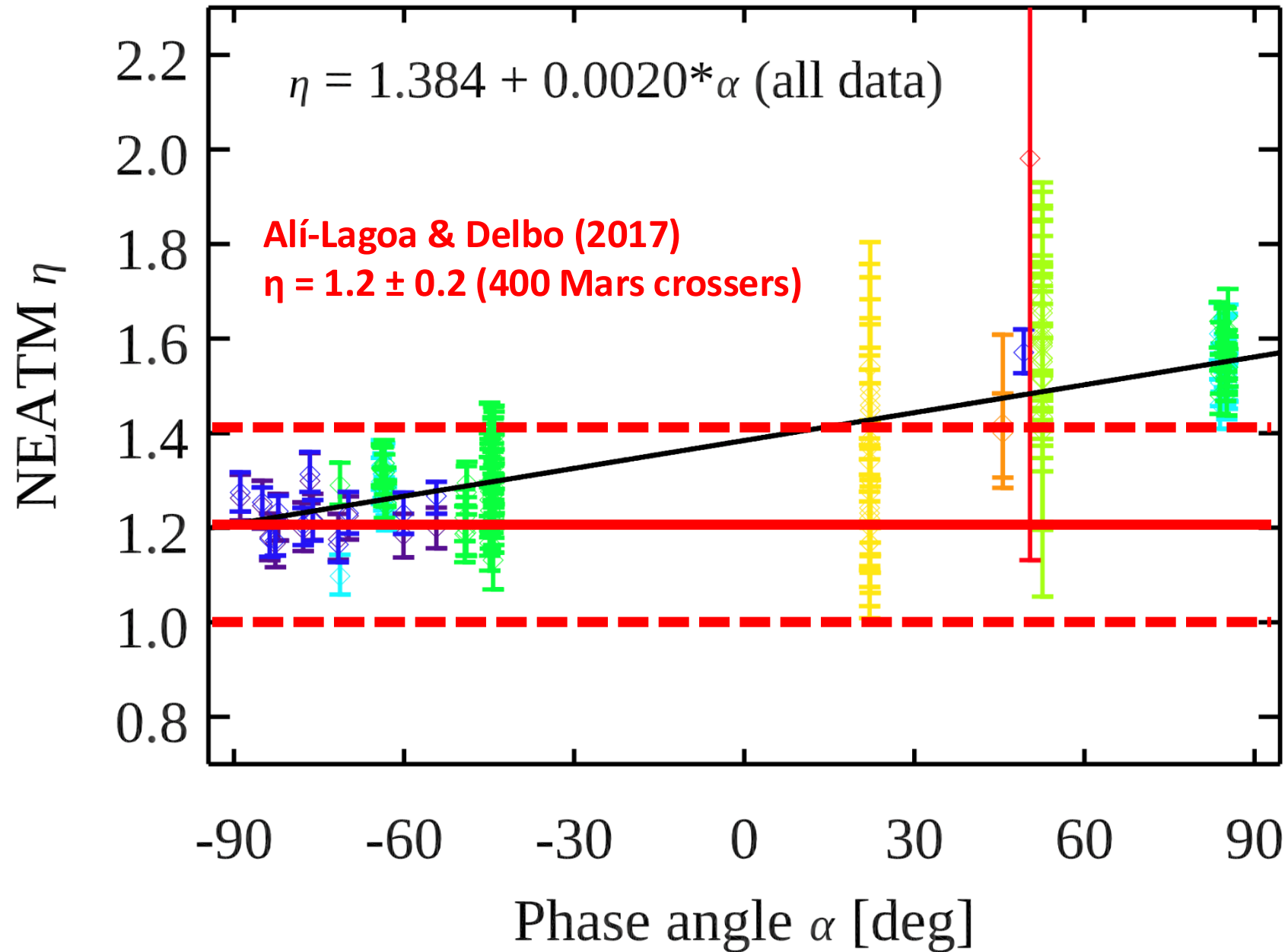
NEATM Beaming phase function



NEATM Beaming phase function



NEATM Beaming phase function



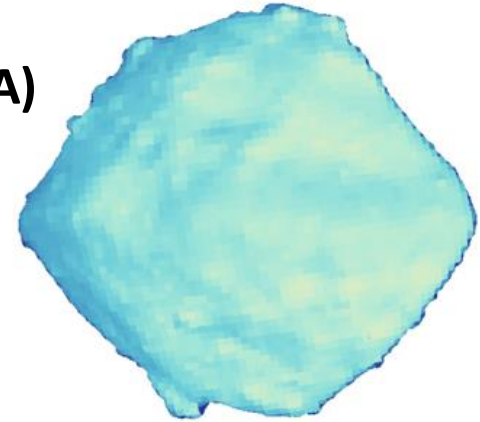
NEATM results: Ryugu

- Using NEATM with $\eta = (1.384 \pm 0.005) + (0.0020 \pm 0.0002)\alpha$ works very well and allows to reproduce Ryugu's size and albedo (**NOTE**: η calculation requires signed α !)
- This NEATM solution explains the available thermal measurements (within their absolute flux errors) over all phase angles ($-90^\circ \dots +90^\circ$), wavelengths (3.3 to 70 μm), rotational phases or heliocentric distances (0.98 ... 1.42 au), the $\chi^2_{reduced} = 1.2!$
- The NEATM η slope is a strong indication for a retro-grade rotator
- The η slope is also indicative of a thermal inertia close to 300 tiu (assuming a rotation period of 7.6 hours and an equator-on viewing geometry)
- Published η -solutions are not working very well for Ryugu (on only for specific phase-angle ranges and wavelengths)
- **But η is not only phase-angle dependent (with the slope depending on the object's thermal and rotation properties), it also changes with wavelengths and heliocentric distance (for the same object!)**

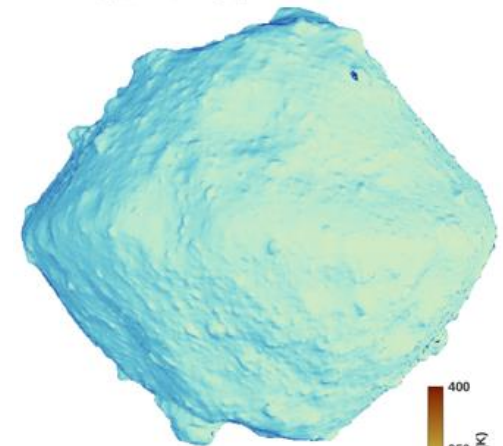
TPM application: Ryugu

(Lagerros 1996, 1997, 1998; Müller & Lagerros 1998, 2002, all in A&A)

- A rough surface model is assumed for each patch (considering only very small segments)
- The effect of heat conduction and surface roughness of the small segment is modelled and the result scaled to the entire patch
- Each small segment of a surface patch is divided into a large number of surface elements
- The **roughness** is simulated by either considering a **hemispherical crater** on otherwise flat surface, or by assuming a **Gaussian random rough surface** (see below)
- The Sun illuminates the segment and its surface elements and moves across the sky as the asteroid revolves around its axis.
- Due to the surface roughness, some elements **shadow** other surface elements which is taken into account
- Each element is heated by the Sun and by visual **light scattered** from neighbouring elements.
- Heat is exchanged with the interior through heat conduction, assuming the elements to be 1-D slabs isolated from neighbouring surface elements
- For temperature calculations the bolometric emissivity ϵ_{bolo} is relevant (integral of the spectral emissivity weighted by the solar spectrum)
- The (disk-integrated) flux (or emittance) is calculated by integrating over all temperatures and by considering the hemispherical spectral emissivity ϵ_{spec} at a given wavelengths (depends on T^4 !)

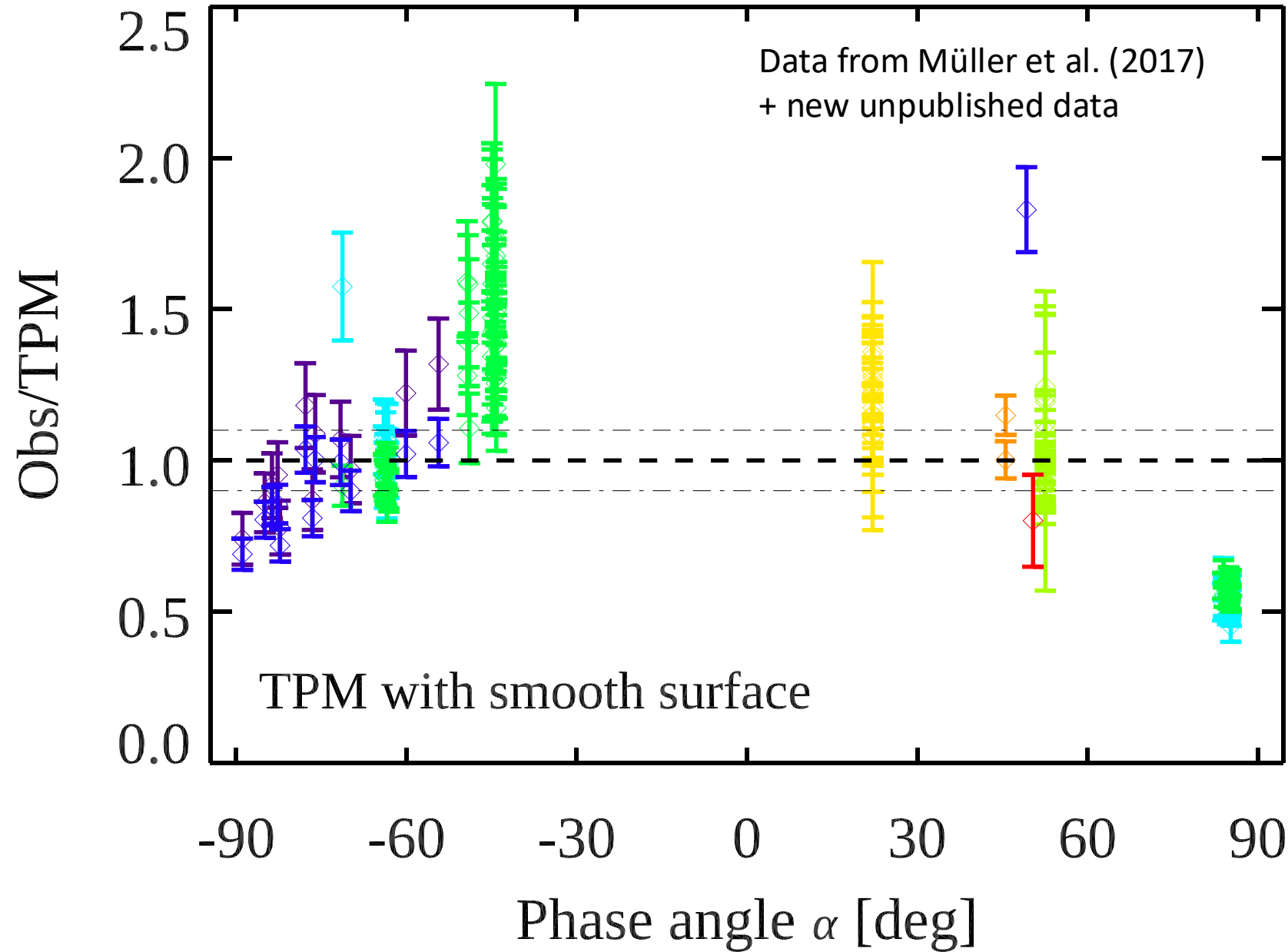


Hyb2_tir_20181114_073228_l3:
 $r=1.335$ au, $\Delta=19.23$ km, $\alpha = -4.8^\circ$,
 $(\lambda, \beta)_{\text{SubSC}} = (356.4, -1.3)$,
 $(\lambda, \beta)_{\text{SubSolar}} = (351.6, -1.9)$



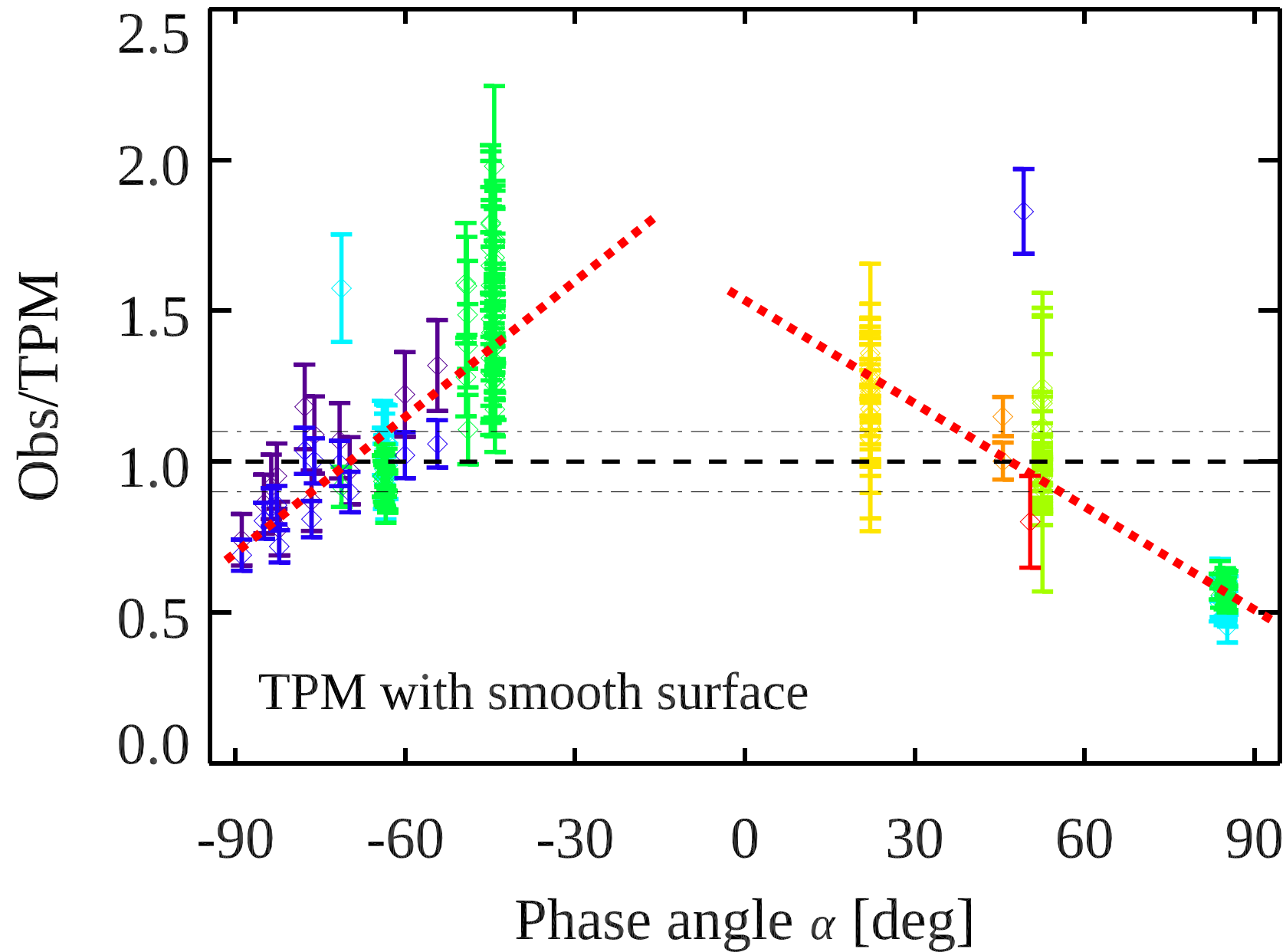
TPM with $\Gamma=350$ tiu,
 $\text{rms} = 45^\circ$ (sph. crater),
 $\epsilon = 0.98$, SPC_200k

TPM roughness effects $f(\alpha)$



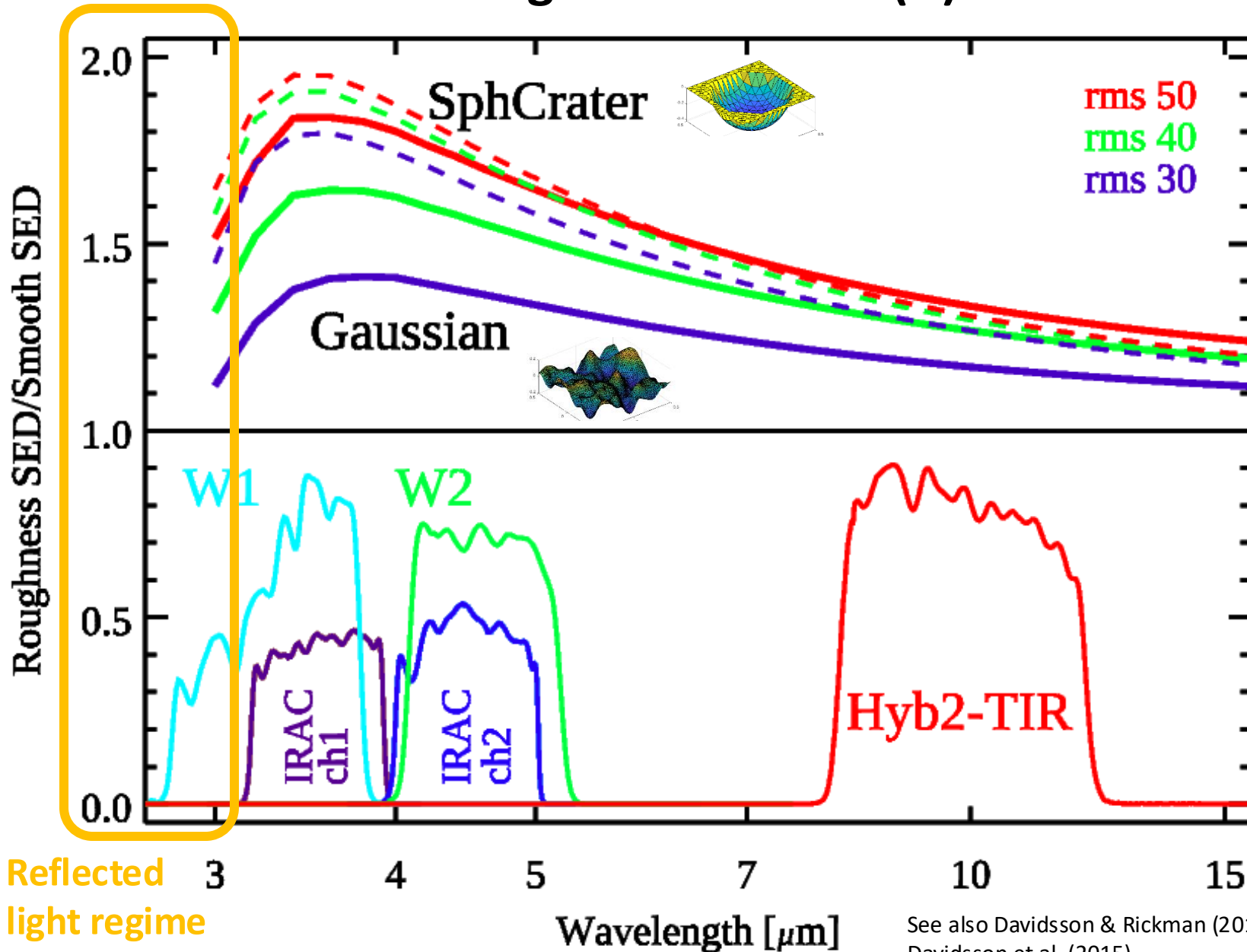
Using the high-resolution shape model and the published thermal inertia for Ryugu

TPM roughness effects $f(\alpha)$



Using the high-resolution shape model and the published thermal inertia for Ryugu

TPM roughness effects $f(\lambda)$

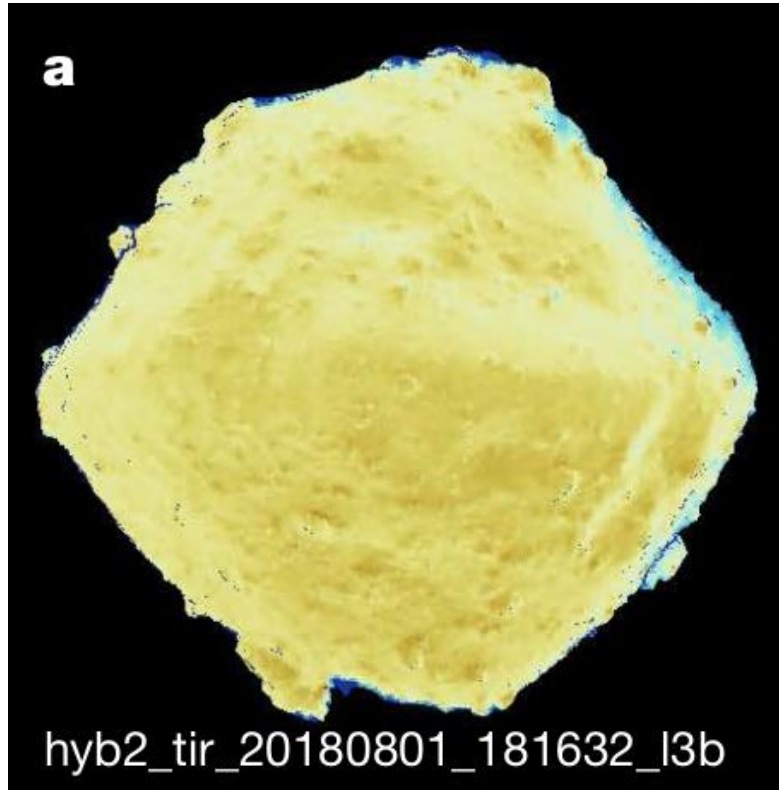


Reflected
light regime

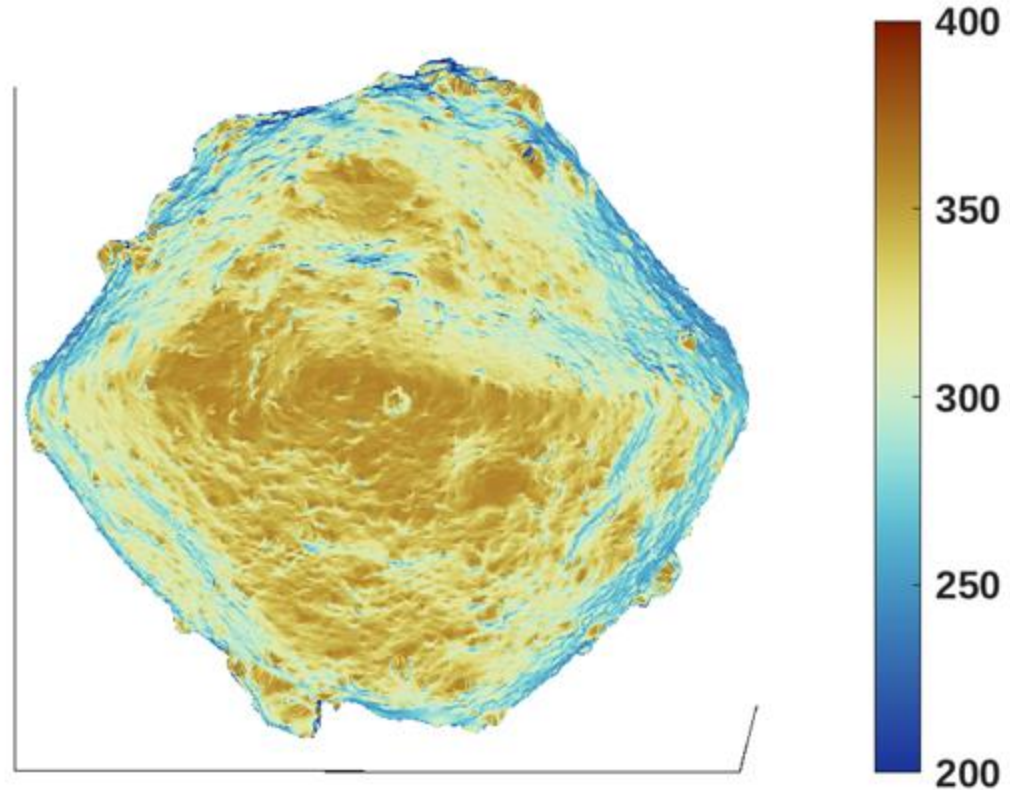
See also Davidsson & Rickman (2014);
Davidsson et al. (2015).

Using the asteroid thermophysical model code by Lagerros (1996, 1997, 1998)
Müller & Lagerros (1998, 2002, all in A&A), no surface roughness added

TIR image from Okada+20



TPM prediction ($\Gamma=225$)

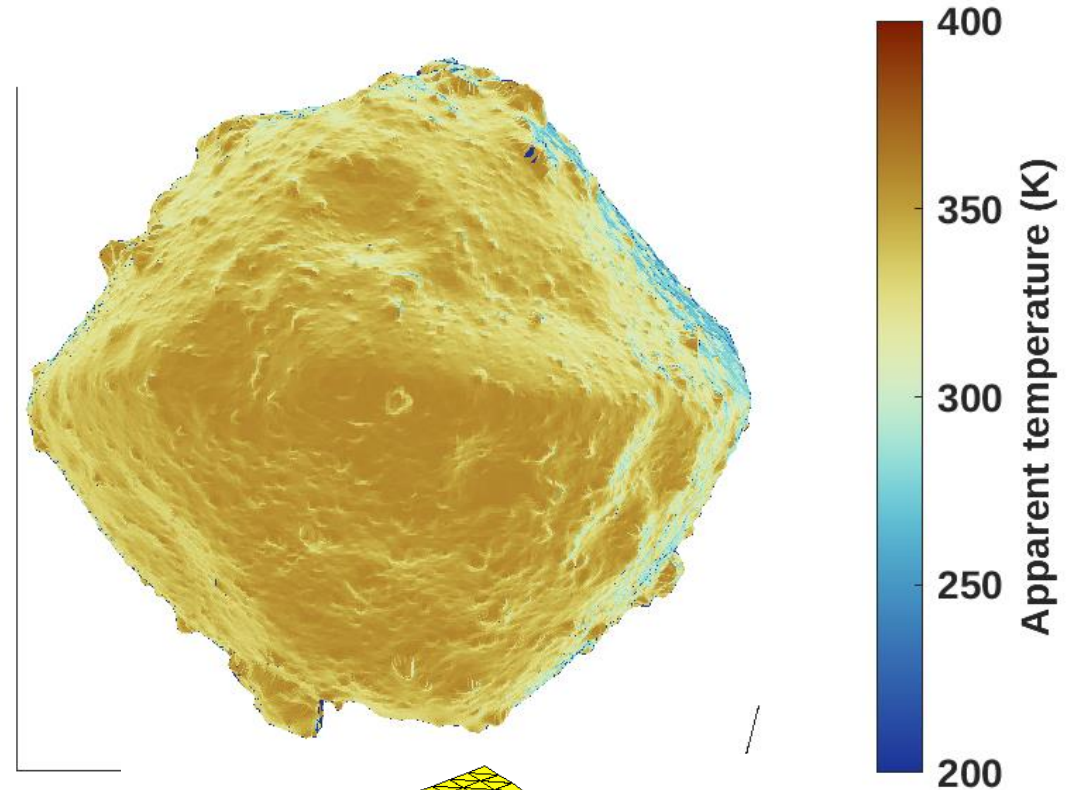
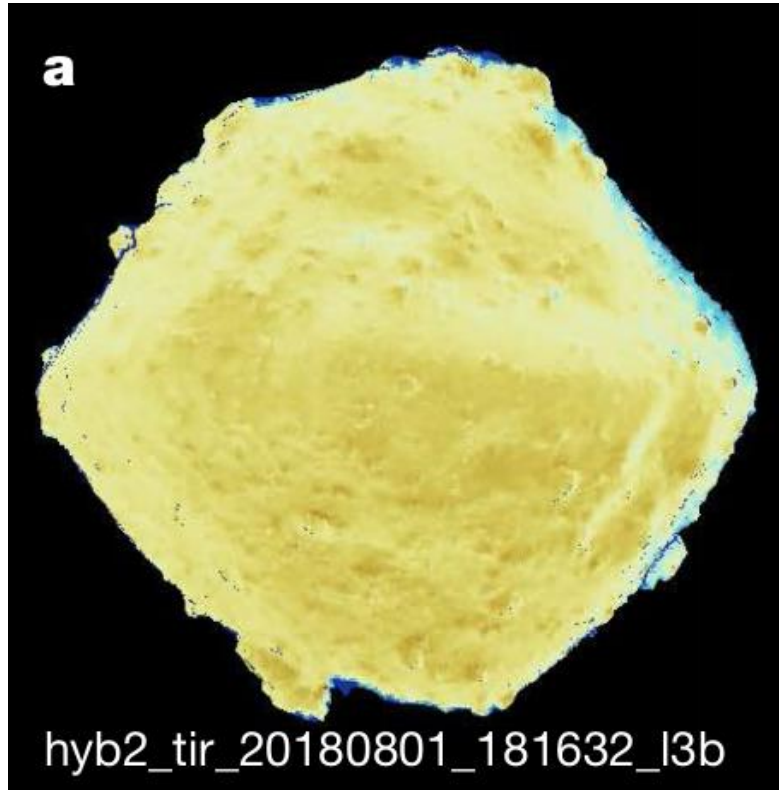


TPM roughness effects

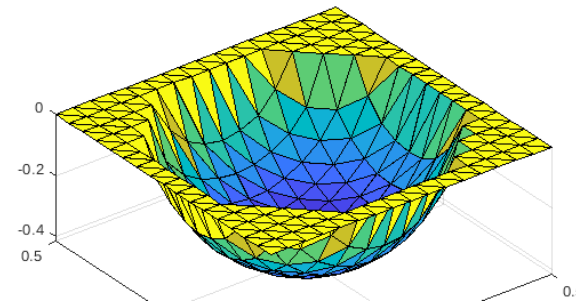
Using the asteroid thermophysical model code by Lagerros (1996, 1997, 1998)
Müller & Lagerros (1998, 2002, all in A&A), adding surface roughness
(hemispherical craters on otherwise flat surface)

TIR image from Okada+20

TPM prediction ($\Gamma=225$, rms 47°)



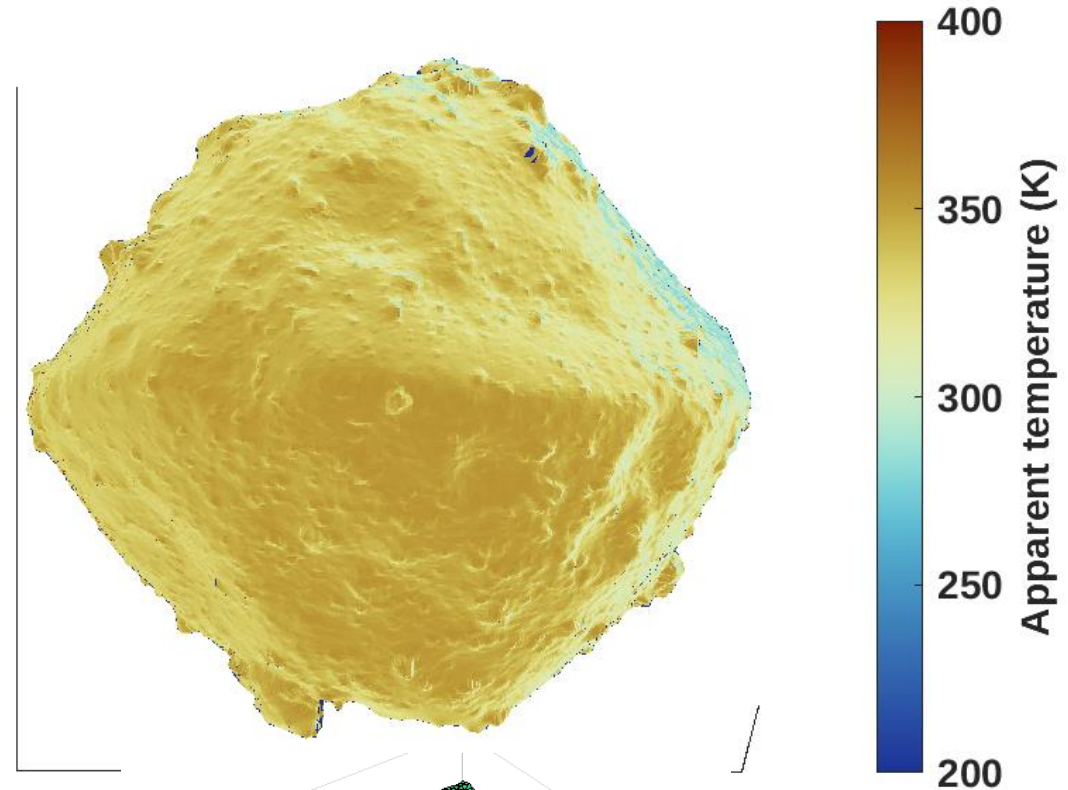
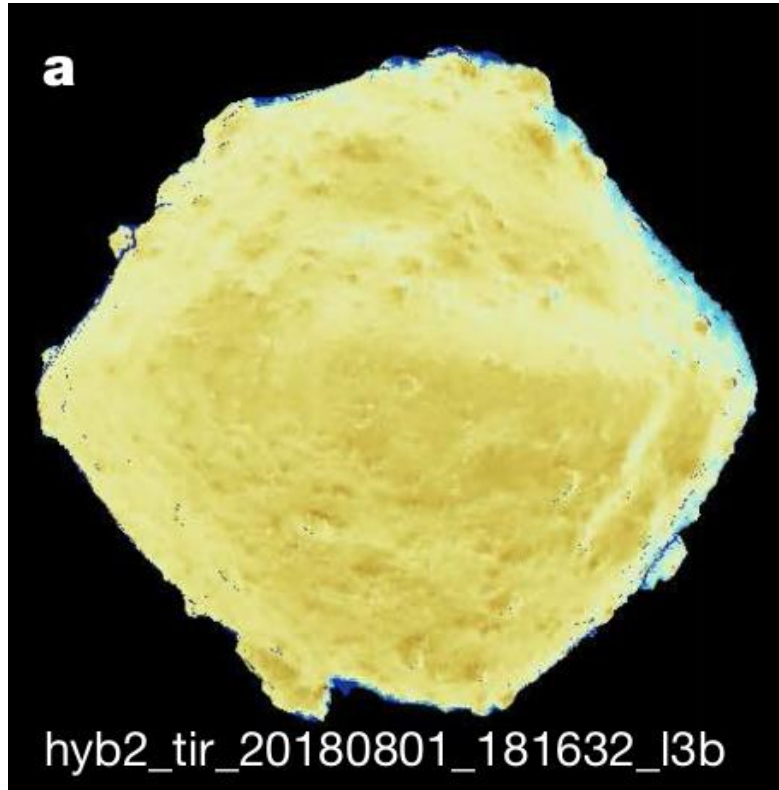
TPM roughness effects



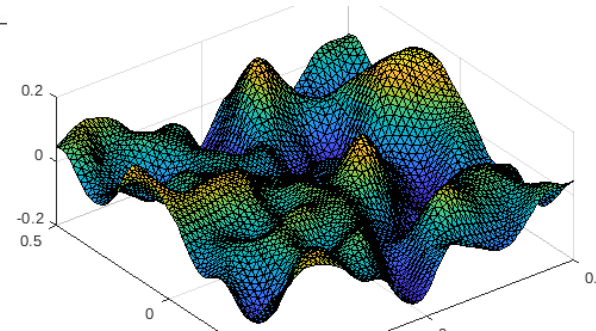
Using the asteroid thermophysical model code by Lagerros (1996, 1997, 1998)
Müller & Lagerros (1998, 2002, all in A&A), adding surface roughness
(assuming Gaussian random rough surface)

TIR image from Okada+20

TPM prediction ($\Gamma=225$, rms 40°)



TPM roughness effects



TPM results: Ryugu

- TPM solution ($\epsilon_{\text{bolo}} \sim 0.98$, TI = 300-400 tiu, roughness: 30-50° rms of surface slopes) **explains the remote data** very well, including before/after opposition effects, short-wavelength data, thermal lightcurves, amplitudes, absolute fluxes, ... AND allows to **explains the surface temperature distribution** (TIR images) in a qualitative way
- Roughness modeling (hemispherical craters or Gaussian random surfaces) is crucial, but results, for exactly the same rms of surface slopes, are not identical
- Thermal inertia as a function of temperature (or r_{helio}) is noticeable for Ryugu in-situ data

Chelyabinsk progenitor (ChPG) orbit (study in the context of NEOMIR project)

- Based on Popova et al. (2013): “*Chelyabinsk Airburst, Damage Assessment, Meteorite Recovery, and Characterization*“, Science 342, 6162
- Orbital elements: $a=1.76 (\pm 0.16)$ au, $e=0.581 (\pm 0.018)$, $i=4.93^\circ (\pm 0.48^\circ)$, $q=0.739 (\pm 0.020)$ au, $T_p=2012-12-31.9 (\pm 2.0)$; Impact: 2013-02-15 03:20 UTC
- 20-m size assumption ($\rightarrow H=27.4/26.2/25.7$ mag for $p_v=0.05/0.15/0.25$)

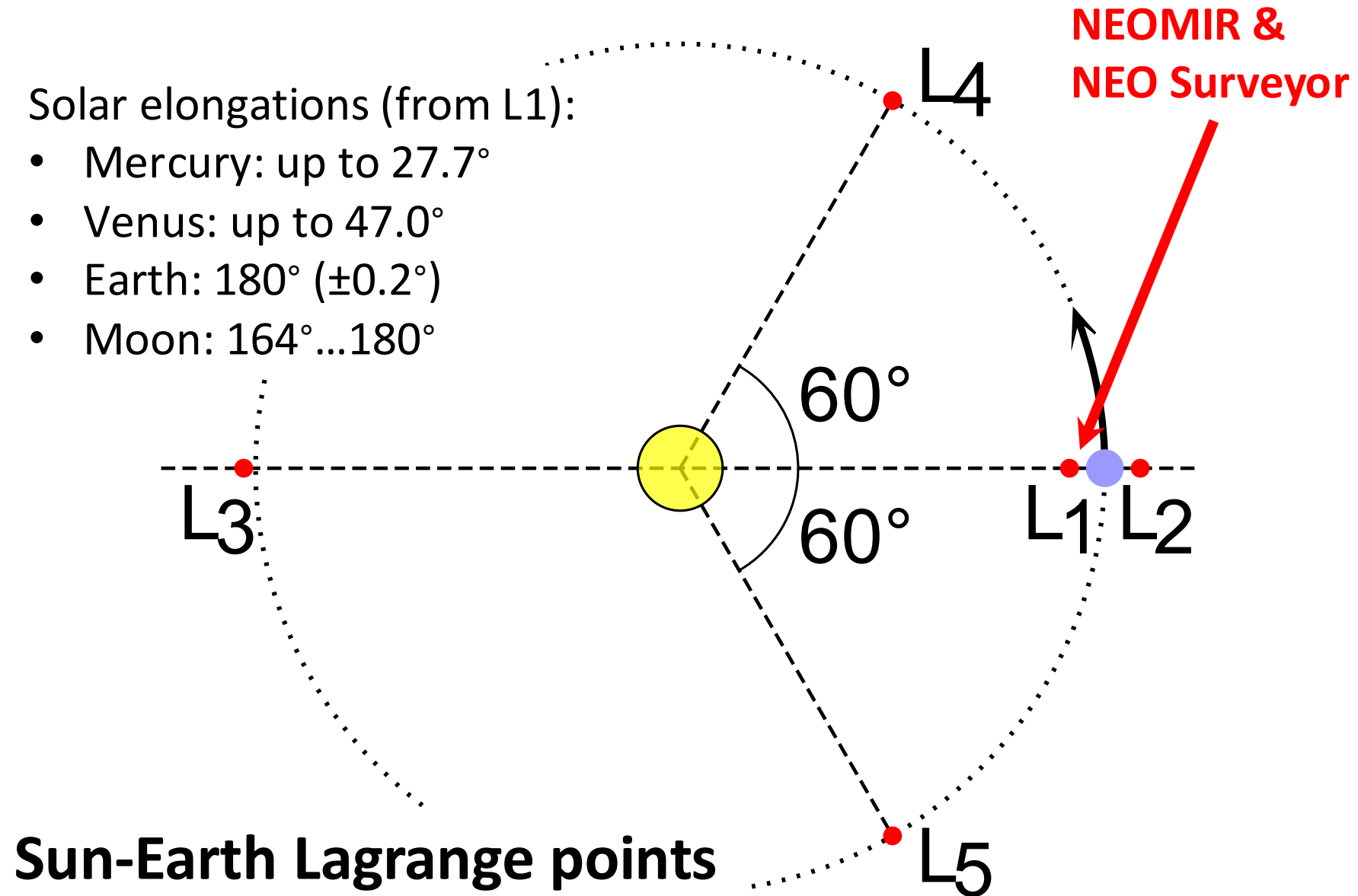
Questions:

- What's the best place to observe (L1,L2,L4,L5, Earth)? \rightarrow L1
- Which wavelengths and why? \rightarrow mid IR (8-12 μm) produces highest SNR
- Is the background a problem? \rightarrow yes, certainly at solar elongations $<60^\circ$
- Is the high apparent motion a problem? \rightarrow yes, in the days before Earth encounter
- What's needed to detect (SNR >5) the ChPG as early as possible? \rightarrow close Sun-proximity observations (down to solar elongation of about 20°); fast detector readout (to avoid saturation) and synthetic tracking techniques (to take advantage of full array-crossing times)
- Where are the problems in the SNR estimates? \rightarrow NEA model calculations

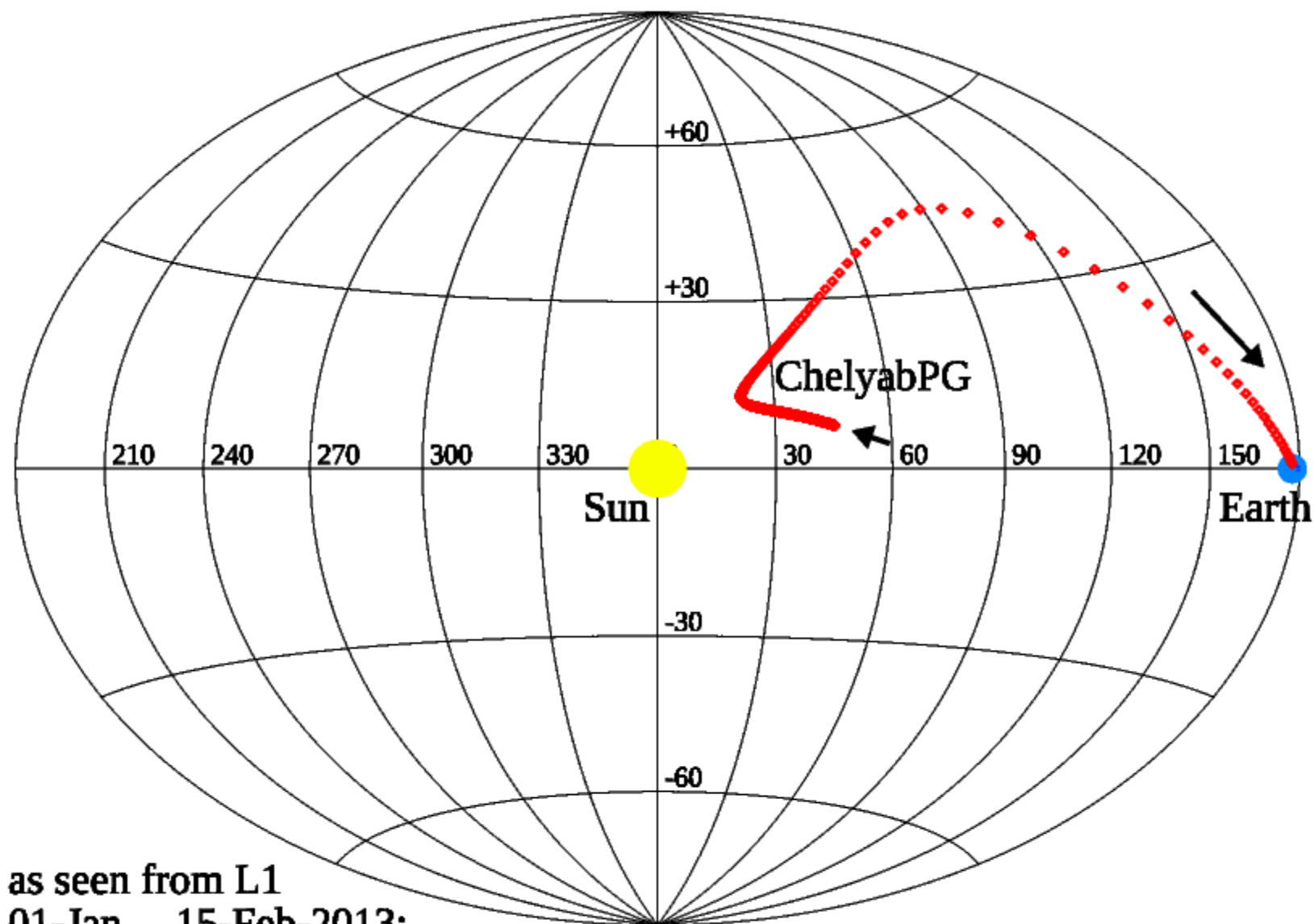
Which is the best place to observe? → L1

Solar elongations (from L1):

- Mercury: up to 27.7°
- Venus: up to 47.0°
- Earth: 180° ($\pm 0.2^\circ$)
- Moon: $164^\circ \dots 180^\circ$

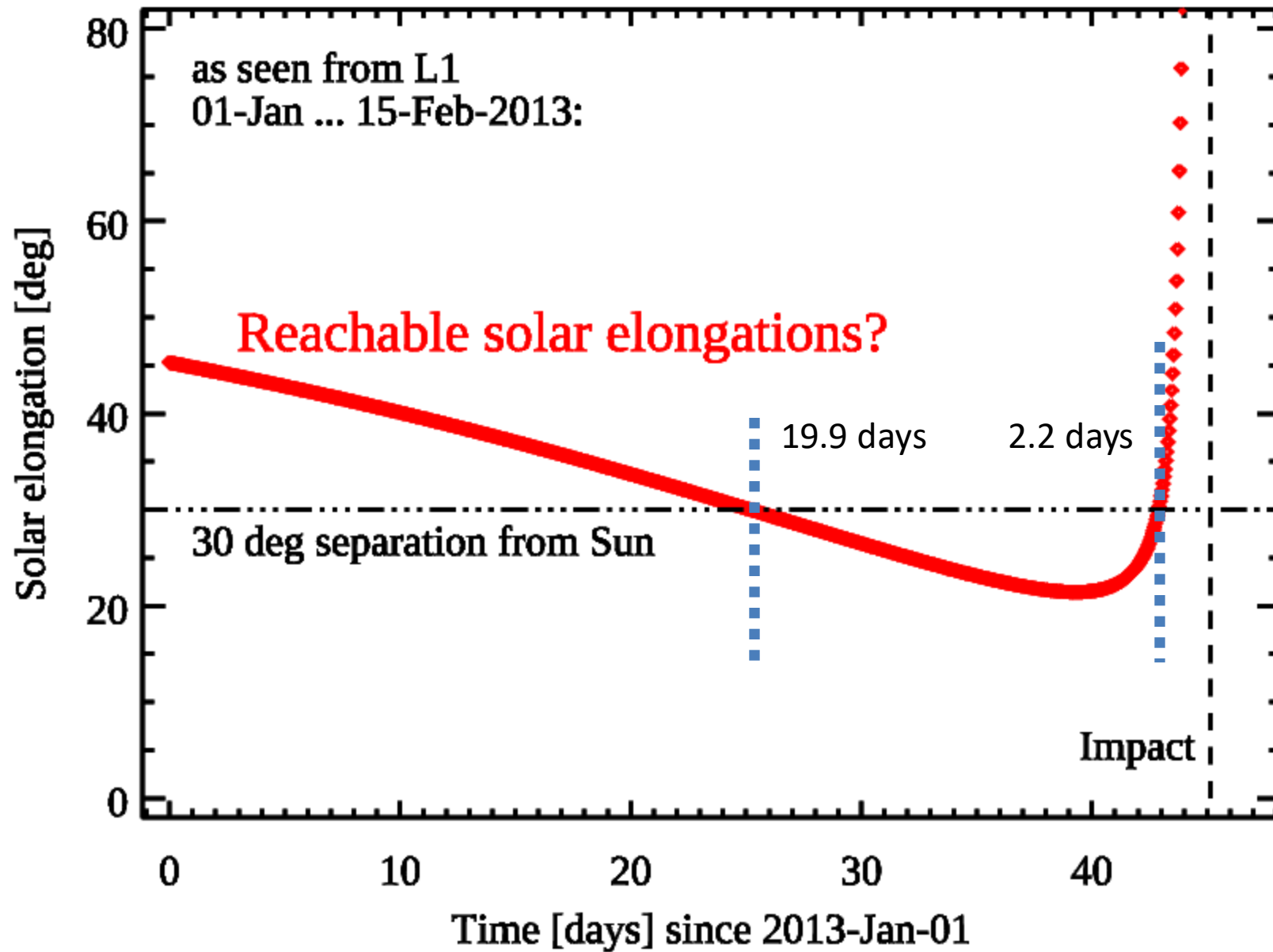


Sun-Earth Lagrange points

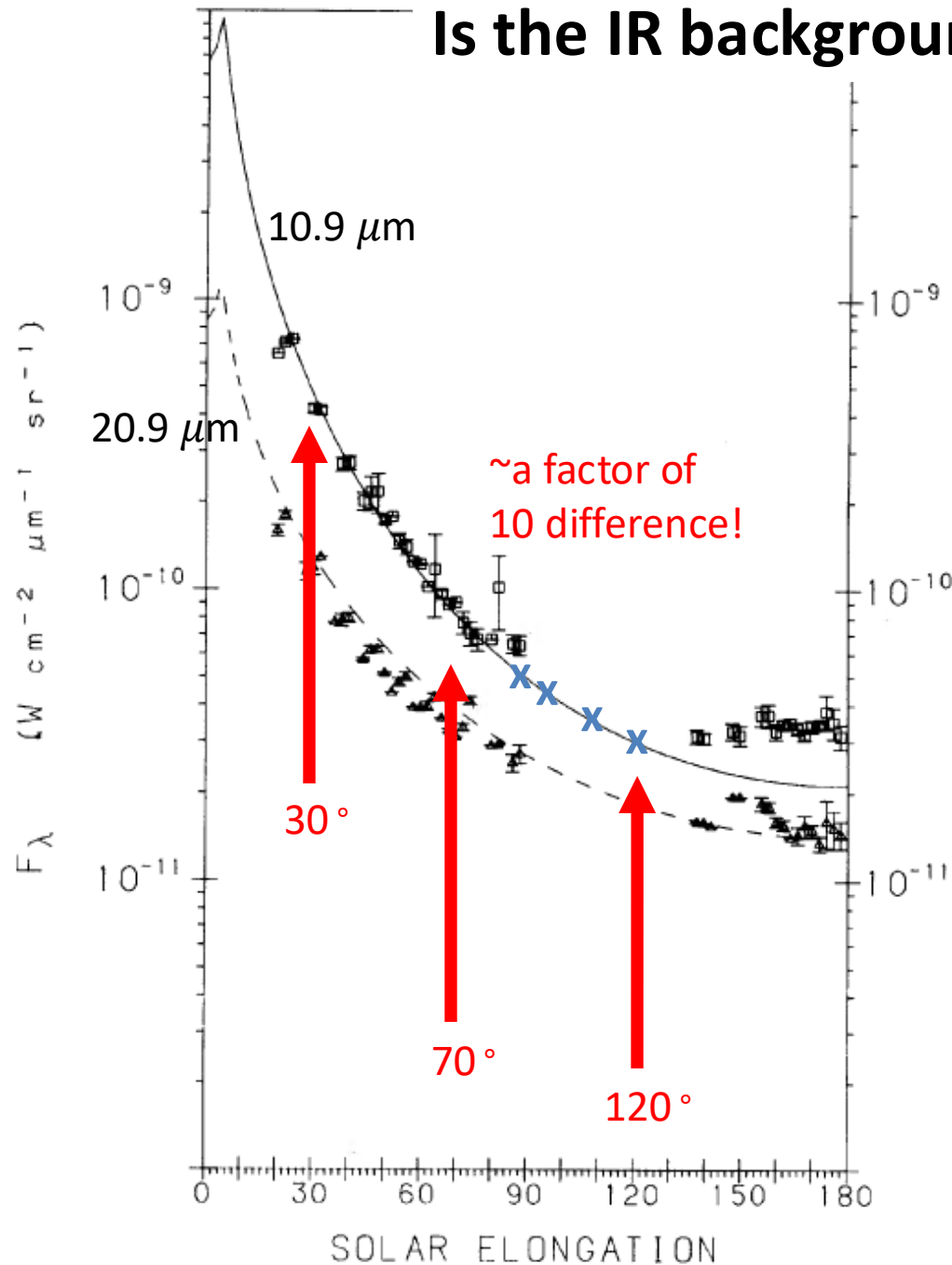


as seen from L1
01-Jan ... 15-Feb-2013:
1-h time resolution

Aitoff Projection
T.Mueller, Aug 24



Is the IR background a problem? → Yes



Murdock & Price (1985),
rocket experiment, Fig. 12:
 F_λ as a function of solar
elongation angle at 10.9 μm
(squares) and 20.9 μm
(triangles) in ecliptic plane.

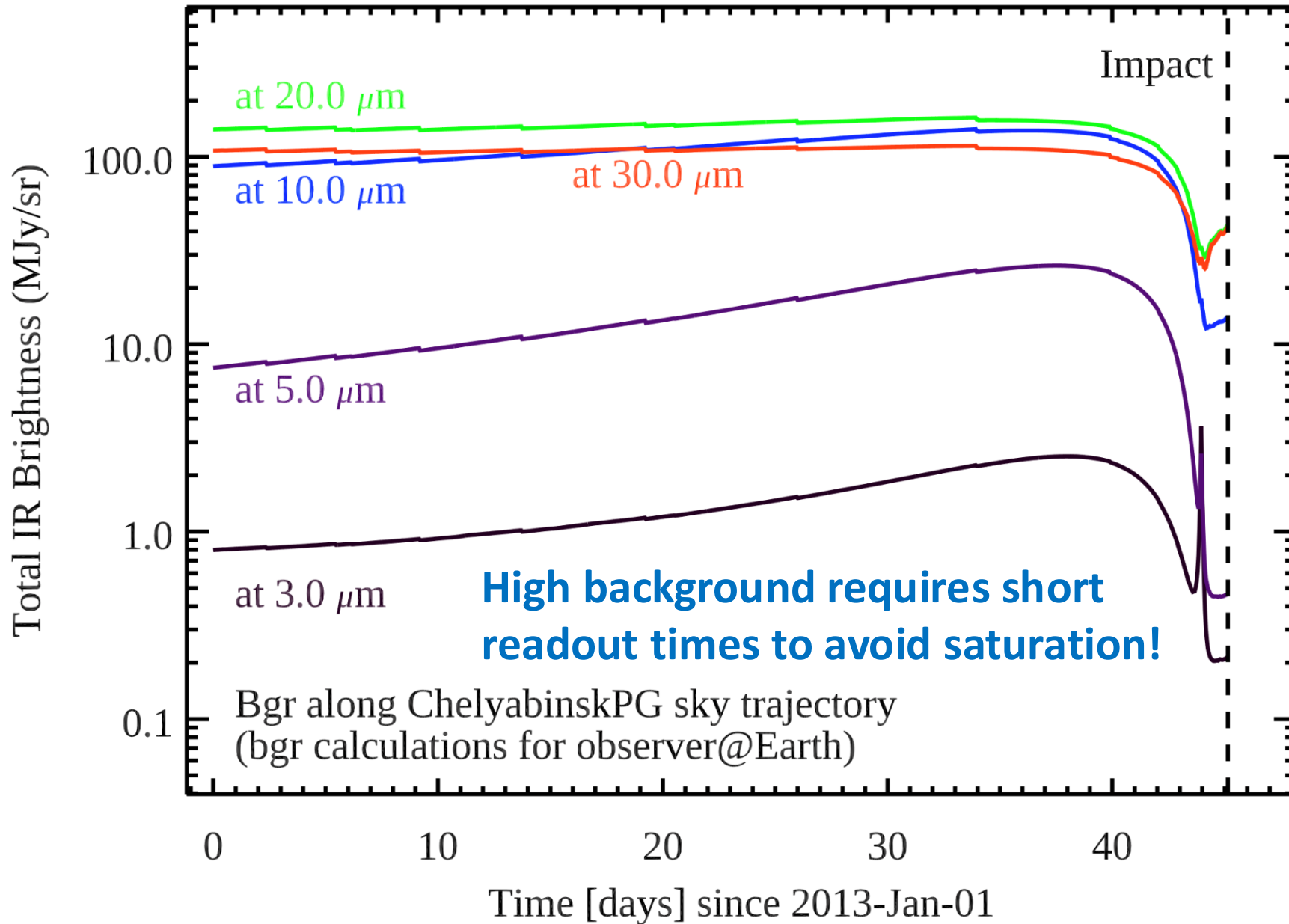
Values from the JWST ETC at
10 μm in the ecliptic plane:

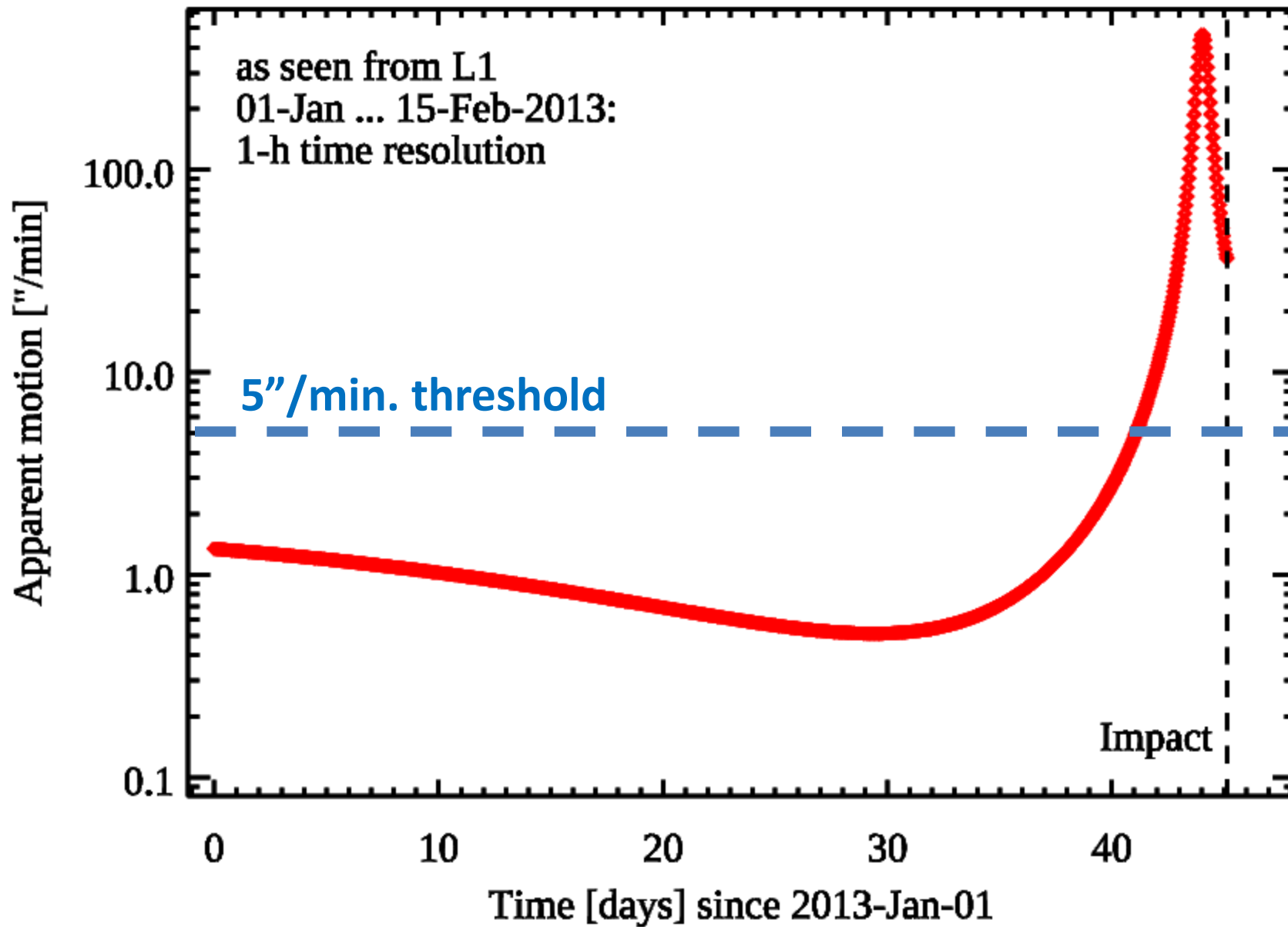
$(\lambda - \lambda_{\text{Sun}})_{\text{ecI}}$	SurfBrightness
85°	40 MJy/sr
90°	37 MJy/sr
100°	31 MJy/sr
120°	25 MJy/sr

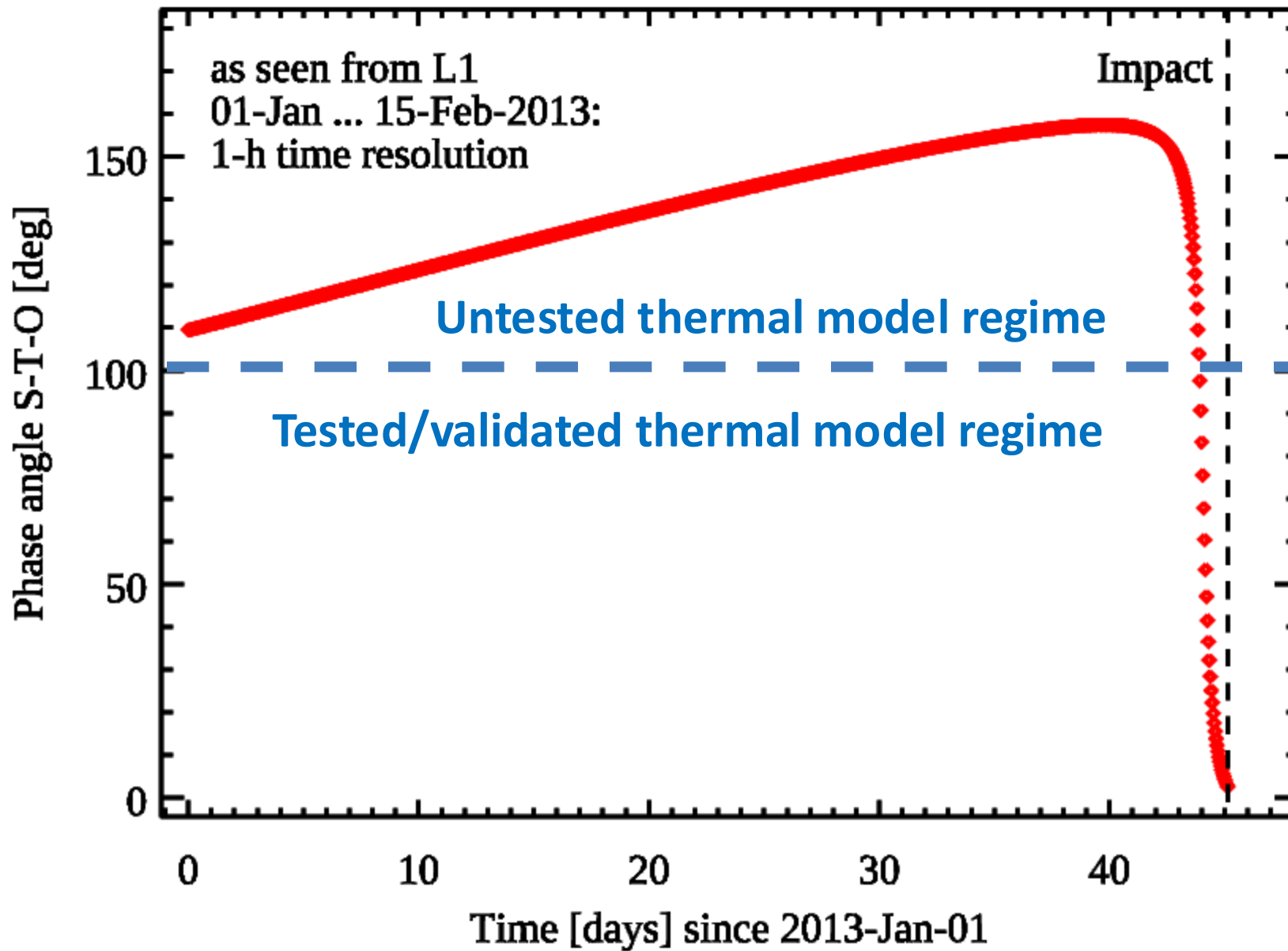
Another tool available at:

[https://irsa.ipac.caltech.edu/
applications/BackgroundModel](https://irsa.ipac.caltech.edu/applications/BackgroundModel)

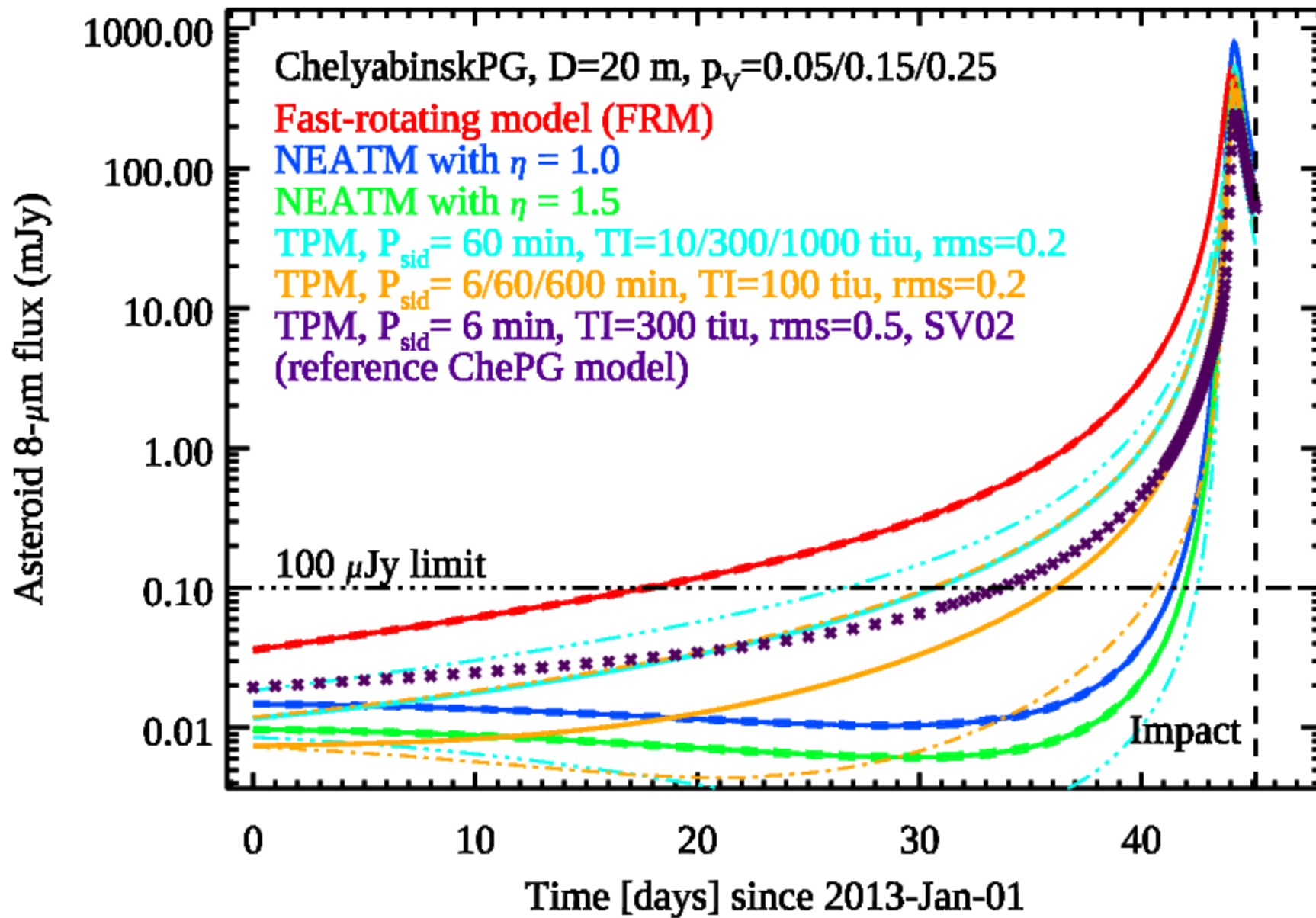
IR Background along apparent trajectory







Different model predictions at 8 μm



Pre-impact detection of Chelyabinsk-type objects in the thermal infrared

A pre-impact detection (and size estimate) of a Chelyabinsk-type object is possible:

- with a 50 cm telescope, large FOV, passively cooled detector
- from L1, at 8-10 μm , 8-10 days before impact (while the object is still fainter than mag 30!)
- but requires observations down to 20° solar elongation, at very high background, produces high data rates and needs synthetic tracking techniques (to take advantage of full array-crossing times)!
- Large uncertainties in the predictions remain due to unknowns in the asteroid IR models!

Difficulty to model IR emission of small asteroids at high phase angles

- The thermal model predictions for (small, fast-rotating?, monolithic?, porous?) NEOs is very difficult and uncertain, especially for high phase angles ($> 90^\circ$)
- There is clearly a lack of IR measurements of asteroids seen under large phase angles!
- NEOs below ~ 200 m in size are rotating faster (Pravec et al. 2008): \rightarrow FRM?
- However, NEATM (with beaming parameters η in the range 1-1.5) seem to work for about 50 NEOs with sizes between ~ 8 to ~ 100 m (Mainzer et al. 2014): \rightarrow NEATM?
- The Yarkovsky drift of a rapidly rotating small asteroid points to an unexpectedly low thermal inertia, indicative for a highly porous or cracked surface (Petkovic et al. 2021; Fenucci et al. 2023): \rightarrow TPM?
- There is also the vector alignment of asteroid spins by thermal torques (YORP), see work by Vokrouhlicky et al. (2003): \rightarrow TPM?
- Our baseline model: TPM with $D = 20$ m, $p_v = 0.15$, $\Gamma = 300$ tiu, roughness rms = 0.5, $P_{\text{sid}} = 6$ min, $\beta_{\text{ecl}} = +45^\circ$, $\varepsilon = 0.9$ (fluxes between NEATM with $\eta = 1.0$ or 1.5 and FRM for wide phase angle ranges)

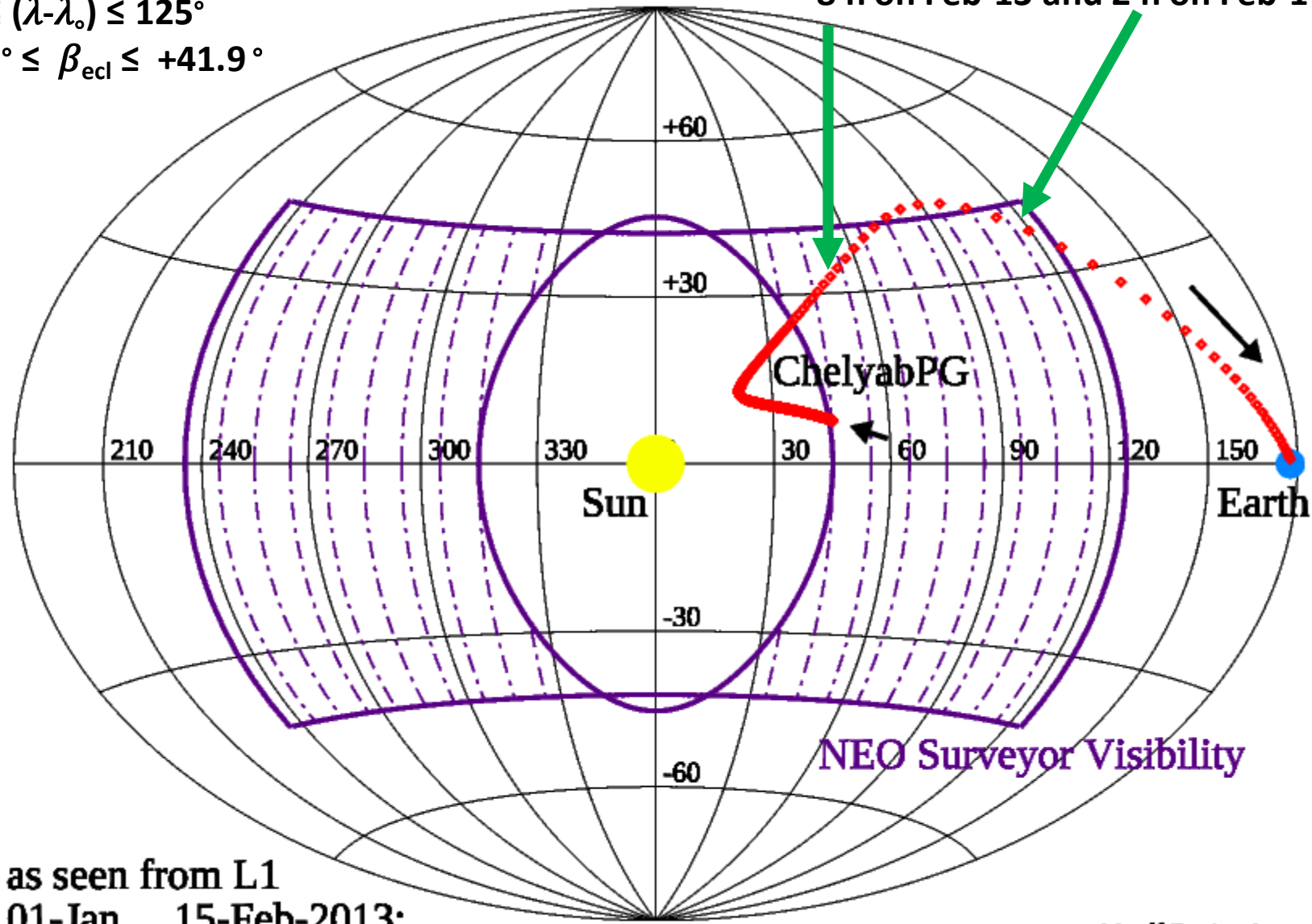
NEOSurveyor visibility zone:

$$45^\circ \leq (\lambda - \lambda_0) \leq 125^\circ$$

$$-41.9^\circ \leq \beta_{\text{ecl}} \leq +41.9^\circ$$

Detectability: about 1 day 16 h before impact

8 h on Feb-13 and 2 h on Feb-14

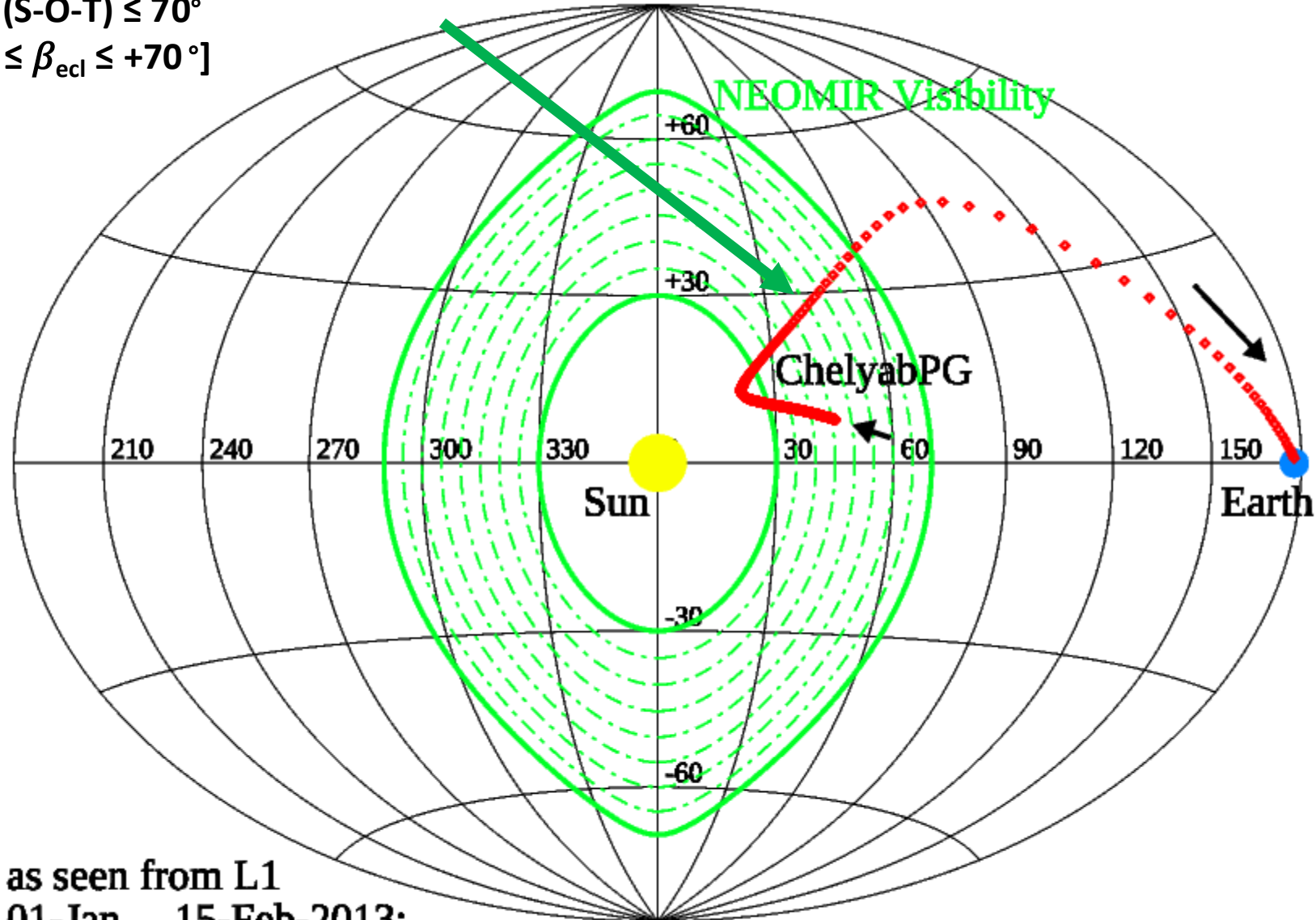


as seen from L1
01-Jan ... 15-Feb-2013:
1-h time resolution

Aitoff Projection
T.Mueller, Aug 24

NEOMIR visibility zone:
 $30^\circ \leq (S-O-T) \leq 70^\circ$
 $[-70^\circ \leq \beta_{\text{ecl}} \leq +70^\circ]$

Detectability: about 2 days 6 h before impact
for about 24 hours



as seen from L1
01-Jan ... 15-Feb-2013:
1-h time resolution

Aitoff Projection
T.Mueller, Aug 24

Different model predictions at 8 μm

